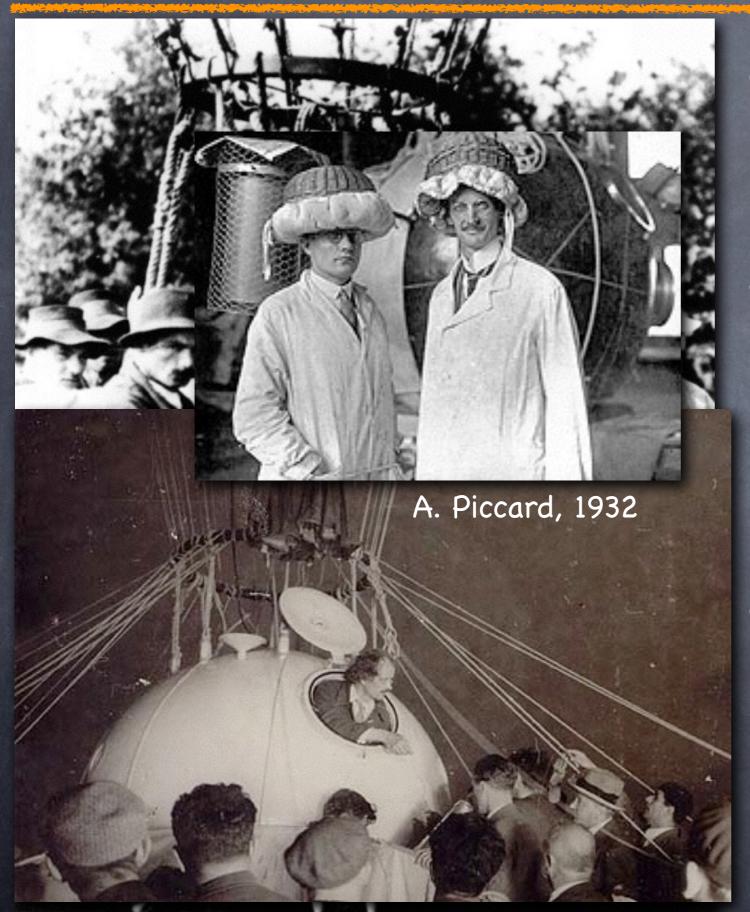
Simulations of Particle Acceleration at Astrophysical Shocks

Damiano Caprioli
Princeton University



An extraterrestrial radiation!

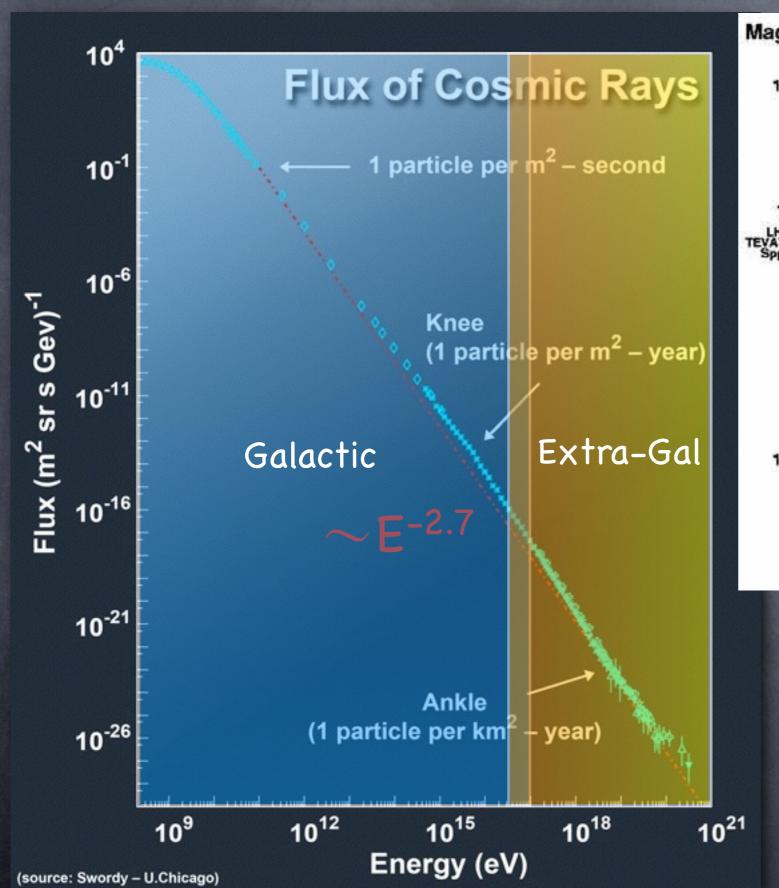


