

Plasma Control of EM Waves/High Energy Density Plasma Jets

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Stanford University



Michigan Institute for Plasma Science and Engineering (MIPSE)
Seminar Series: March 28, 2018

Acknowledgements

- AFOSR - MURI
- DOE (NNSA-SSAP)
- Collaborators: L. Raja, T. Ito
- Graduate students

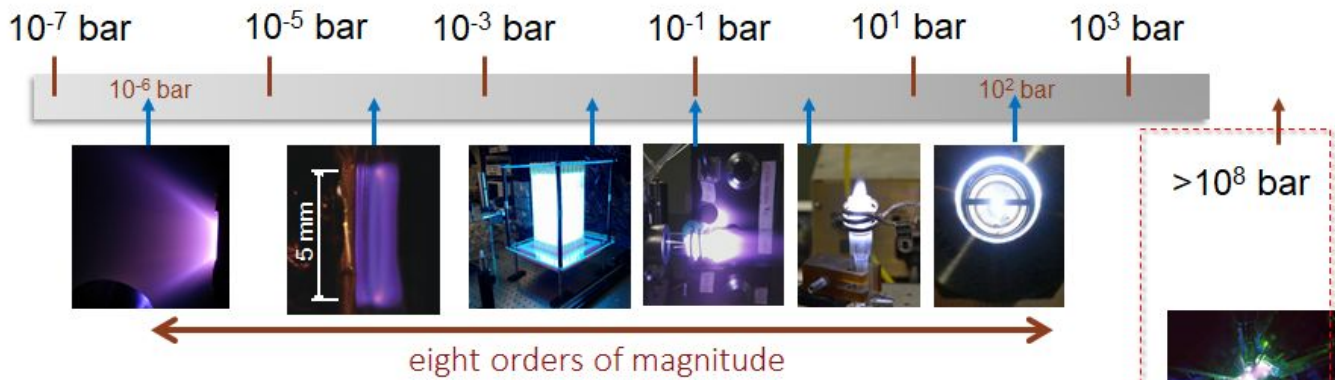


Go Blue!

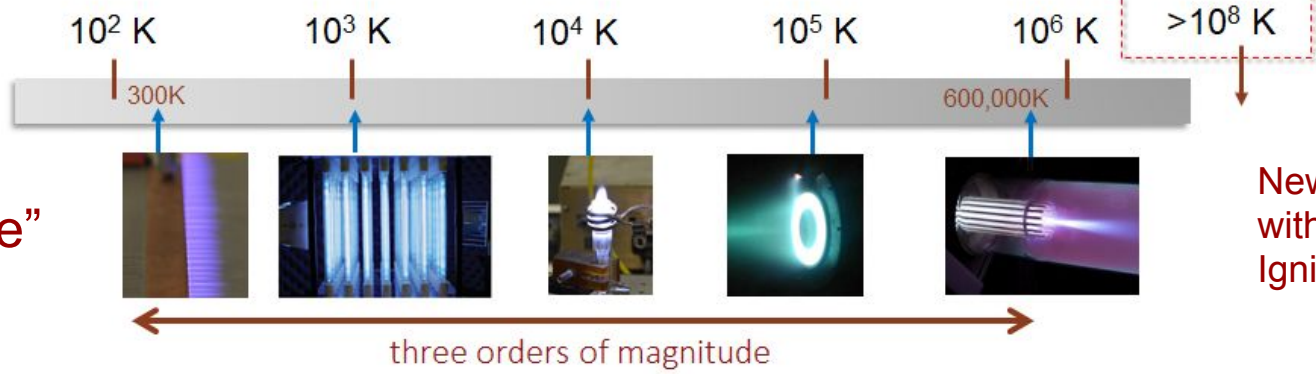


Expose Students to a Wide Range of Conditions

Pressure



Ion
"Temperature"



New Collaboration
with the National
Ignition Facility

Research Interests

Plasma Aerophysics

Space Plasma
Propulsion

Plasma-Assisted
Combustion

Plasma
Aerodynamics



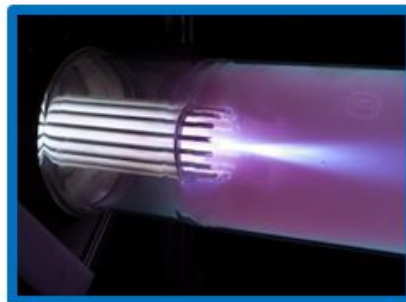
Plasmas in Extreme Conditions

High Pressure
Discharge Physics
and Chemistry

T. Yong

High Energy Density
Plasma Physics

W. Riedel
T. Underwood



EM Wave-Plasma Interactions

Plasma
Metamaterials

R. Colon
R. Lee

Plasma Photonic
Crystals

B. Wang
F. Righetti

Plasma Control of
Guided
Microwaves

A. Marcovati
D. Biggs



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MIPSE 2012



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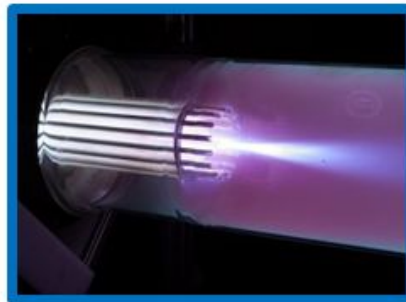
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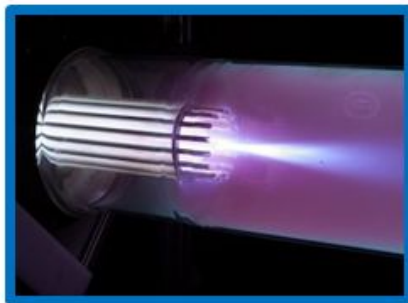
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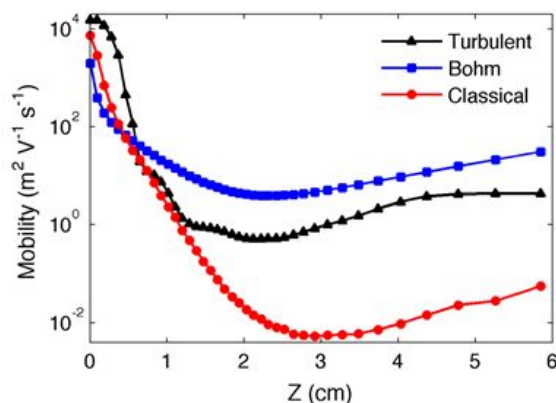
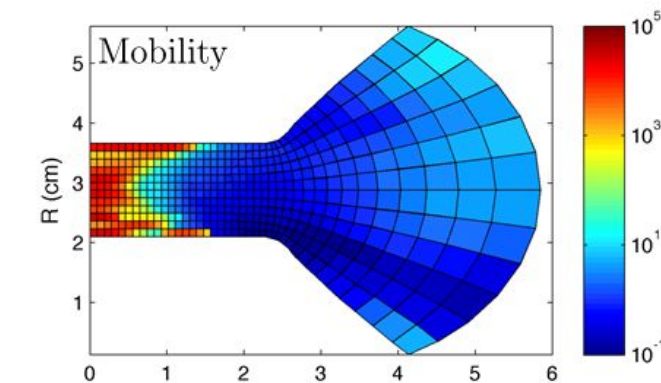
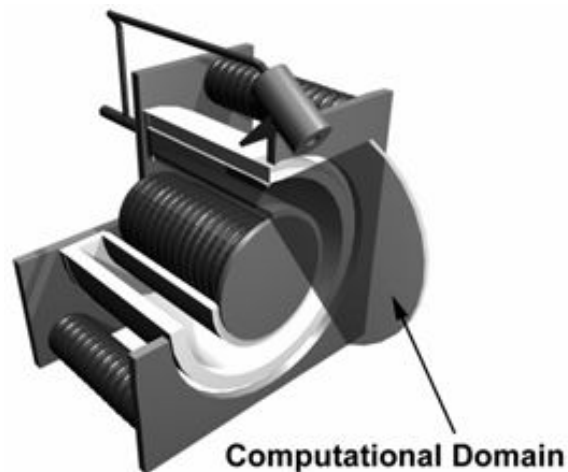
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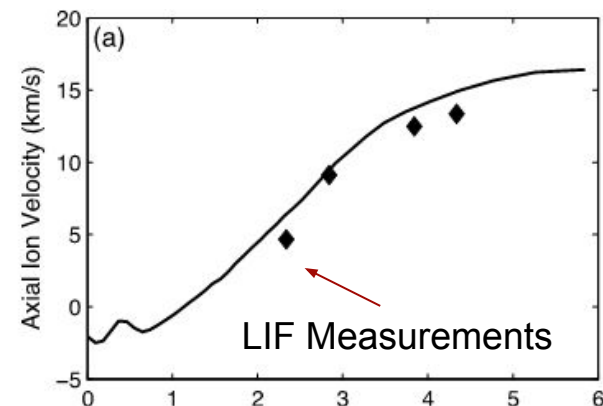


Hall Thrusters

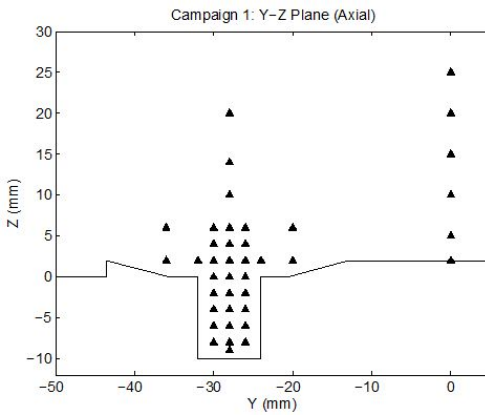
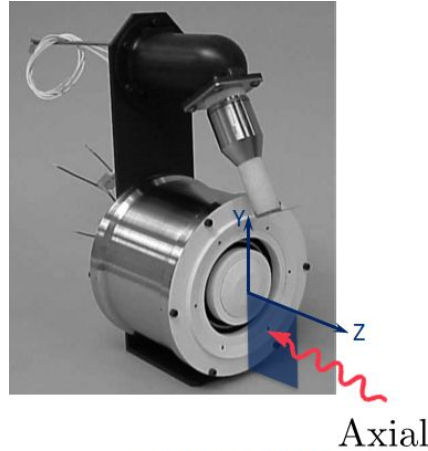
Electron Transport Models for Hall Thruster Simulations



$$\mu_{\text{eff}} = \left(\frac{|J_{e\perp} E_{\perp}|}{K n_e e c_e^2 B^3} \right)^{1/2}$$

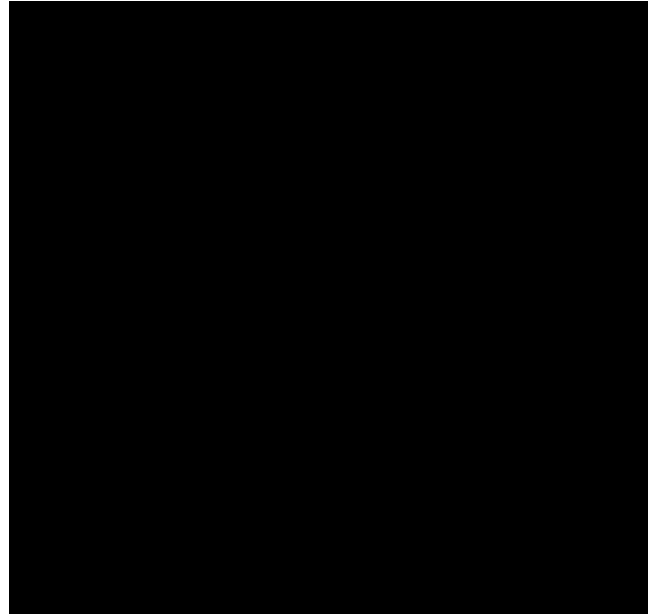


Hall Thruster Diagnostics

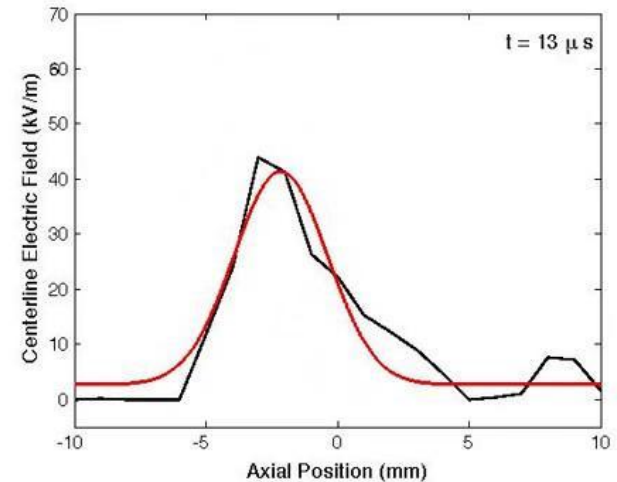


Time-Current Synchronized LIF Velocimetry

Axial Velocities within Channel



Inferred Electric Field



$$E_z = \frac{m_i}{e} \left(\frac{\partial v_z}{\partial t} + v_z \cdot \frac{\partial v_z}{\partial z} \right)$$

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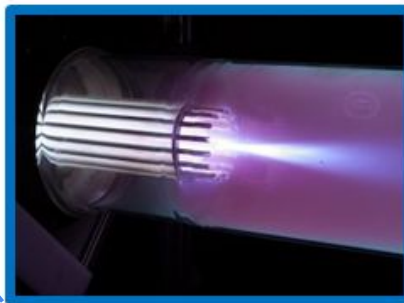
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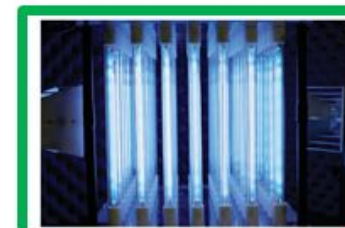
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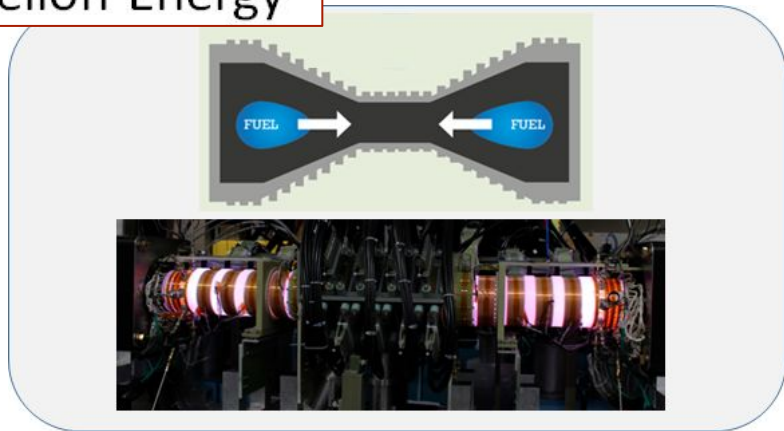
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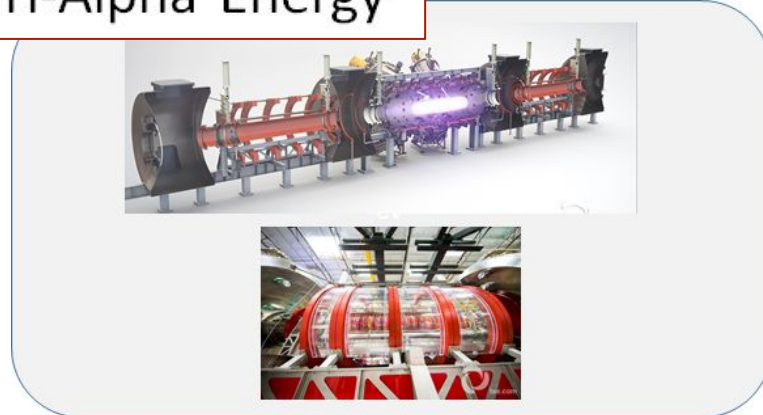
Fusion Quo Vadis? High Energy Density Plasma Jets?

Helion Energy

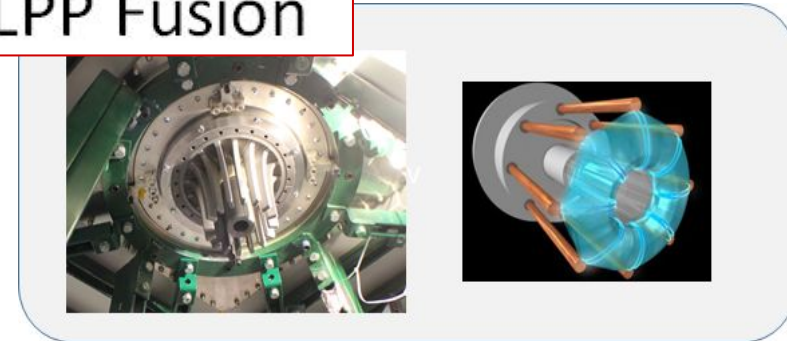


- Colliding/converging and magnetically-compressed plasma jets
- conditions intermediate to MCF and ICF
 - $n_e = 10^{17} - 10^{18} \text{ cm}^{-3}$ ($10^{23} - 10^{24} \text{ m}^{-3}$)
 - $T = 10^6 - 10^7 \text{ K}$
 - $T \sim 0.1 - 5 \text{ ms}$

Tri-Alpha Energy

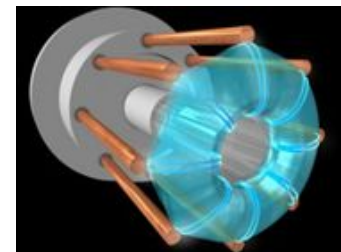
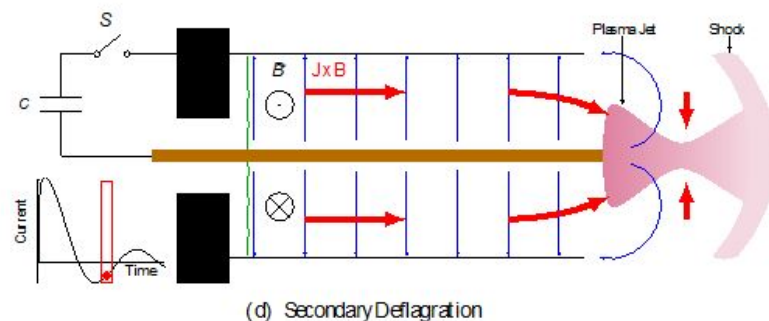
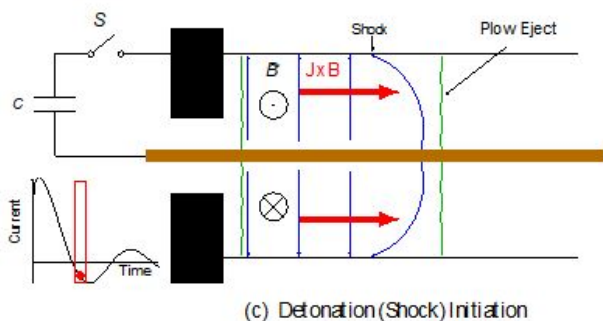
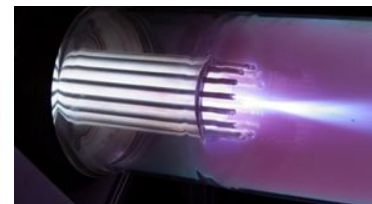
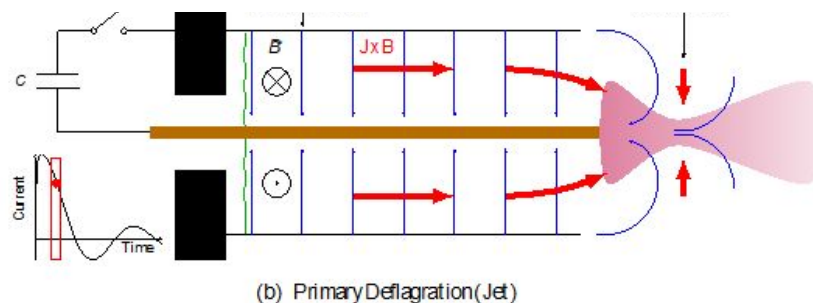
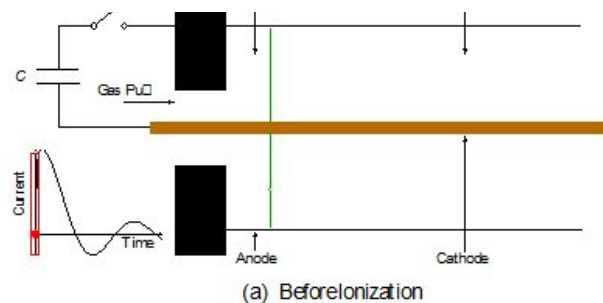


LPP Fusion



Deflagration Plasma Accelerator (Mathers-type)

- two branches of operation (Loebner et al, PRL 2015)
 - smooth expansion into vacuum or shock
- distinguished by prefill or vacuum fill initial conditions

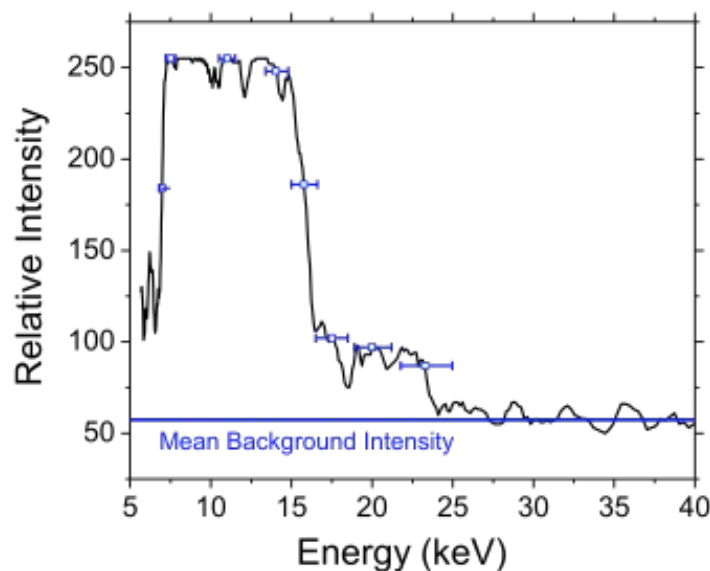
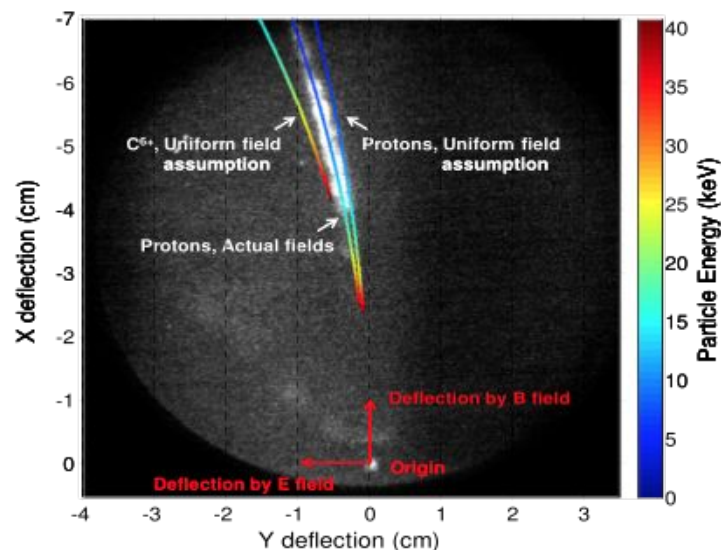


An Interesting (and Unresolved) Characteristic

In the deflagration mode the plasma is a source of high energy protons

- small fraction of the total hydrogen flux
- evidence of this also reported by Mather in 1969

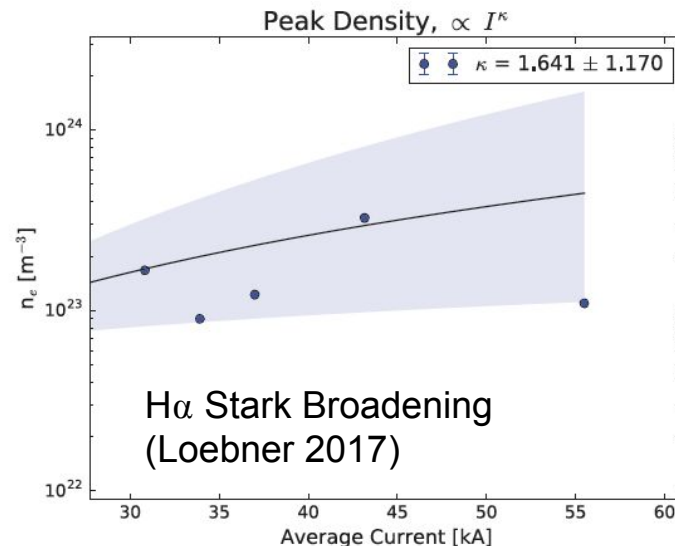
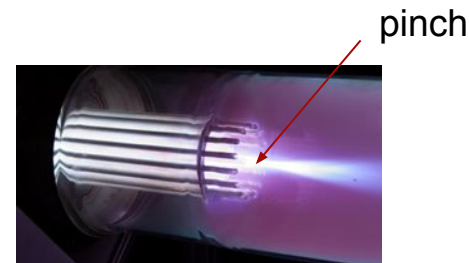
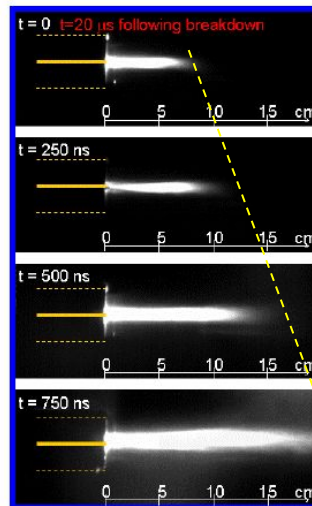
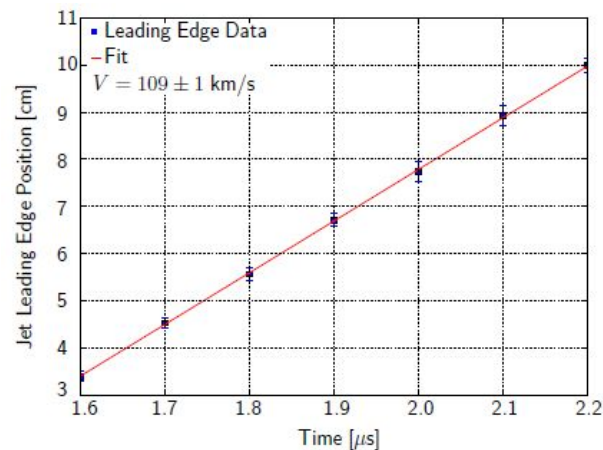
Poehlmann et al, PoP 2013



Interesting Properties at Exit of the Gun

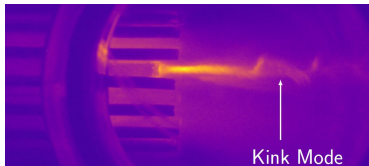
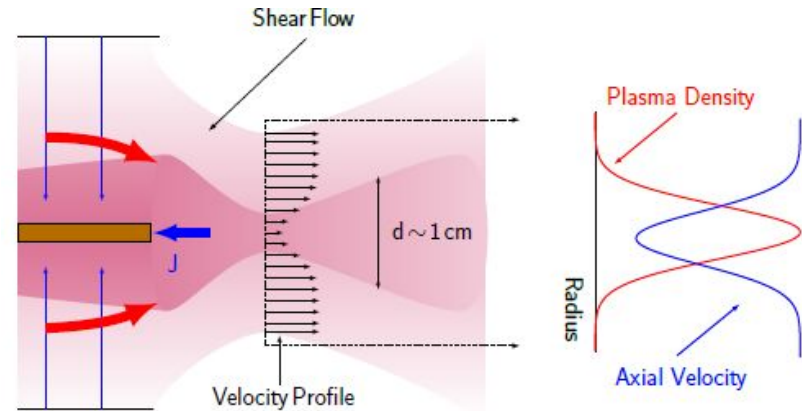
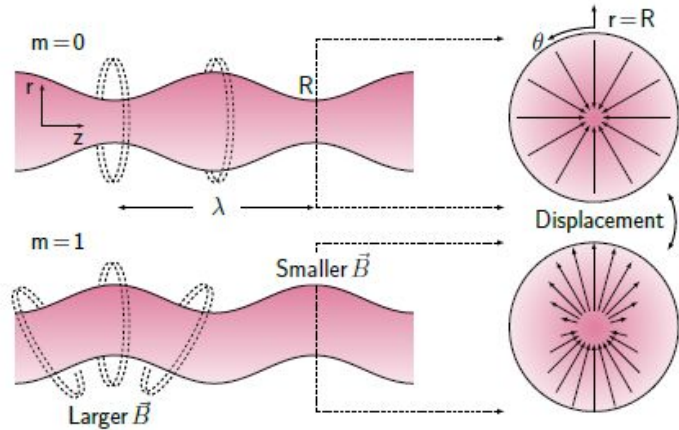
Measured:

- jet velocity, pinch current, n_e
- infer self-generated B (~ 1 T), T (20 – 50 eV)
- flow is supermagnetosonic: $M = \frac{V}{V_A} \approx 1 - 1.5$



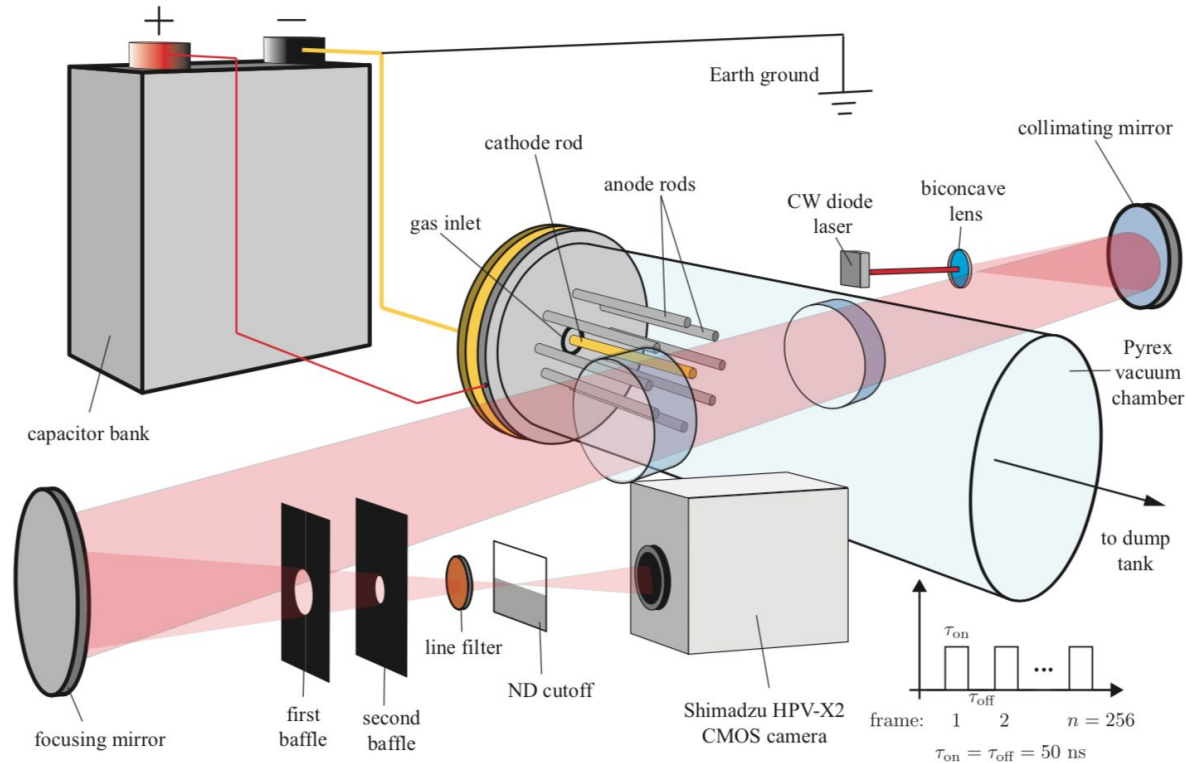
Expect Classic Pinch Instabilities

- plasma at exit is swept into the cathode wake where pinch assemblies
- surrounding flow of higher velocity
- velocity shear serves to stabilize the pinch (Shumlak, PRL 1995)



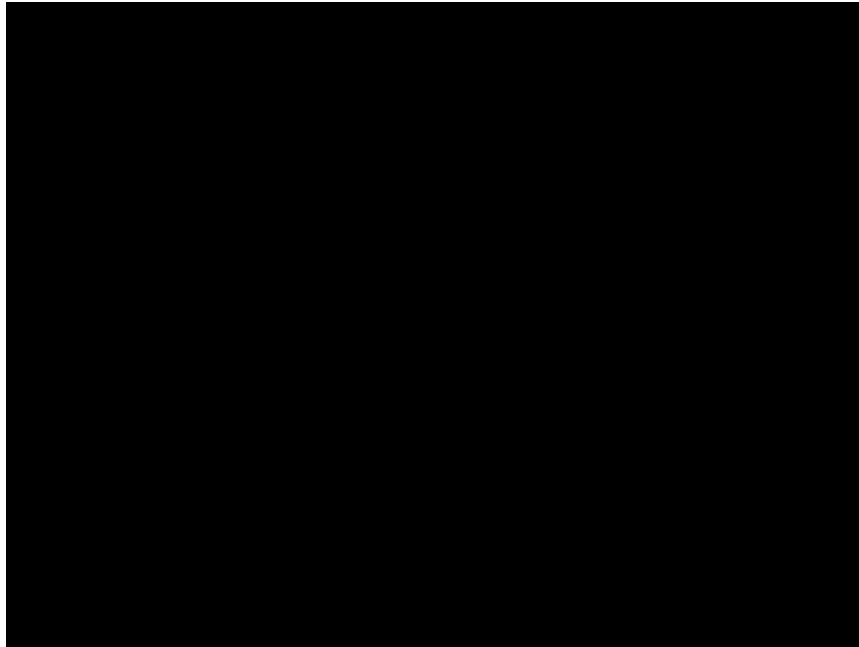
evidence of kink ($m = 1$) instability

Schlieren Videography of Jet and Pinch Region

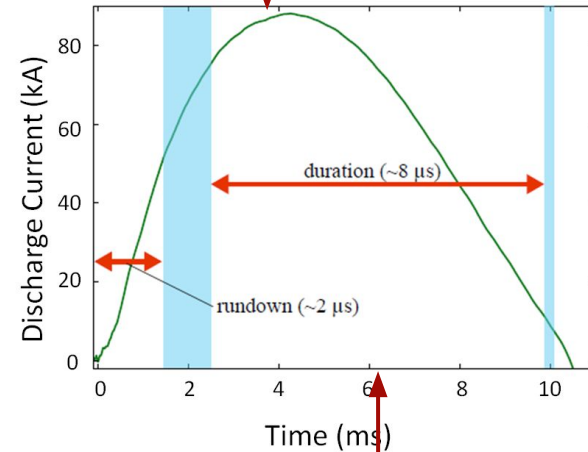
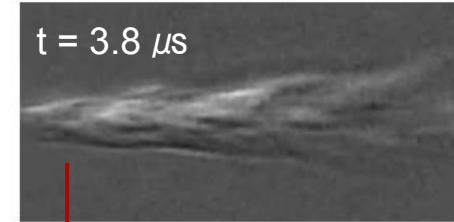


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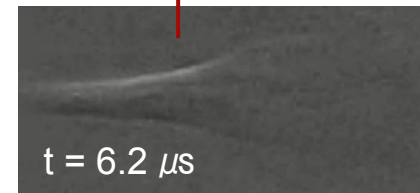
Full 10 MHz Video



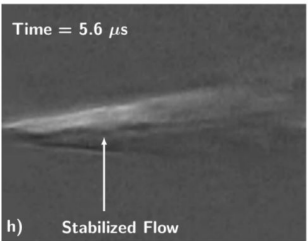
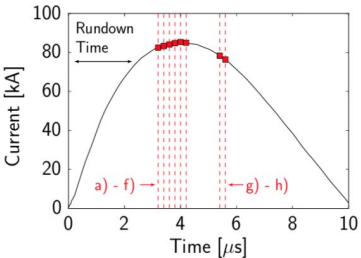
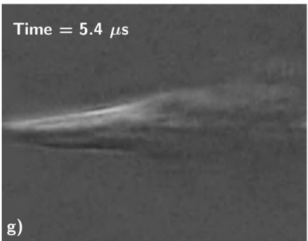
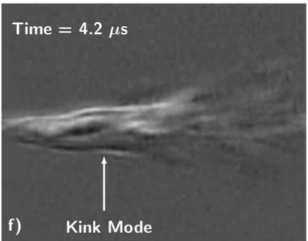
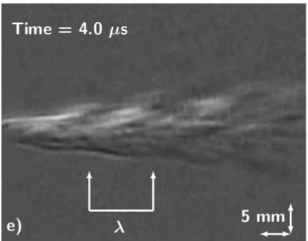
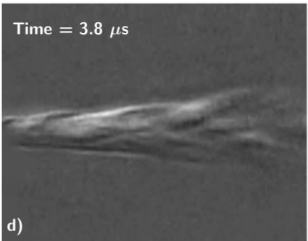
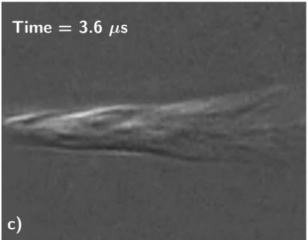
low current/low shear



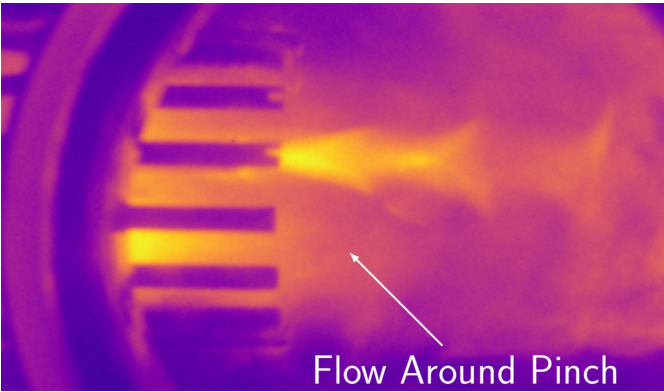
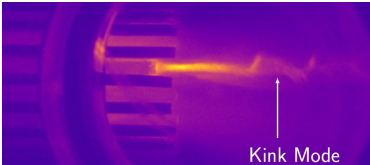
high current/high shear



9 kV (Hydrogen)



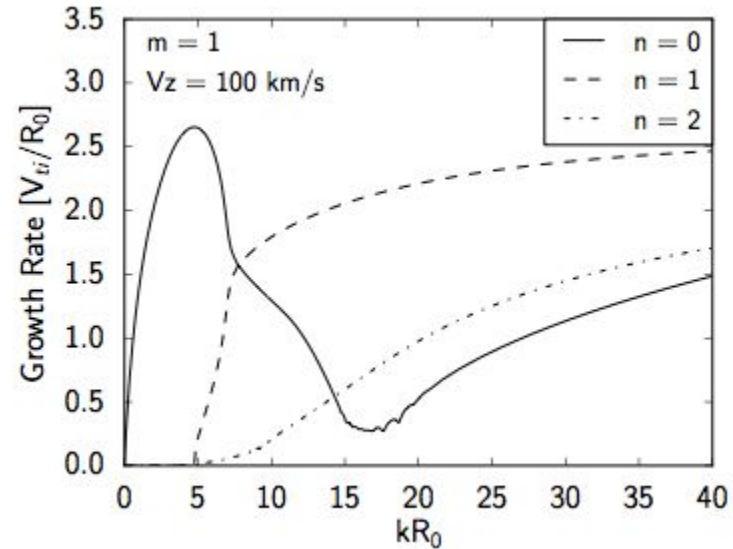
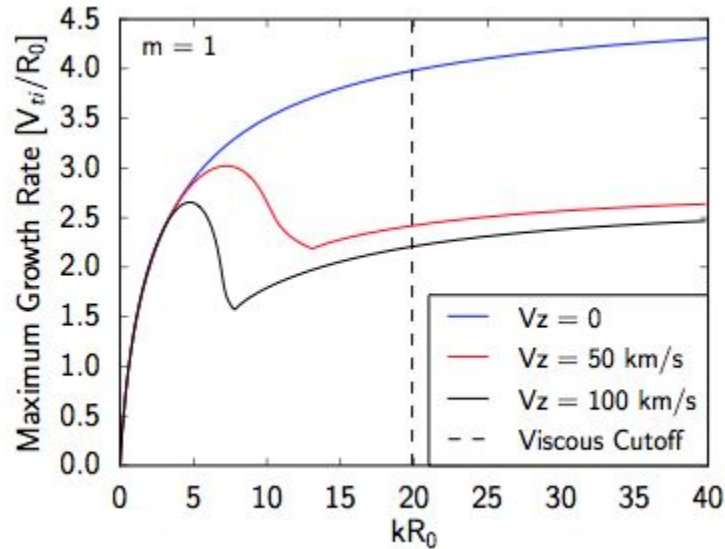
Lower voltages (4 kV): lower external flow velocity and current:



5 ns frames from ICCS camera

Linear Analysis of Jet Stability

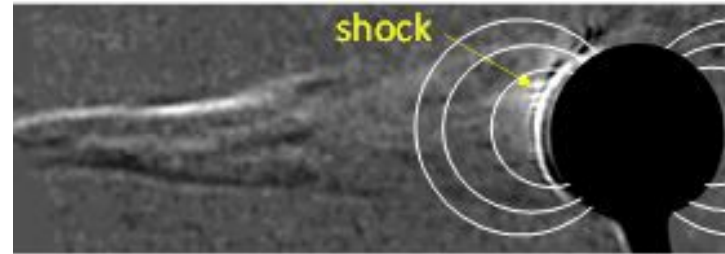
- Fourier analyzed the linearized ideal MHD equations
- Identified the fastest-growing axial mode ($m = 1$, kink)
- Surrounding flow found to damp higher k contributions (shorter wavelength)



Future Work

- What is the source of these fast ions?
- Can we extend the stable period by extending current pulse duration and total energy?
- shocks (for astrophysical relevance)

Interstellar
shock
around
young star
L.L. Ori



Schlieren image of shock around
magnetized ($B = 1\text{T}$) sphere (terrella)

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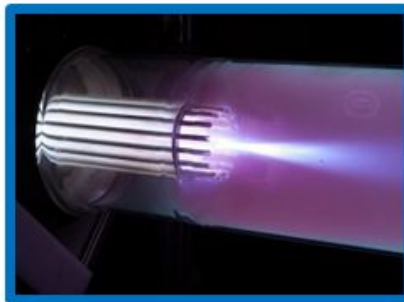
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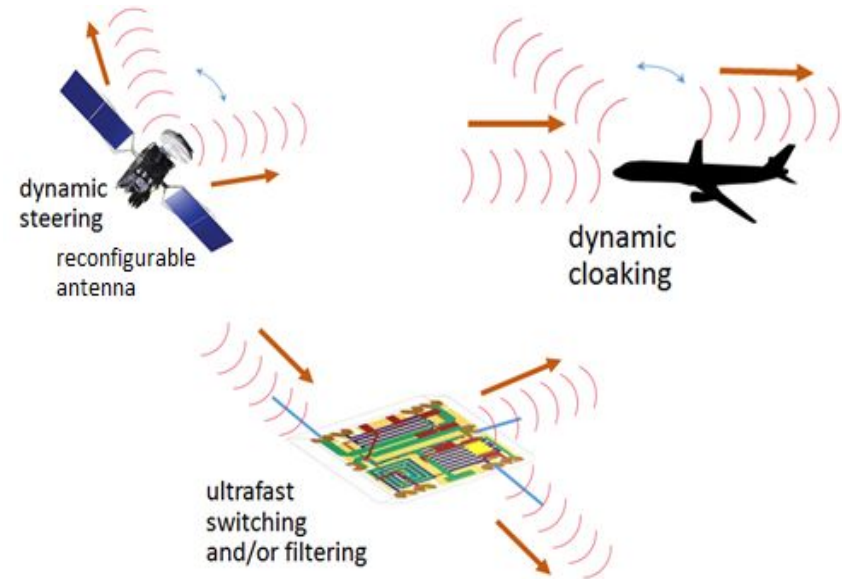
Manipulation and Control of mm-Waves with Plasmas

Currently a field of much commercial and military interest

- reflect/steer/tune
- switch or filter
- amplify or mix
- modulate...

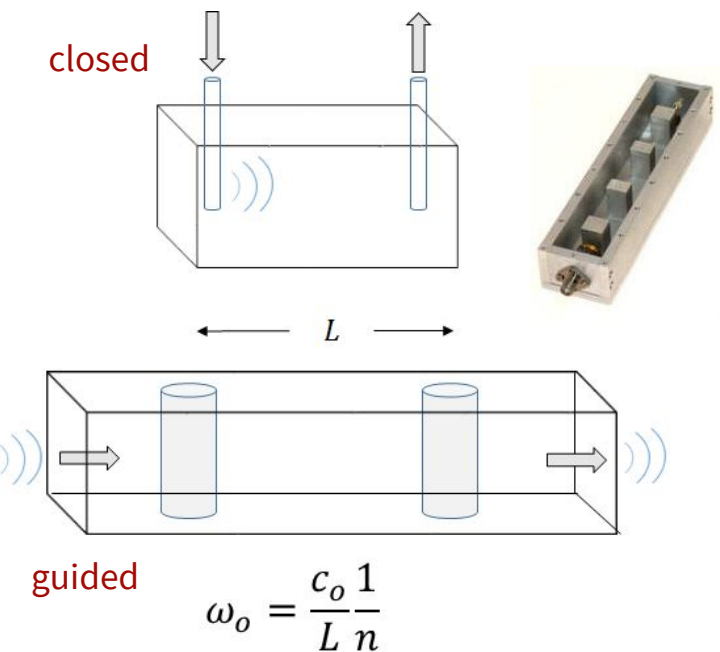
... EM waves with high dynamic fidelity

- not prone to disruption by natural or other EM interferences



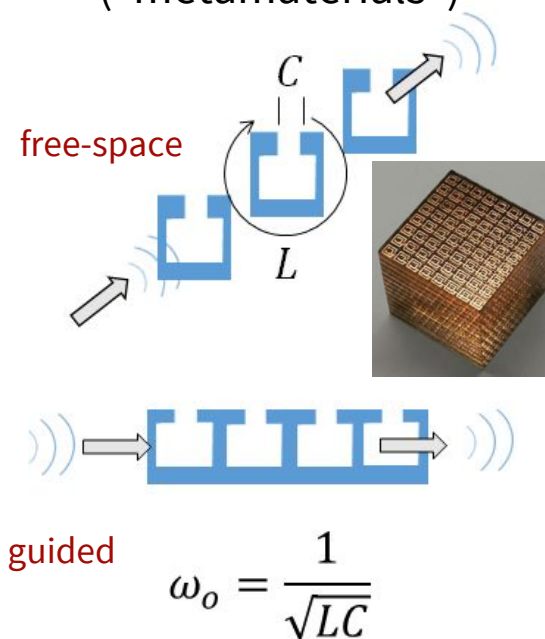
Example EM Platforms

Cavity Resonators

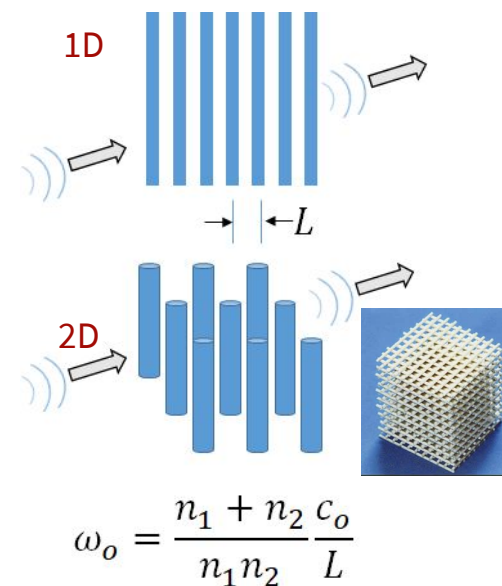


“Artificial” Materials

Structural Resonators (“metamaterials”)



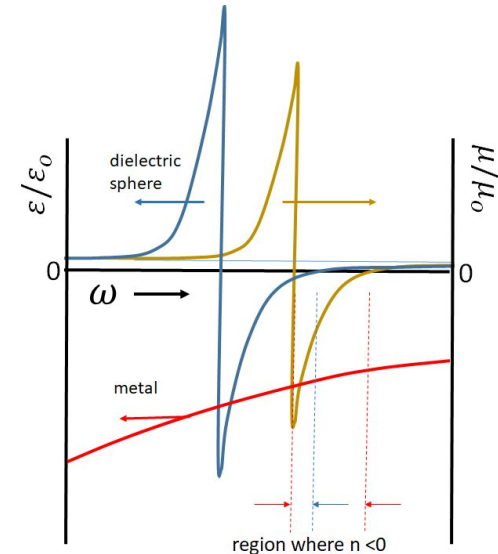
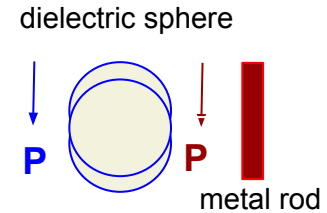
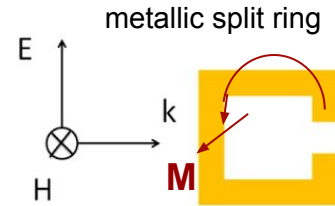
Bragg Resonators Bandgap Materials (“photonic crystals”)



(mid-gap frequency for maximum gap width)

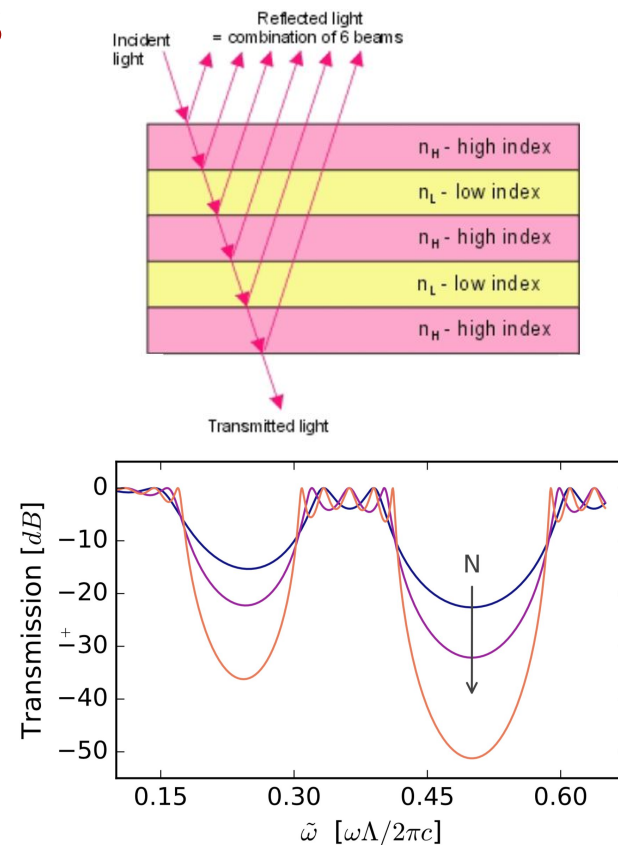
“Artificial” Materials: Structural Resonators

- “elements” (“atoms”) have a geometry that elicits EM response
- array or results in macroscopic transmission properties defining an ϵ , μ
- some configurations can combine elements to engineer unique properties (e.g. negative refraction)
- response is often limited to narrow ranges of ω and not tunable



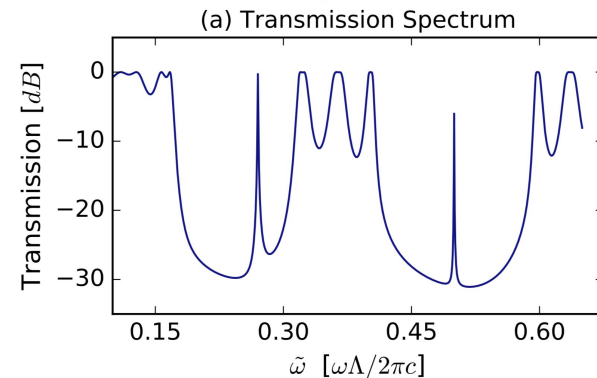
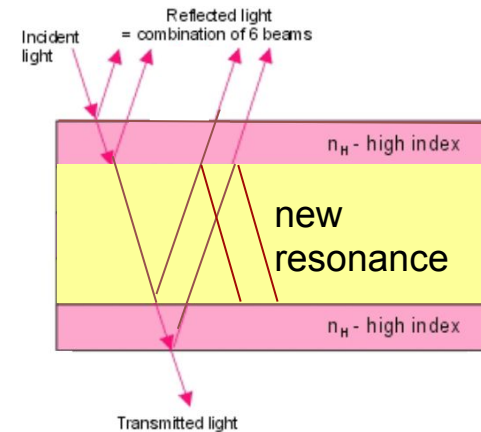
“Artificial” Materials: Bragg Resonators

- “Bragg scattering” - used in understanding atomic structure in materials
- “elements” (“atoms”) have an ordered structure resulting in interface reflection
- constructive/destructive interference produces bandgaps/transmission windows



“Artificial” Materials: Bragg Resonators

- “Bragg scattering” - used in understanding atomic structure in materials
- “elements” (“atoms”) have an ordered structure resulting in interface reflection
- constructive/destructive interference produces bandgaps/transmission windows
- configurations can be engineered to have “defects”
- generally not tunable/reconfigurable



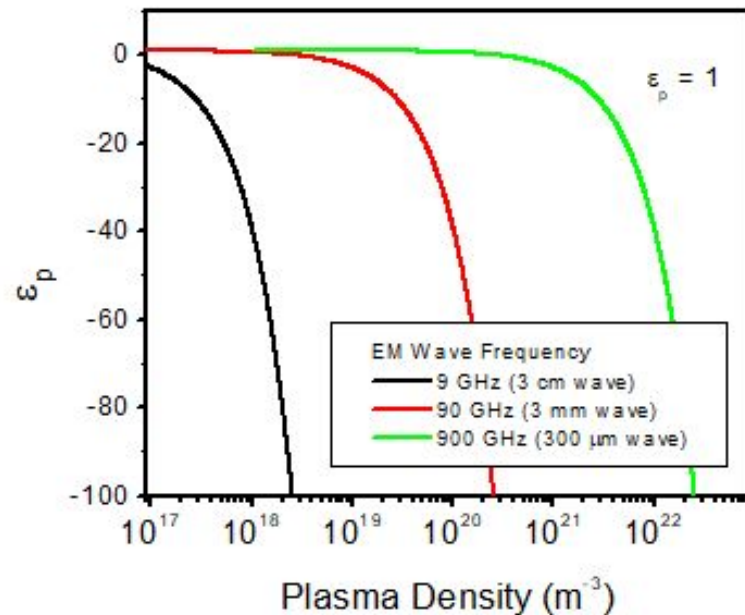
Can a plasma be used to dynamically reconfigure metamaterials or photonic crystals?

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$$\epsilon_p = 1 - \frac{\omega_p^2}{\omega^2}$$

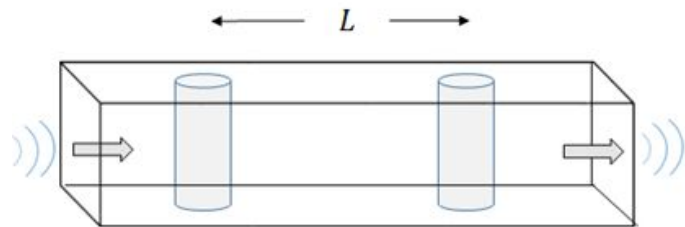
Refractive
index:

$$n = \sqrt{\epsilon_p}$$



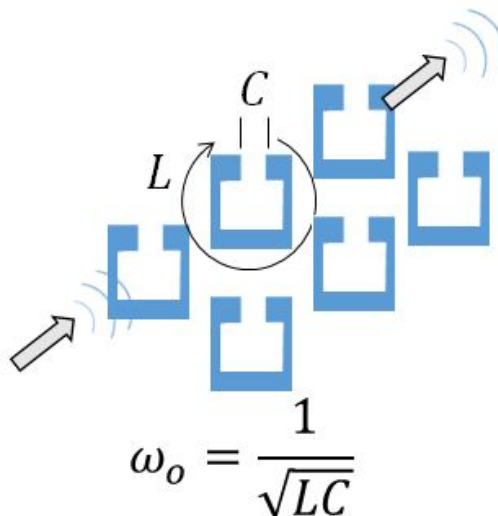
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$$\omega_o = \frac{c_o}{L} \frac{1}{n}$$

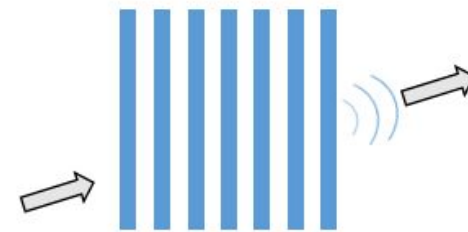
Structural Resonators ("metamaterials")



$$\omega_o = \frac{1}{\sqrt{LC}}$$

"Artificial" Materials

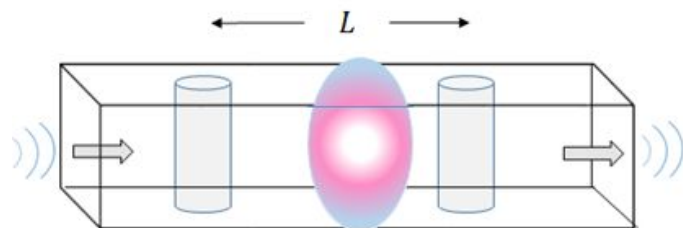
Bragg Resonators ("photonic crystals")



$$\omega_o = \frac{n_1 + n_2}{n_1 n_2} \frac{c_o}{L}$$

Plasma Functionalization

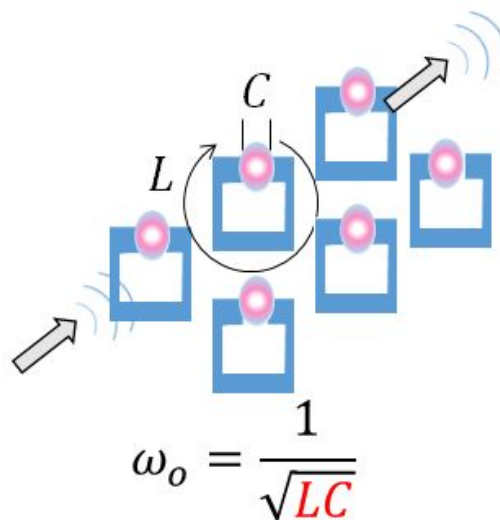
Cavity Resonators



$$\omega_o = \frac{c_o}{L} \frac{1}{n}$$

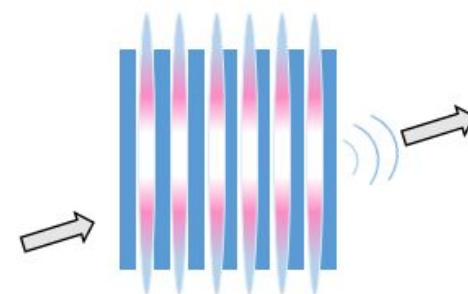
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$$\omega_o = \frac{1}{\sqrt{LC}}$$

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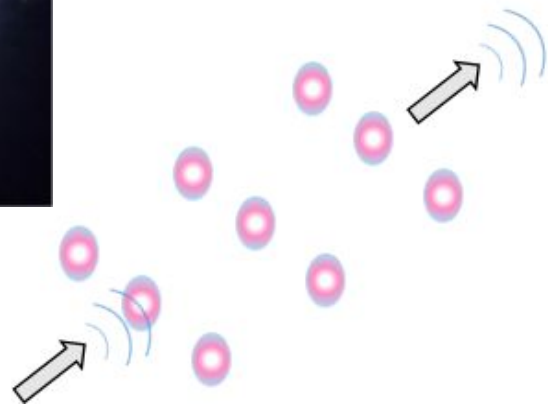
$$\omega_o = \frac{n_1 + n_2}{n_1 n_2} \frac{c_o}{L}$$

Plasmas Themselves can Serve as the Structures and Scatterers

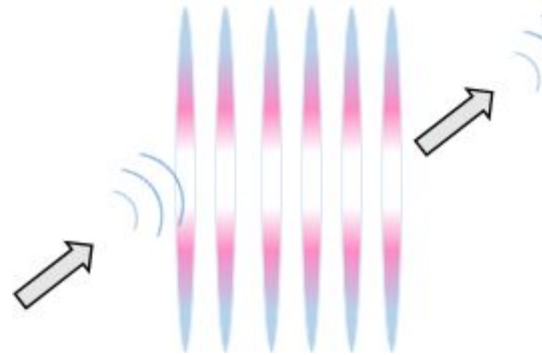
“Artificial” Materials



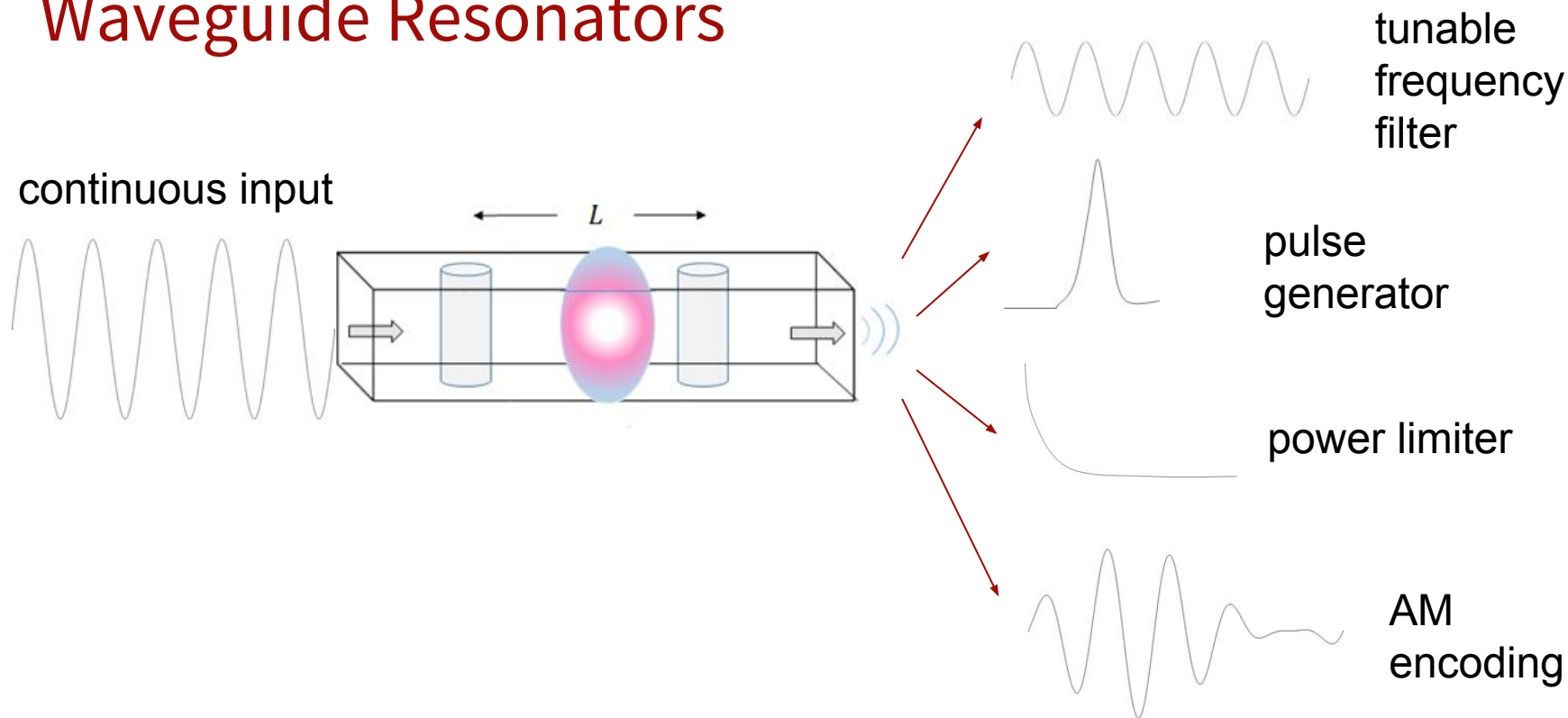
Structural
Resonators
 (“metamaterials”)



Bragg Resonators
 (“photonic crystals”)



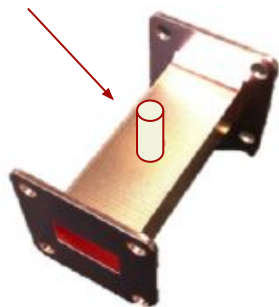
Waveguide Resonators



Plasma-Tuned Waveguide Filter

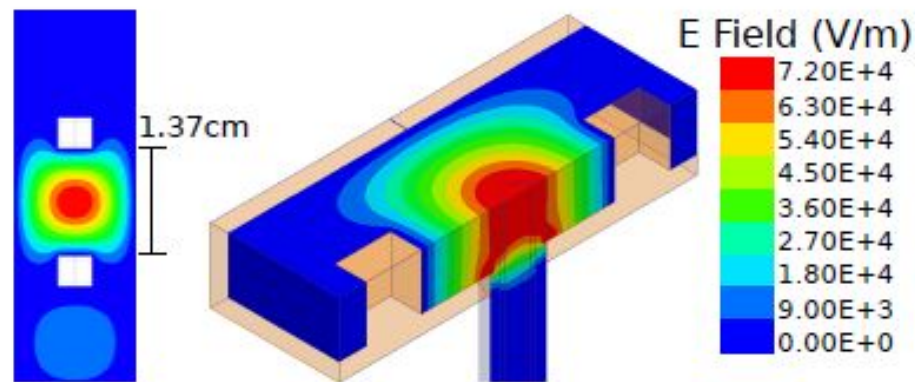
Enables tunable transmission in the range of the empty cavity resonance

5 mm diameter
discharge tube



- Cavity dimensions:
15 x 13.7 x 7.5 mm
- Frequency: 14 GHz
- Quality: 1800

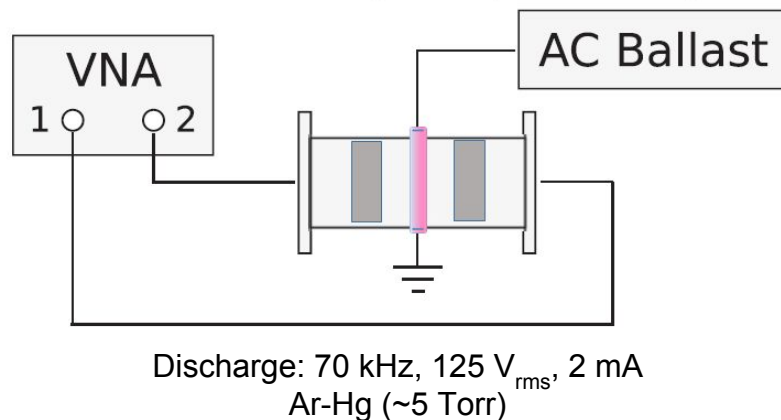
TE_{01} Resonant Mode
 $\omega_o = 13.73$ GHz



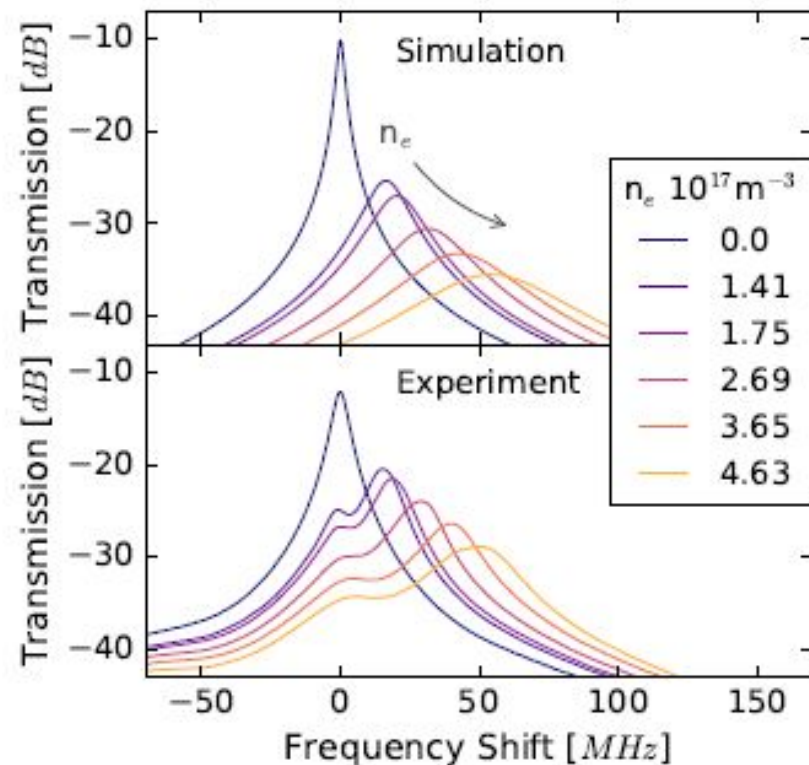
ANSYS (HFSS) simulations of
cavity fields

Tunable Filter

(a) Steady Exp. Setup

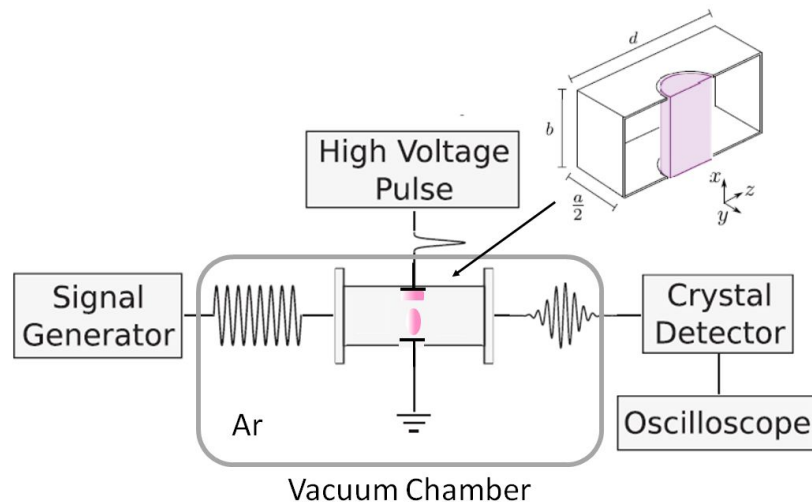
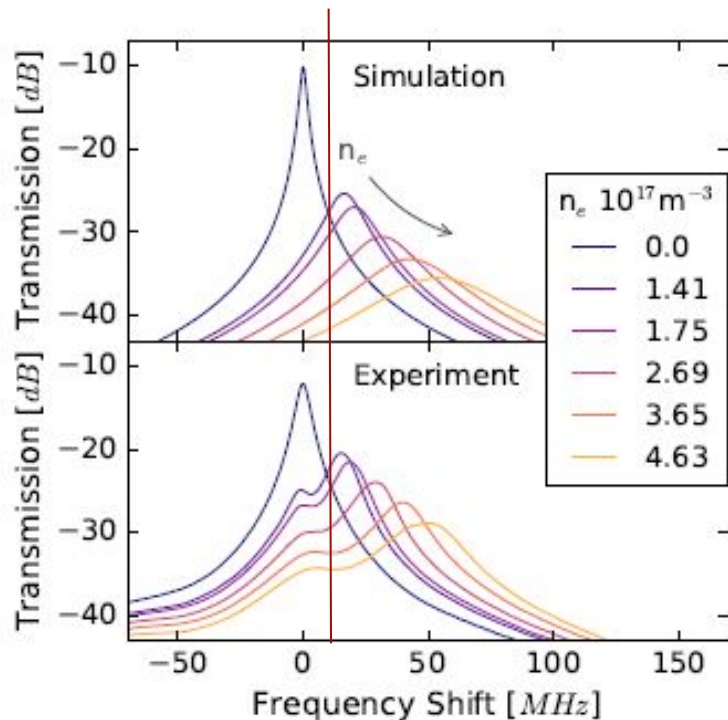


- simulated shifts for plasma column of uniform density
- measure shifts for various discharge currents
 - average plasma density estimated from current and voltage
- good agreement between measured and simulated shifts



Pulse Generation

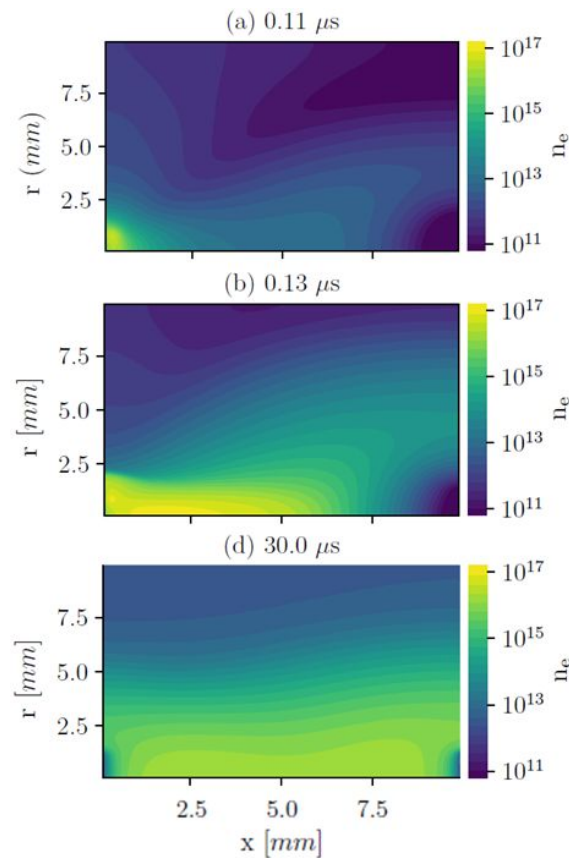
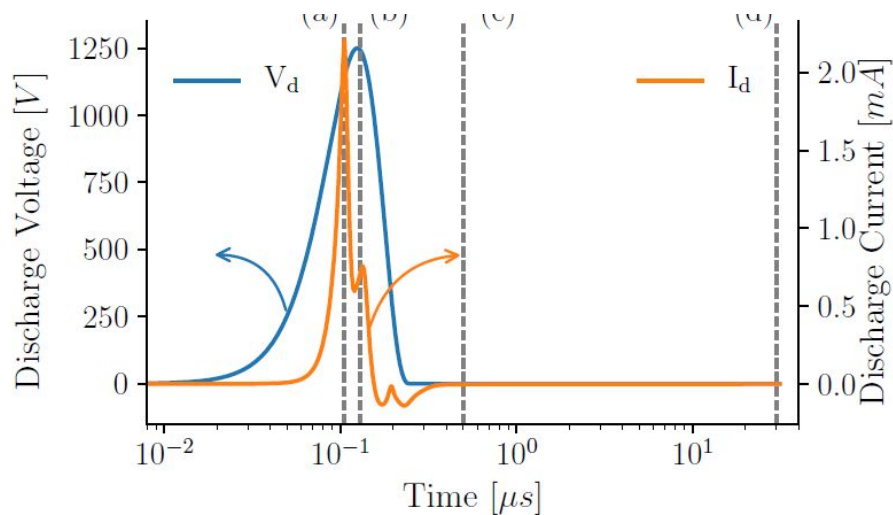
source parked at fixed frequency



- plasma rise time $\sim 0.01 \mu\text{s} <$ cavity response time
- plasma decay time $\sim 1\text{-}10 \mu\text{s}$

Pulse Generation (simulations)

- 2D 2nd order space 4th order time
- SG scheme for fluxes
- transport lookup table from EEDF solver
- finite rate kinetics (including metastable)



Pulse Generation (simulations)

Model the conductor-filled resonator response
(Slater, 1946)

Shift:

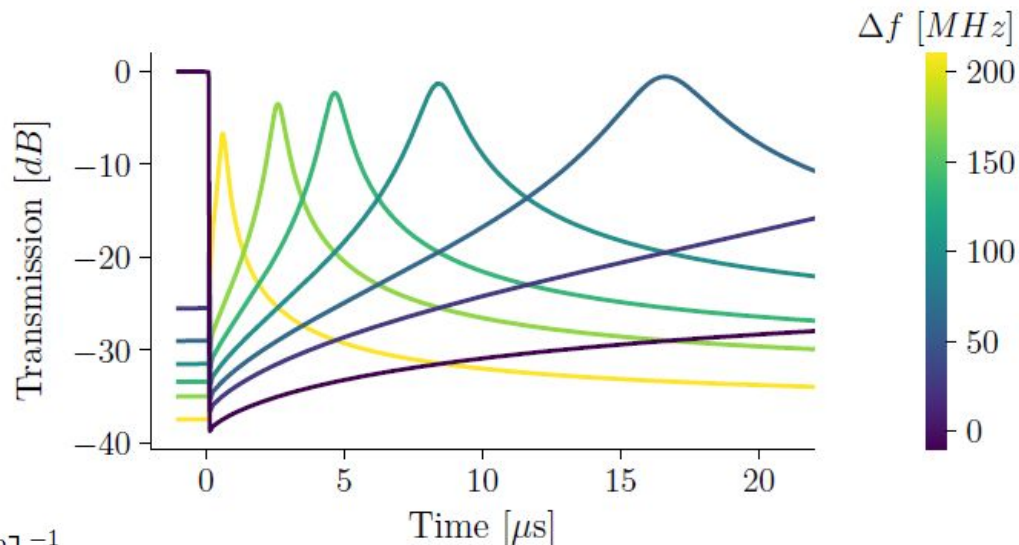
$$\Delta\tilde{\omega}_{\text{res}} = \left(1 - \frac{i+1}{Q_0} - \int_V \frac{\omega_p^2}{\omega_0^2 + i\nu_e\omega_0} E_0^2 dV\right)^{\frac{1}{2}} - 1$$

Quality:

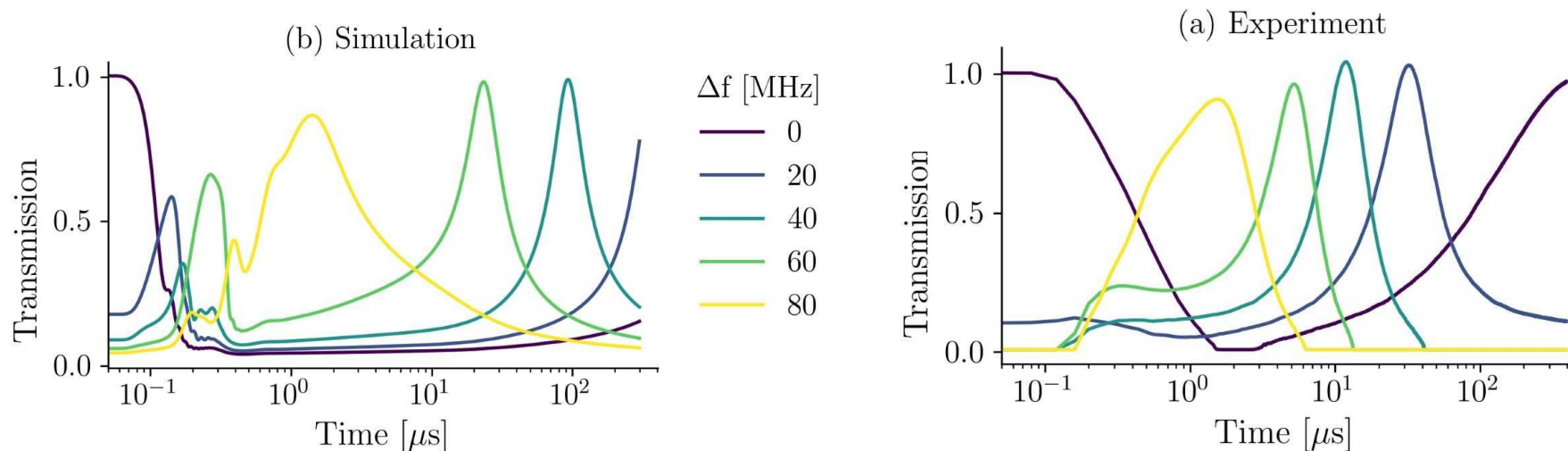
$$Q = -\frac{1}{2} [1 + \text{Re}(\Delta\tilde{\omega}_{\text{res}})] \text{Im}(\Delta\tilde{\omega}_{\text{res}})^{-1}.$$

Transmission:

$$T(\omega) = \left(\frac{1}{Q_0}\right)^2 \left[\left(\frac{\omega \omega_0^{-1}}{1 + \Delta\tilde{\omega}_{\text{res}}} - \frac{1 + \Delta\tilde{\omega}_{\text{res}}}{\omega \omega_0^{-1}} \right)^2 + \left(\frac{1}{Q} \right)^2 \right]^{-1}$$



Comparison to Experiments

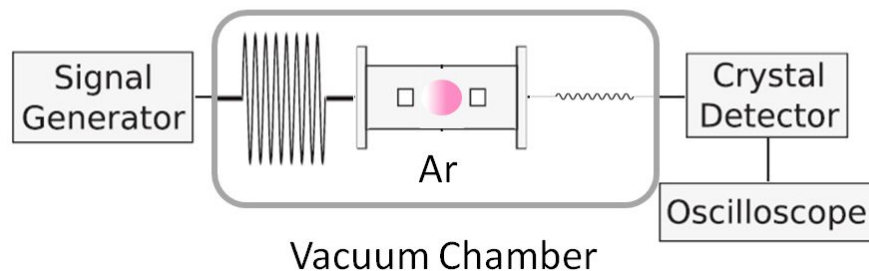


- resolve ionization and recombination signature
- model is underpredicting Ar recombination rates
 - means of measuring recombination rates?

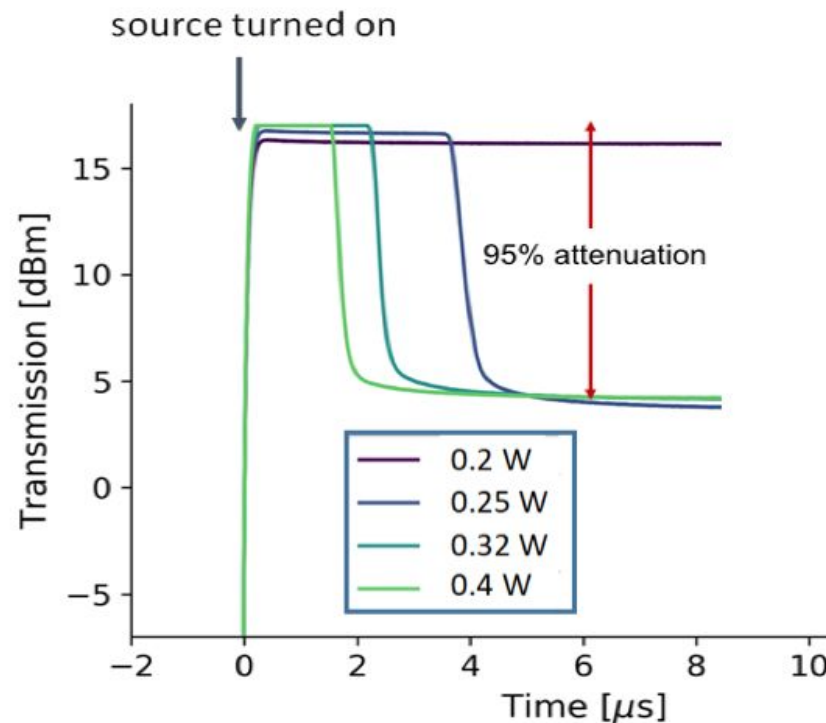
Power Limiter



image: Popular Mechanics 2009

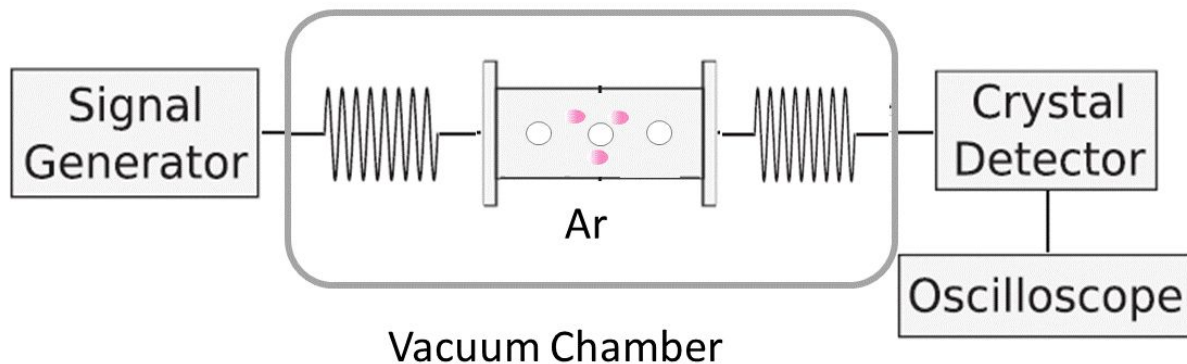


Experimental Results



Amplitude Modulator

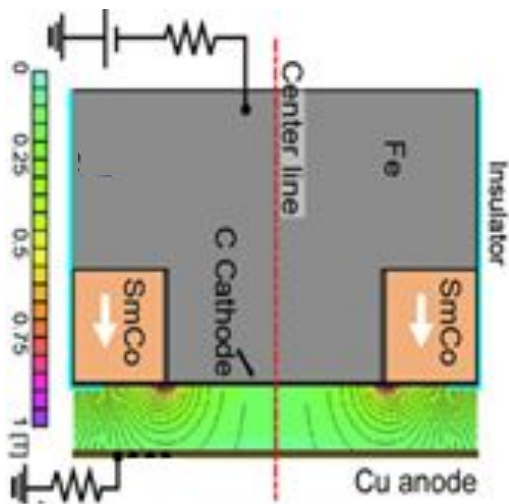
Location of plasma within resonator posts strongly affects transmission



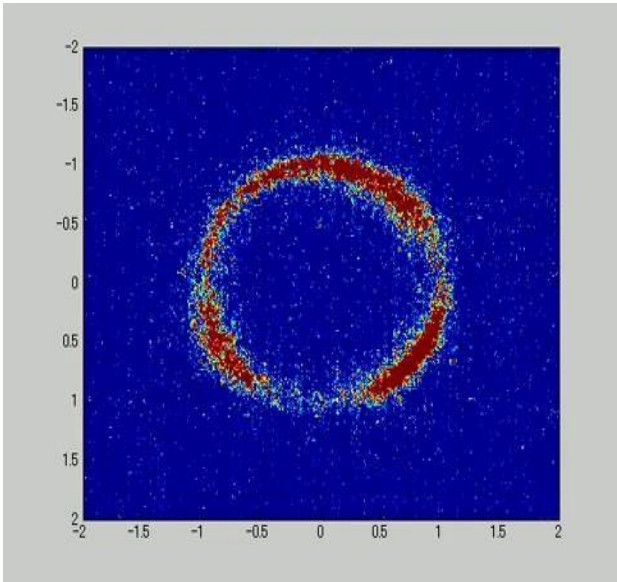
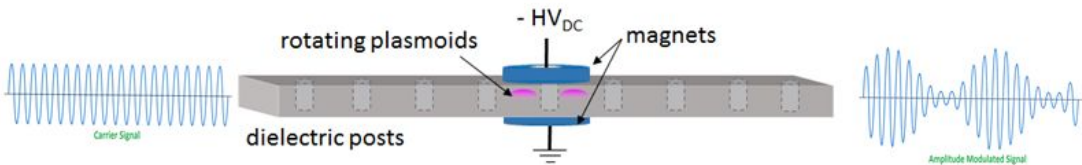
Can we generate coherent structures that rotate at high frequencies (MHz)?

(and fit within a Ku-Band waveguide - 12-18 GHz)

Amplitude Modulator - Micro Magnetron Discharge

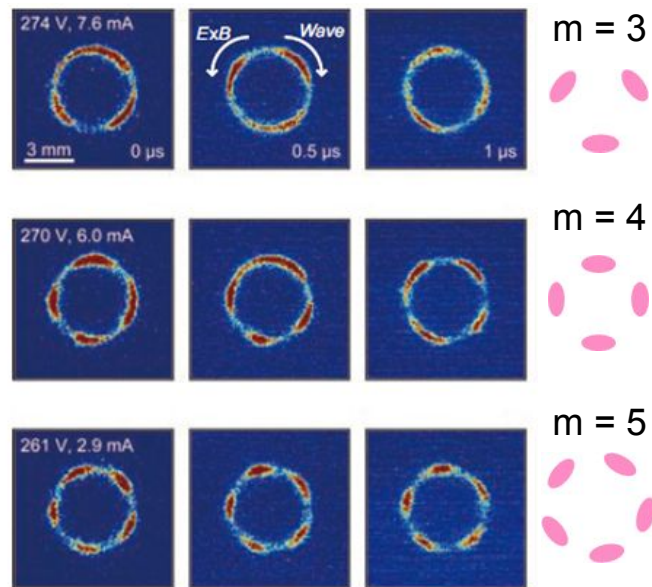


→ 5 mm ←

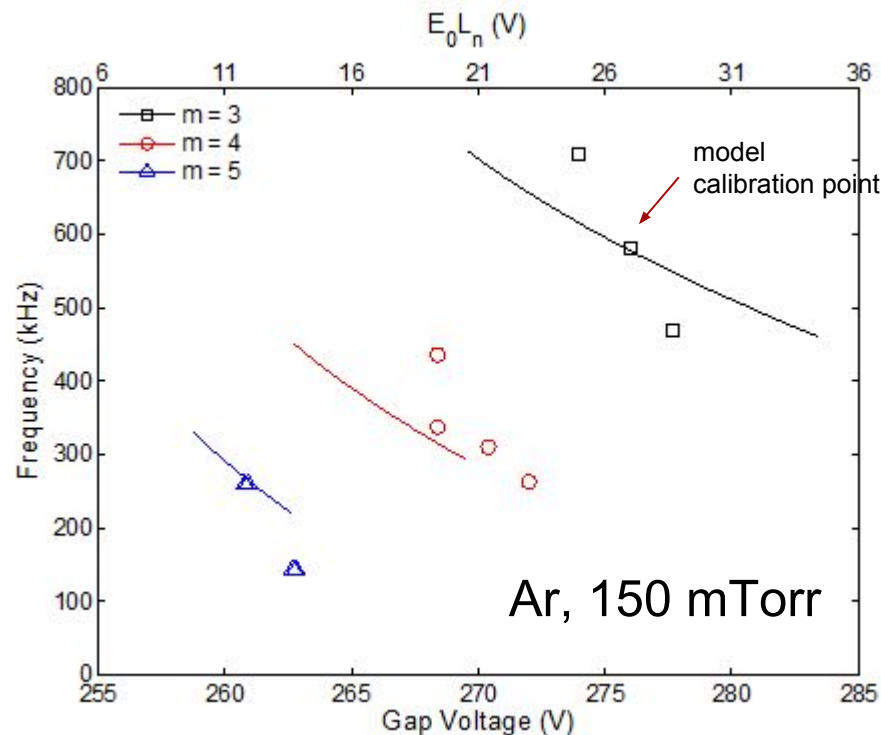


Frequency and Mode Controlled by Voltage

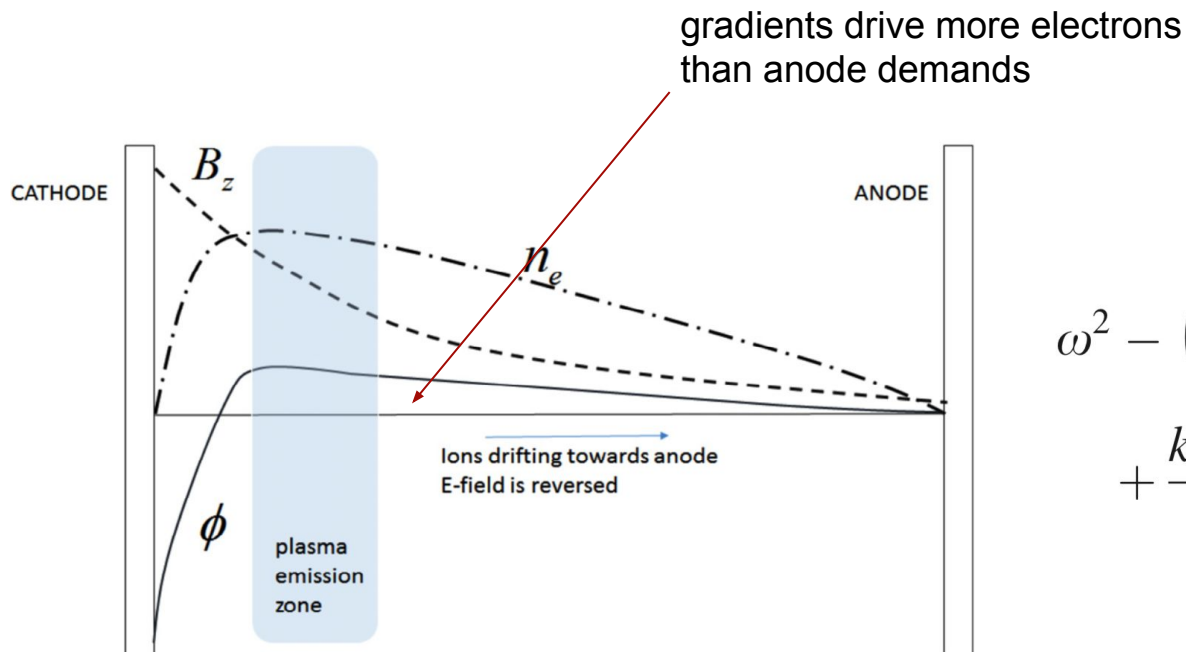
frames from high speed video camera



Model: Gradient-driven drift instability driven by the local $E \times B$ electron rotation (Ito et al, 2016)



Local Field Reversal

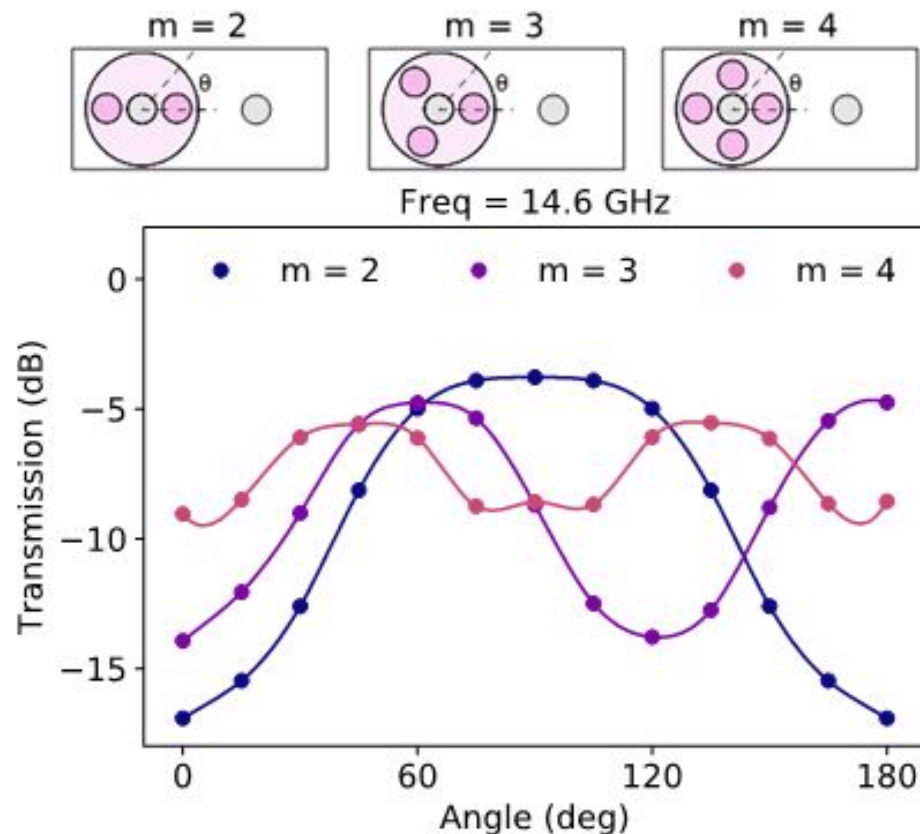


$$\omega^2 - \left(2k_x v_{i0} + \frac{k_{\perp}^2 c_s^2}{\omega^* - \omega_D} \right) \omega + k_x^2 v_{i0}^2 + \frac{k_{\perp}^2 c_s^2 (\omega_0 + \omega_D)}{\omega^* - \omega_D} - i \frac{k_{\perp}^2 c_s^2 \nu_c}{\omega^* - \omega_D} = 0$$

Gradient-drift wave dispersion

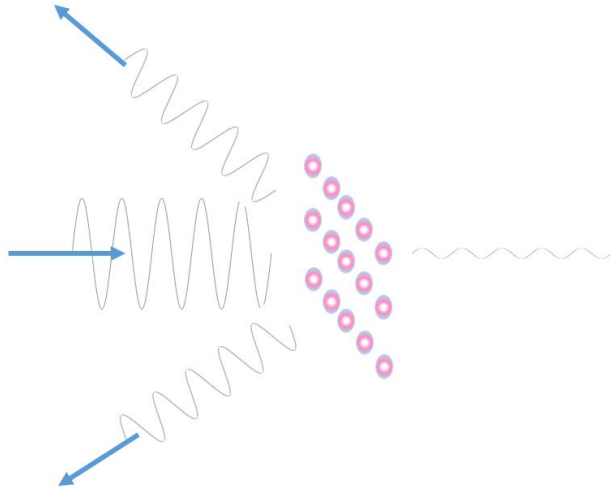
HFSS Simulations

- assumes uniform plasma blobs of circular cross-sections
- indicates a favoring of the $m = 2$ mode (higher modulation amplitudes)

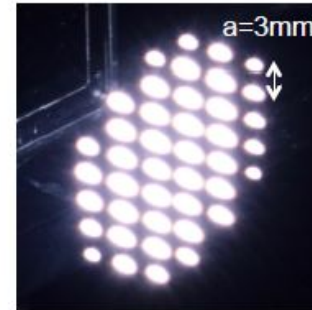
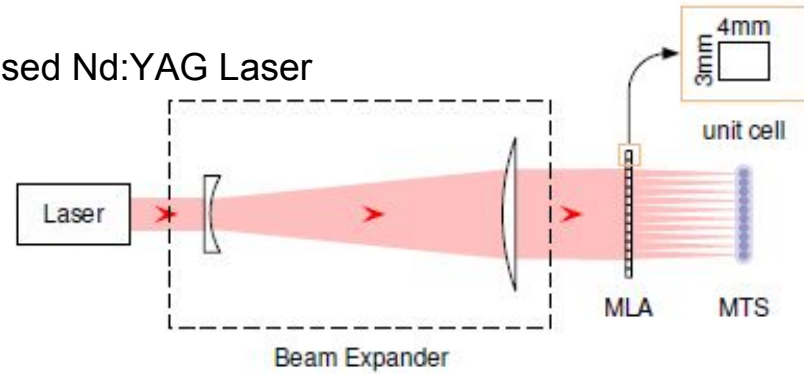


Can we generate a mm-wave mirror in air on demand ?

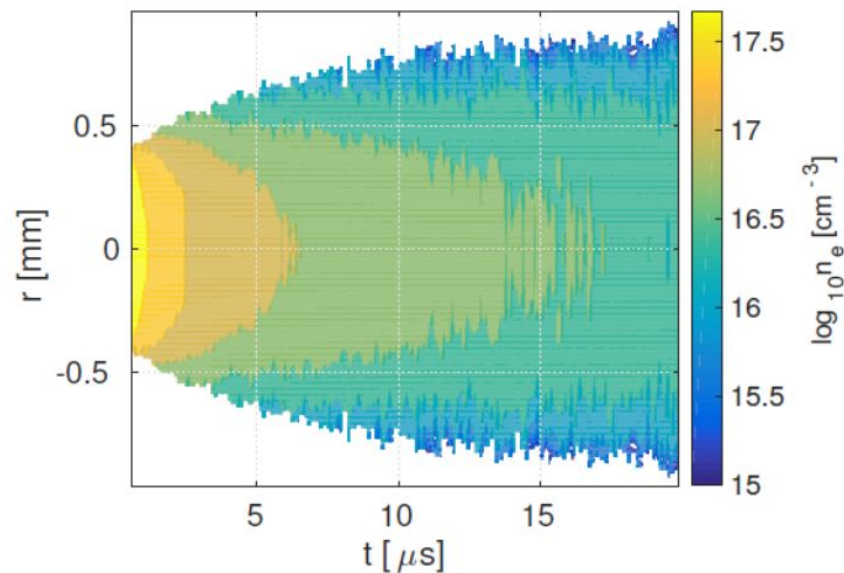
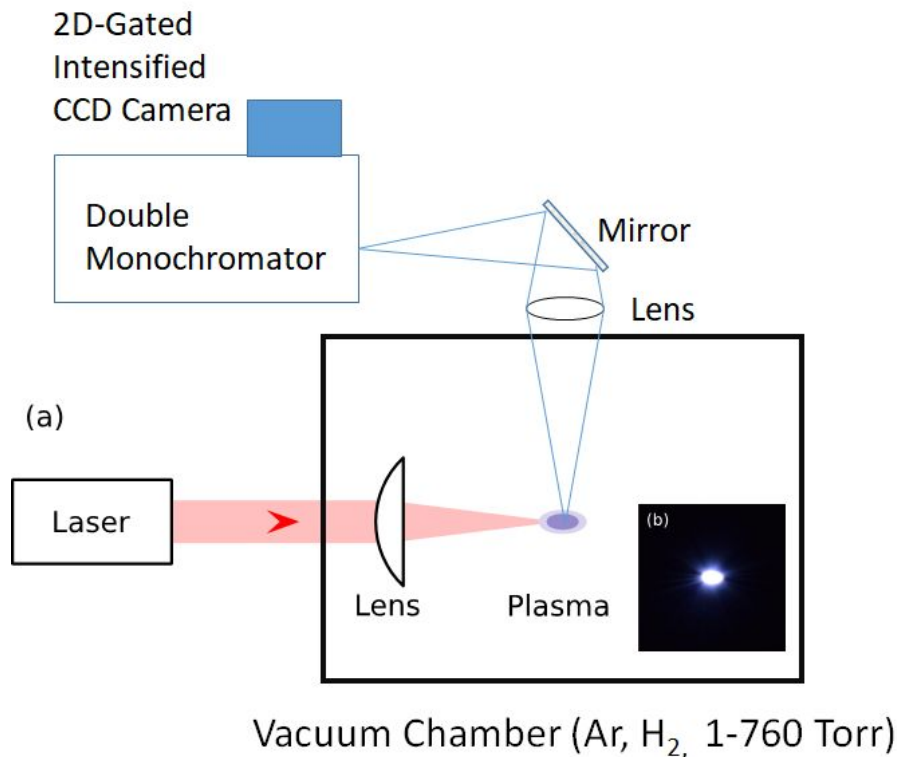
Laser Plasma Metamaterial Reflectors



Pulsed Nd:YAG Laser



Characterizing Single Laser Plasma Kernel



20 Torr Ar (H₂ Seed)
Ne from H α Stark Broadening

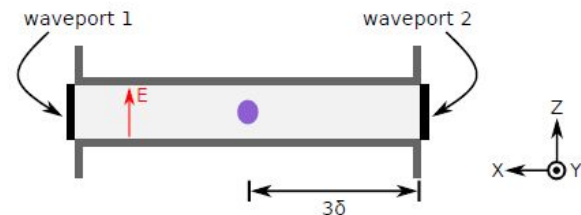
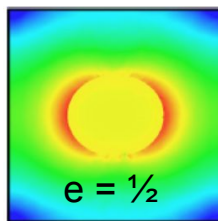
Modeling the Mie Resonance of a Plasma Kernel

Following Mie Theory:
Electric polarization $\mathbf{P} = \alpha \mathbf{E}$

Polarizability:

$$\alpha = v \frac{\epsilon_p - 1}{1 + L_e(\epsilon_p - 1)}$$

Fields in plasma ellipsoid



Eccentricity-dependent parameter:

$$L_e = 1/3 \quad \text{Sphere}$$

$$L_e = 1/2 \quad \text{Cylinder}$$

Absorption and scattering cross sections:

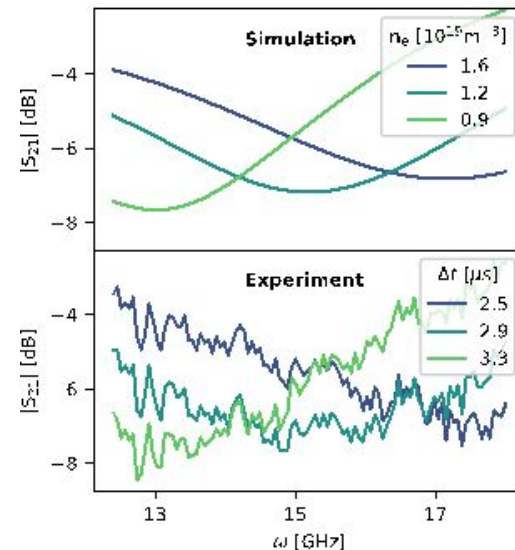
$$\sigma_{abs} = k \operatorname{Im}\{\alpha\}$$

$$\sigma_{scatt} = \frac{k^4}{6\pi} |\alpha|^2 \quad \xrightarrow{\text{collisionless resonance}}$$

$$\omega_o = \sqrt{L_e} \omega_p$$

localized surface
plasmon mode

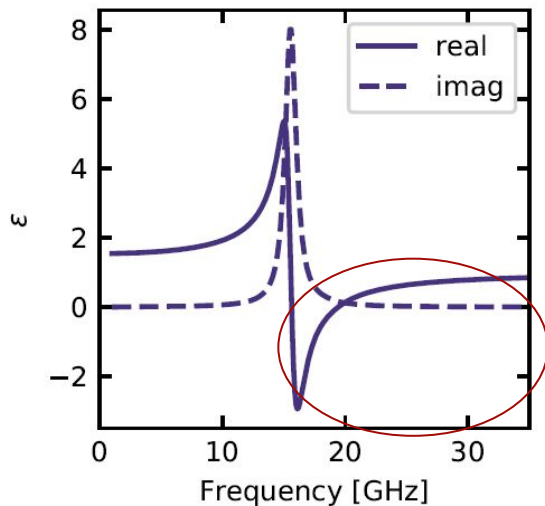
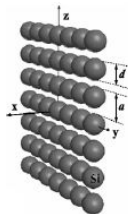
Validating
Experiments for
Single Plasma
Particle in
Waveguide



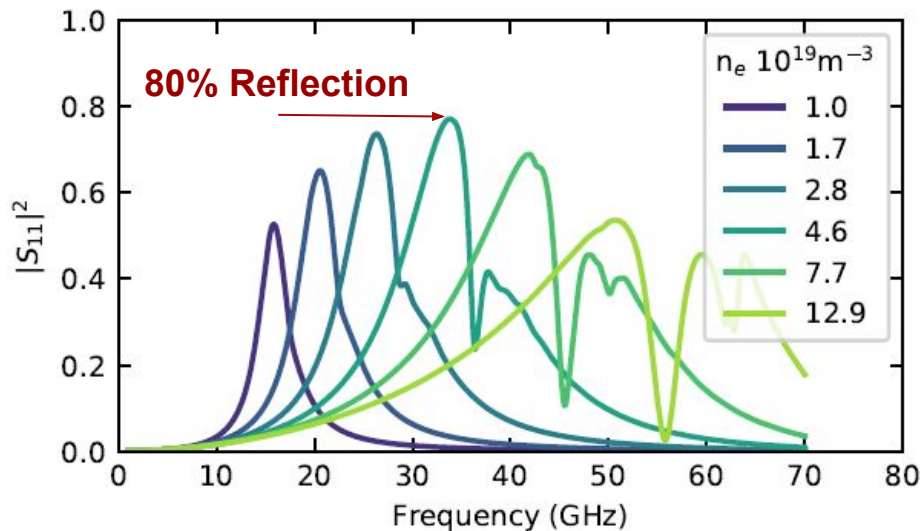
Response of the all-Plasma Metasurface

Model the array as volume-filling Lorentz Oscillators (fill fraction, f)

$$\epsilon = 1 + \frac{f\omega_p^2}{\omega_0^2 - \omega^2 - i\gamma\omega}$$

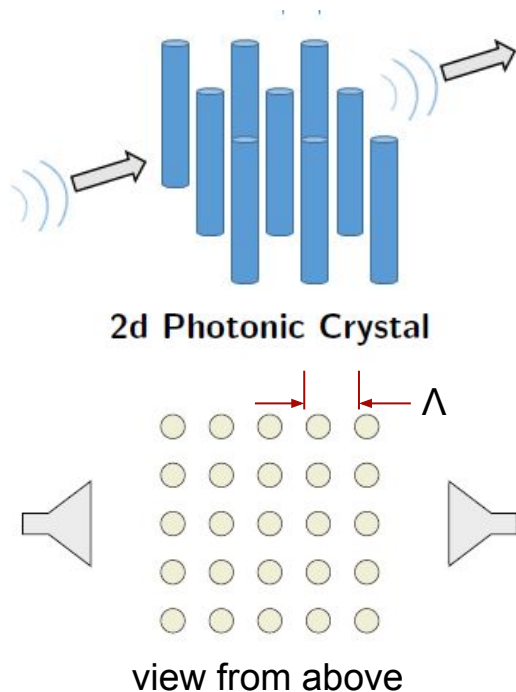


ANSYS simulations of reflection from an infinite array of $a = 3$ mm plasma ellipsoids

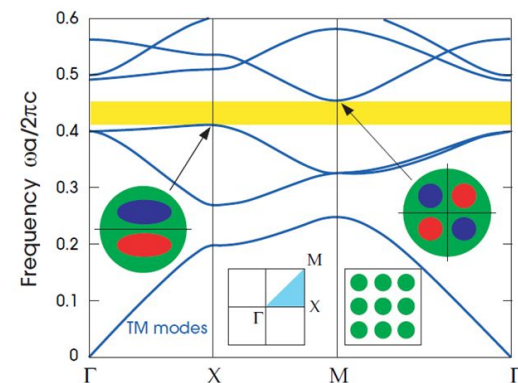
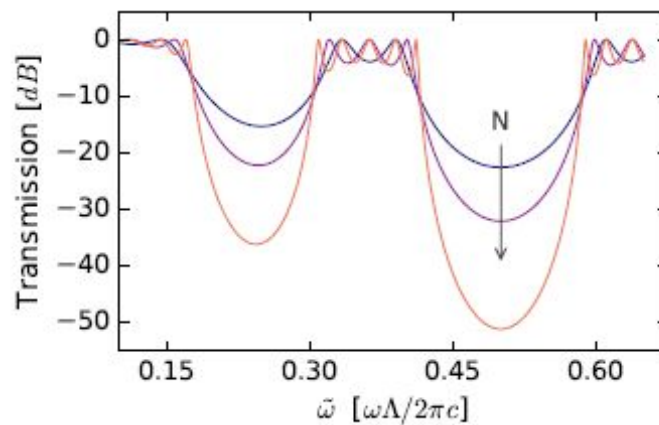


Scattering experiments from the arrays are currently underway

Bragg Resonators: 2D Photonic Crystals



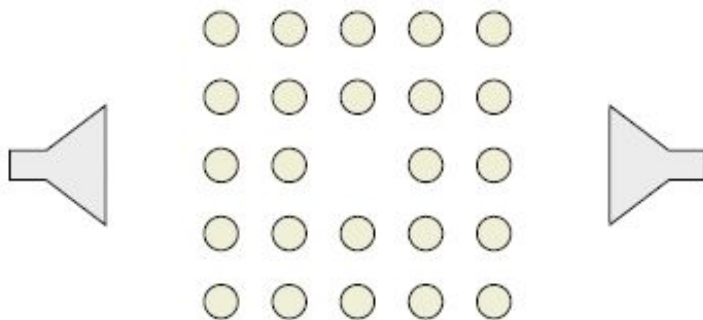
Periodic Band-Gaps



- band gaps are a result of an imaginary wavenumber (dispersion)
- strength of bandgap depends on ratio of dielectric constants (rod and matrix)

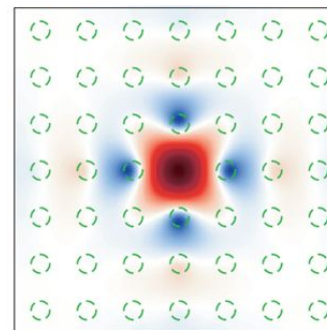
Vacancy Defects

2d Photonic Crystal

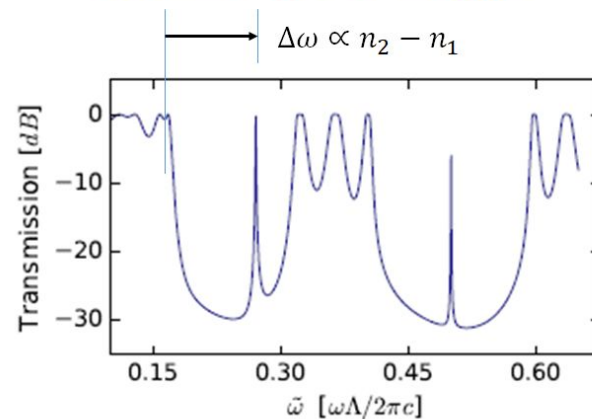


- vacancy defect acts as a resonant cavity
- incident waves at defect resonance “tunnel” in/out through evanescent coupling
- resonant frequency

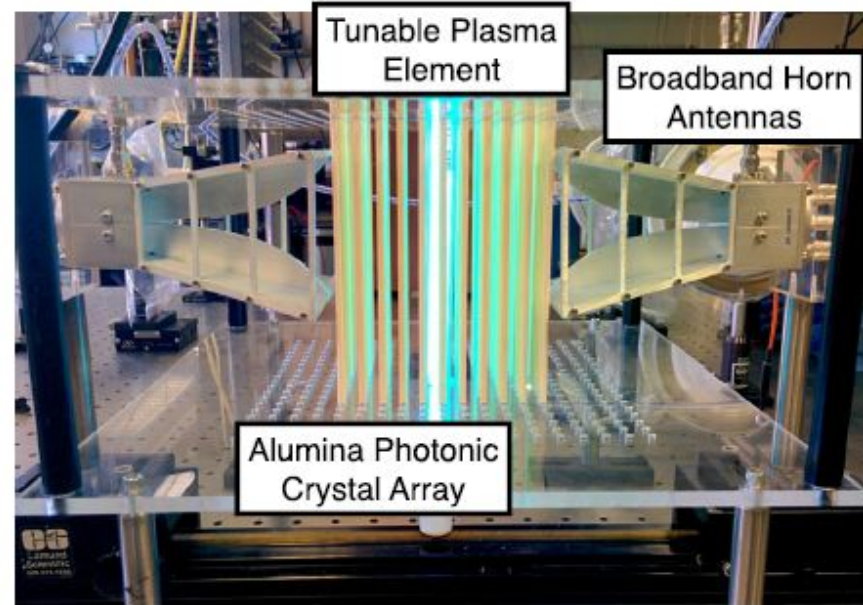
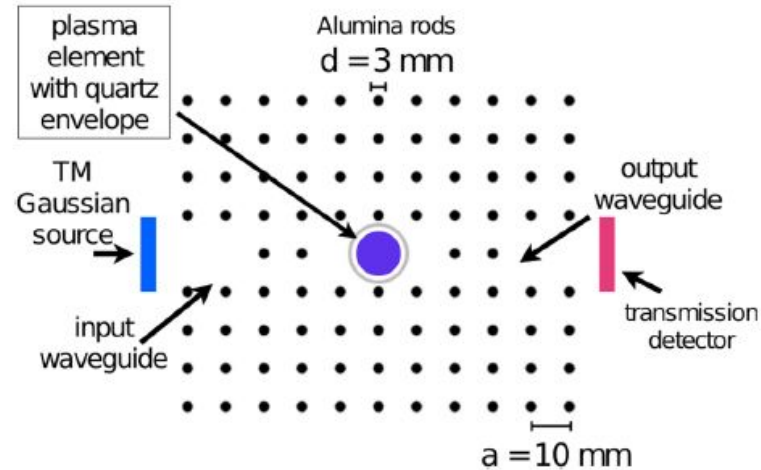
Vacancy Defect Fields



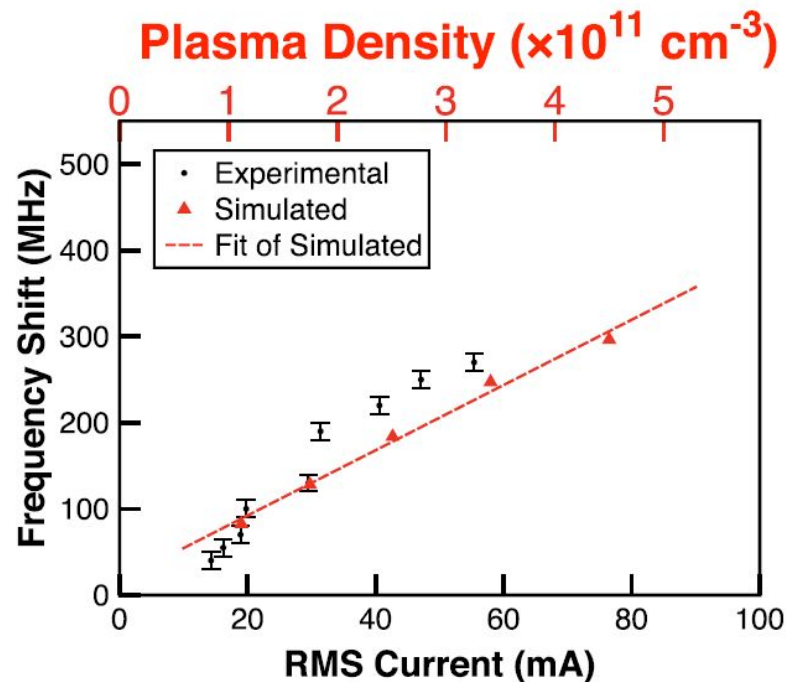
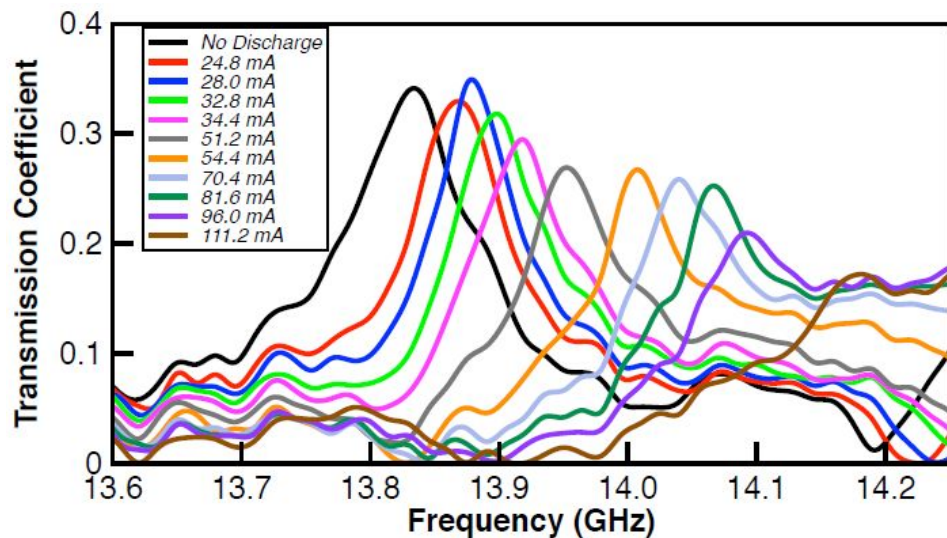
Defect State Resonance



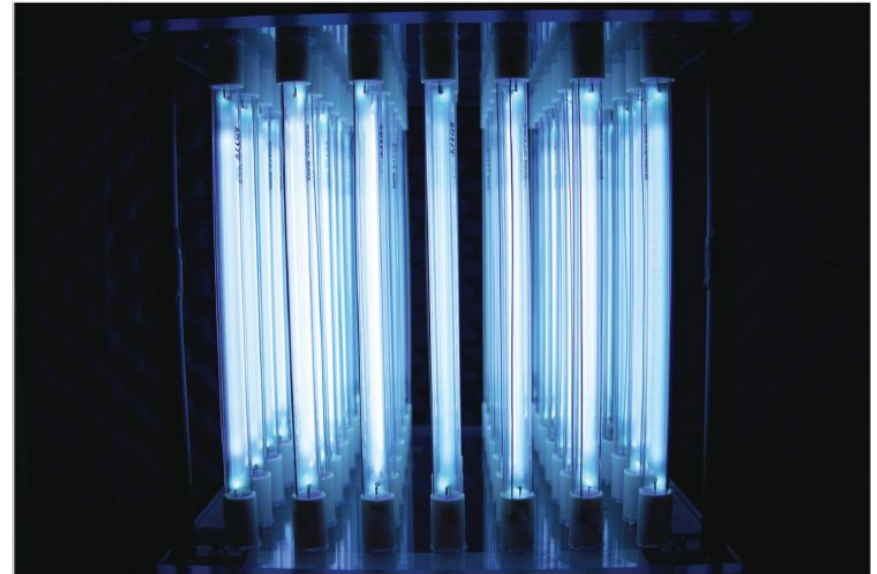
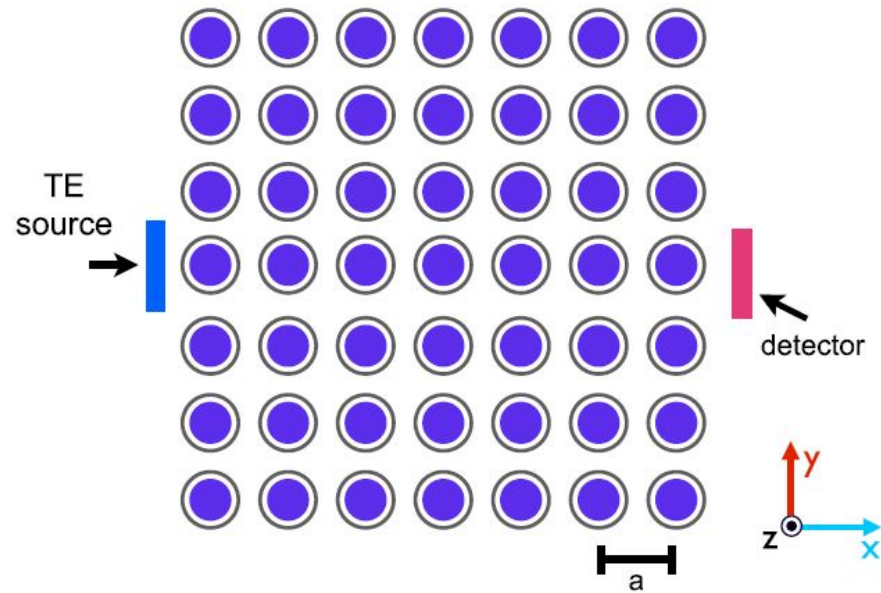
2D Photonic Crystal - Plasma Defect



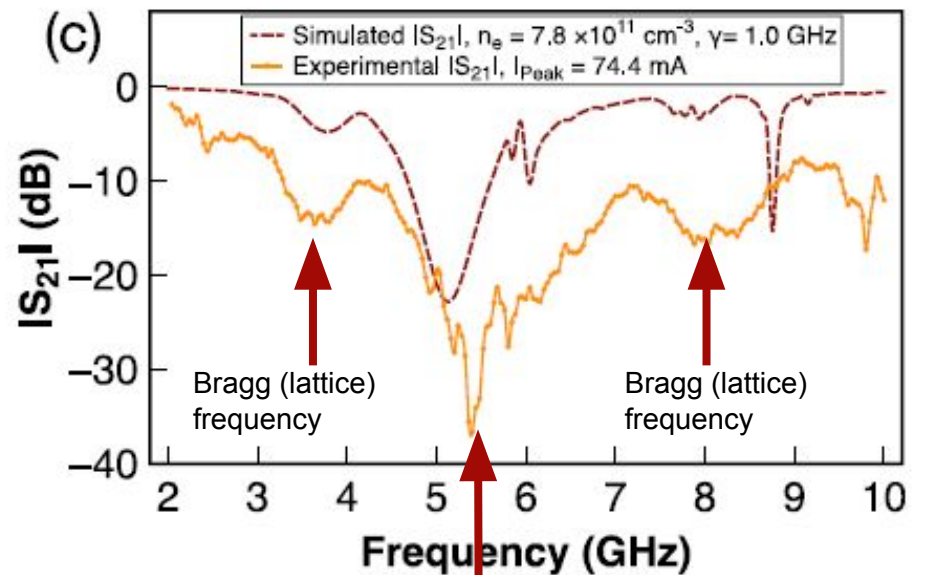
Plasma Tuning of Defect Transmission



All-Plasma Photonic Crystal

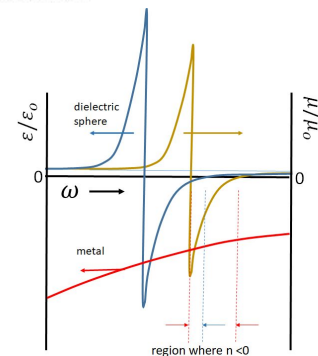
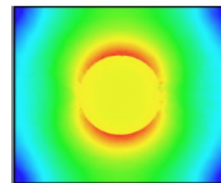


Multiple Gaps - Bragg and Surface Plasmons



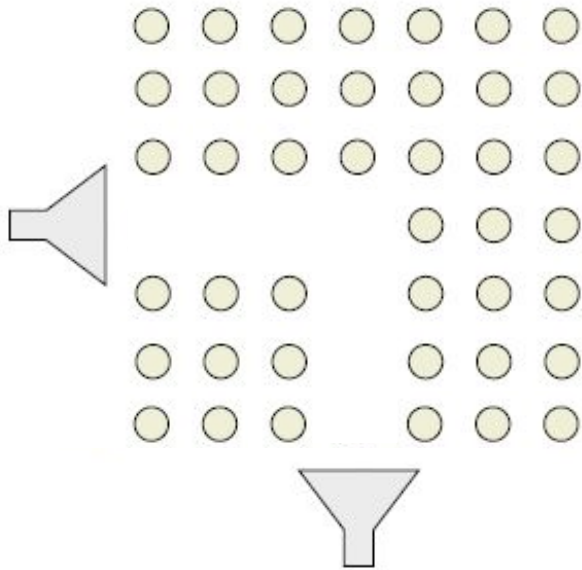
Mie resonance frequency
 (localized surface
 plasmon) $\omega_o = \omega_p / \sqrt{2}$

- Bragg gap too weak to engineer for waveguiding
 - 7×7 crystal size too small
 - refractive index difference (at most 1) is too small
- Surface plasmon gap strong
 - dynamic waveguide?



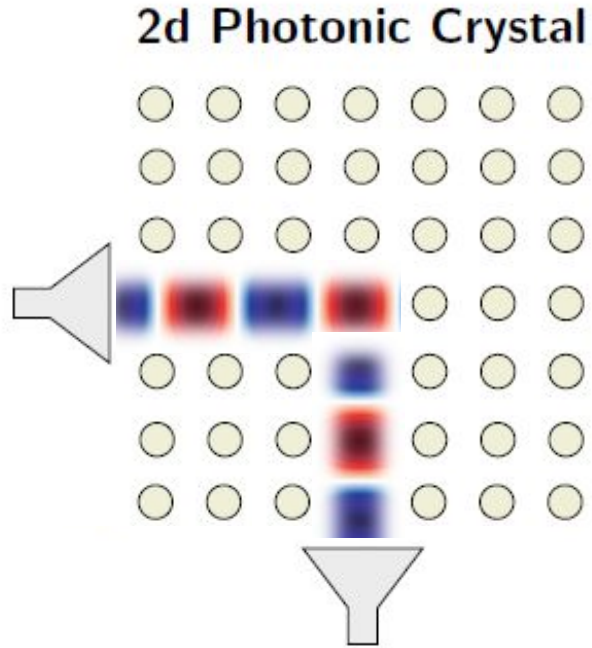
Waveguiding

2d Photonic Crystal



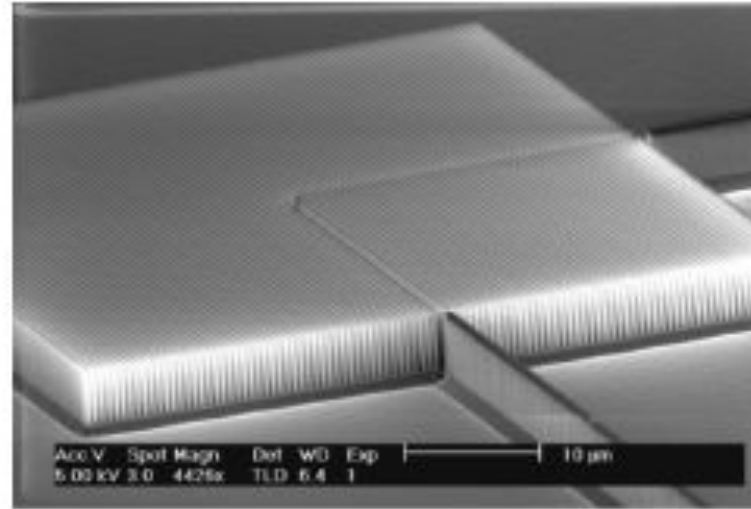
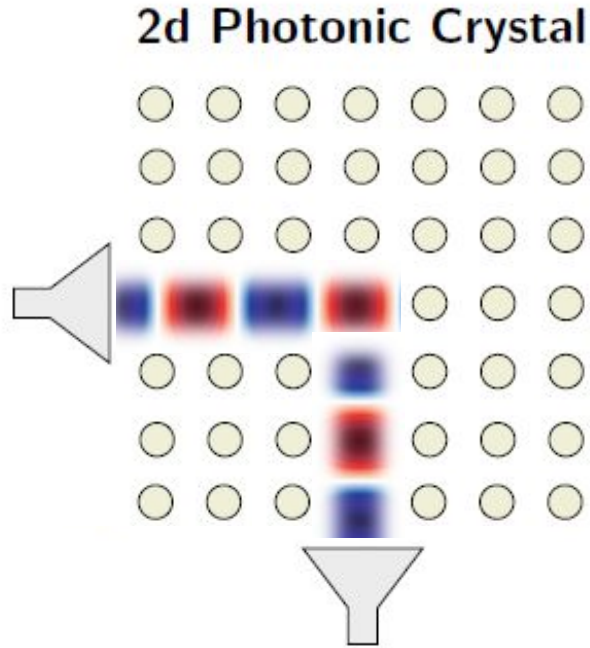
- line vacancies admit waves tuned to the defect state within the bandgap

Waveguiding



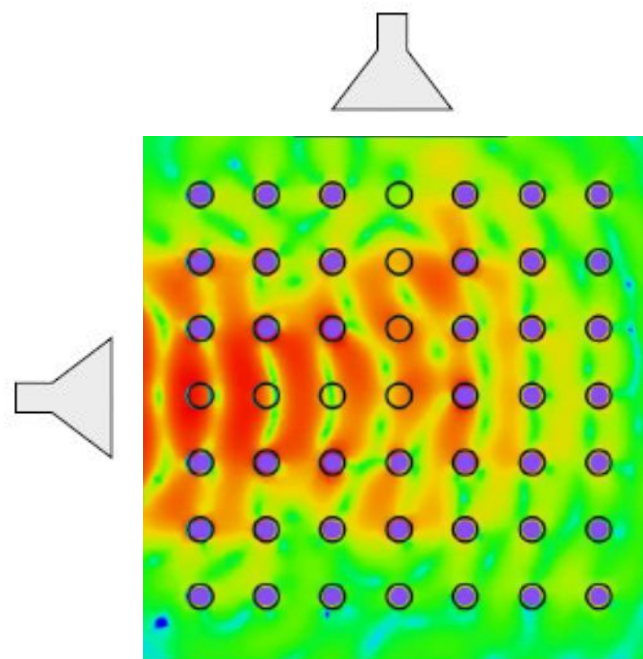
- line vacancies admit waves tuned to within the bandgap
- waves cannot propagate beyond the gap and are guided

Waveguiding

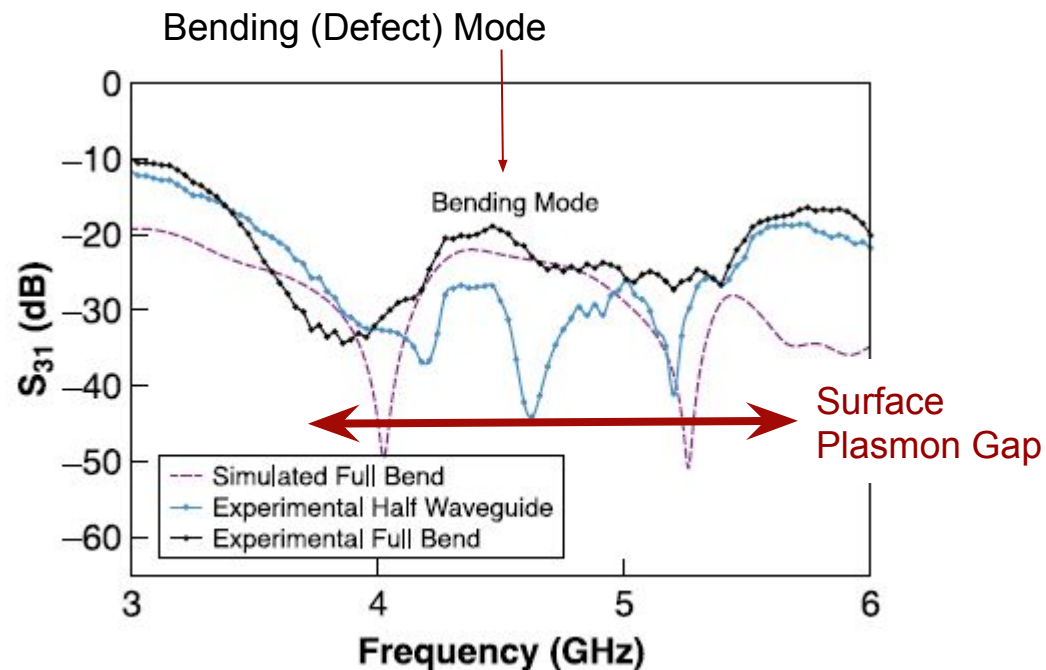


- line vacancies admit waves tuned to the defect state within the bandgap
- waves cannot propagate beyond the vacancy and are guided
- basis for on-chip optical devices

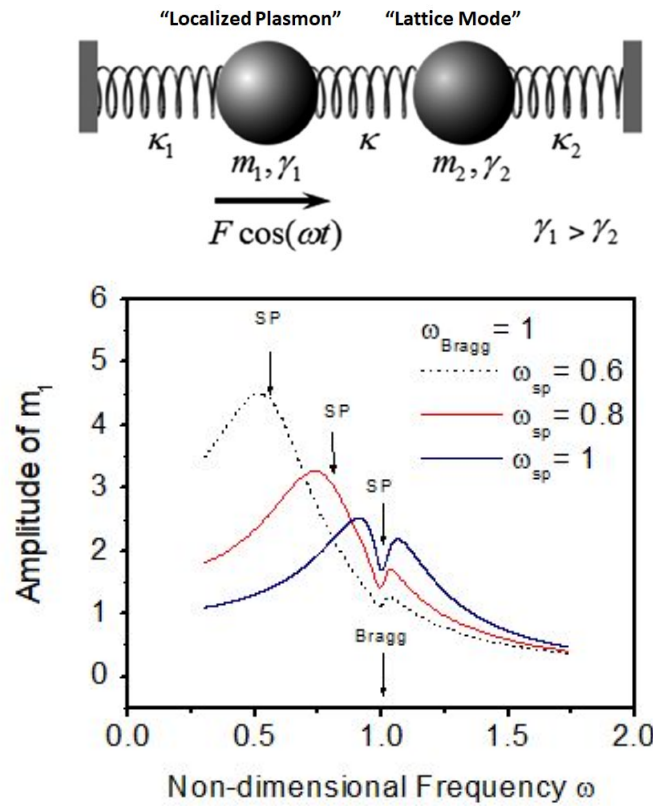
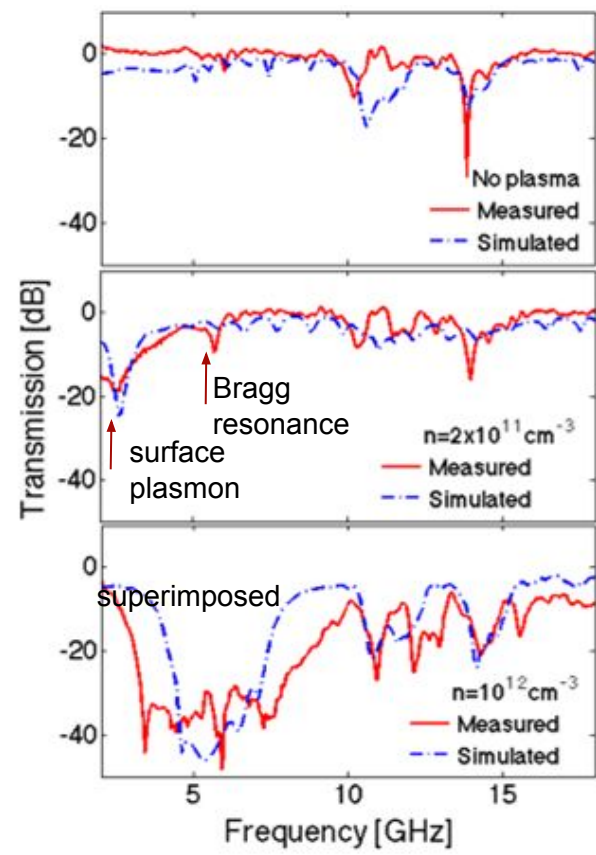
Multiple Gaps - Bragg and Surface Plasmons



Simulated Fields



Fano (Lattice) Resonance Enhancement



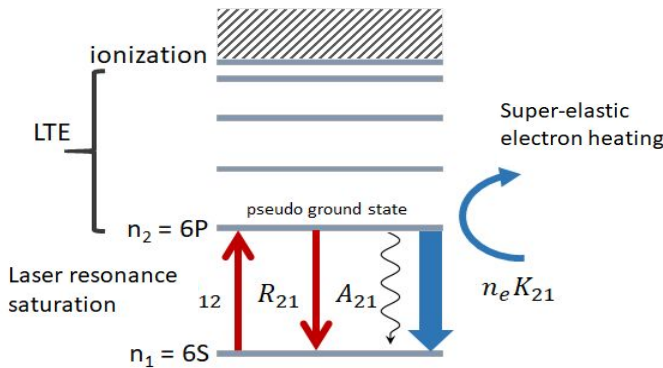
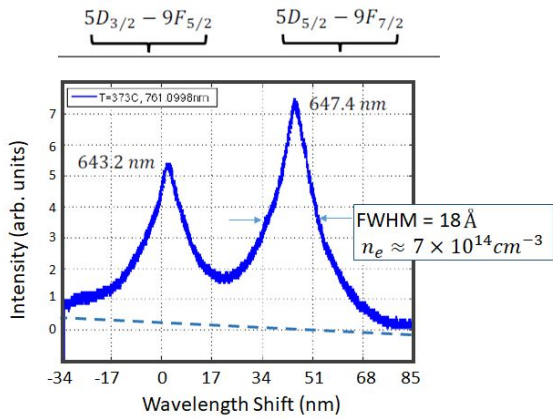
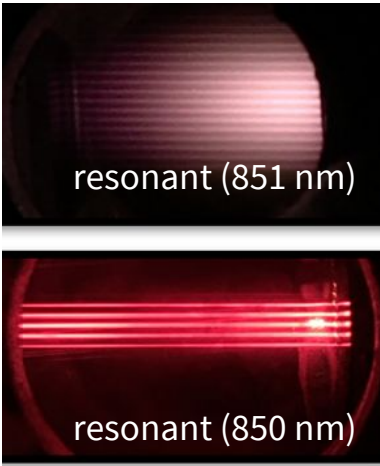
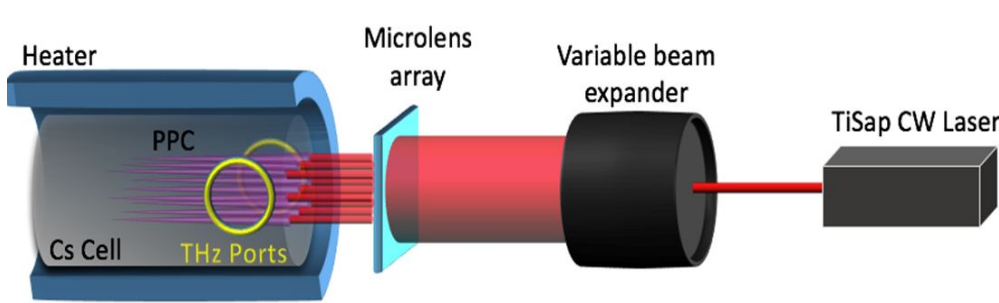
What's on the Horizon?

- THz All-Plasma PCs (Higher Plasma Density)
- 3D Plasma Photonic Crystals

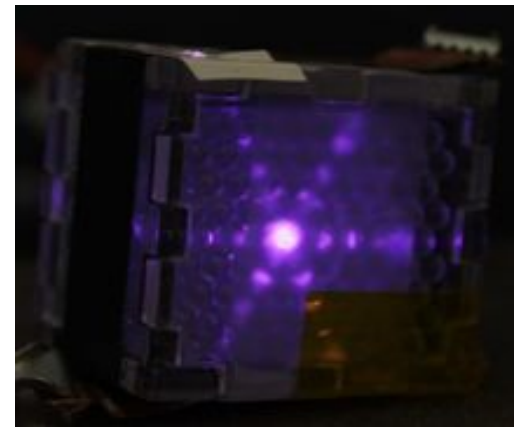
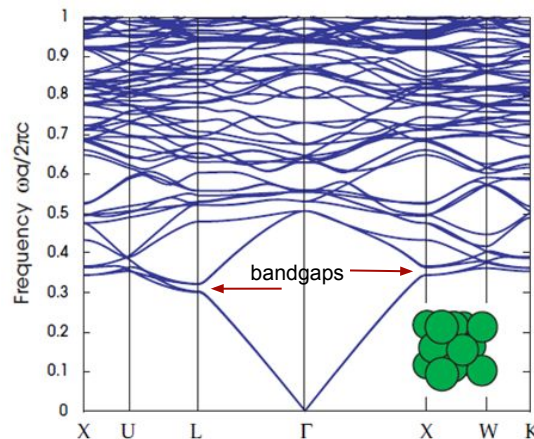
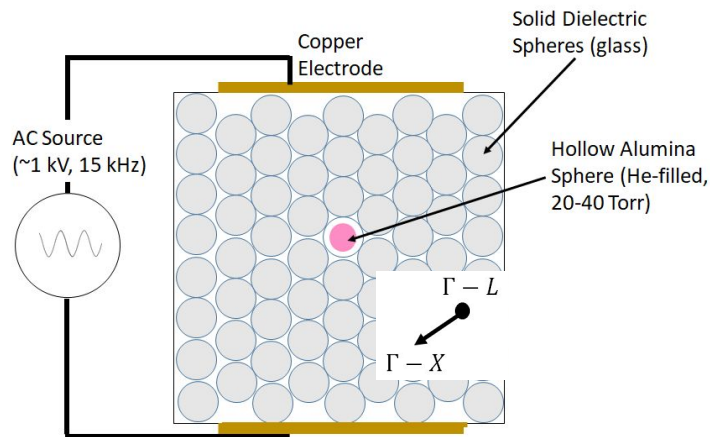
THz All-Plasma Photonic Crystal?

- Requires high plasma densities
 - $n_e = 10^{15} - 10^{16} \text{ cm}^{-3}$
- Small lattice constant ($< 300 \text{ } \mu\text{m}$)
- Continuous generation

Laser shaping and ionization



3D Photonic Crystal with Plasma-Filled Defect?



- closed-packed dielectric spheres (glass) with plasma-filled defect
- hollow alumina (He-filled) core excited with a DBD
- band-gaps expected but weak for glass ($n_{\text{glass}}/n = 1.3$)
- demonstrate plasma excitation
- redesign for solid alumina spheres with hollow-core defect

Thank you for your attention and patience!

Questions?

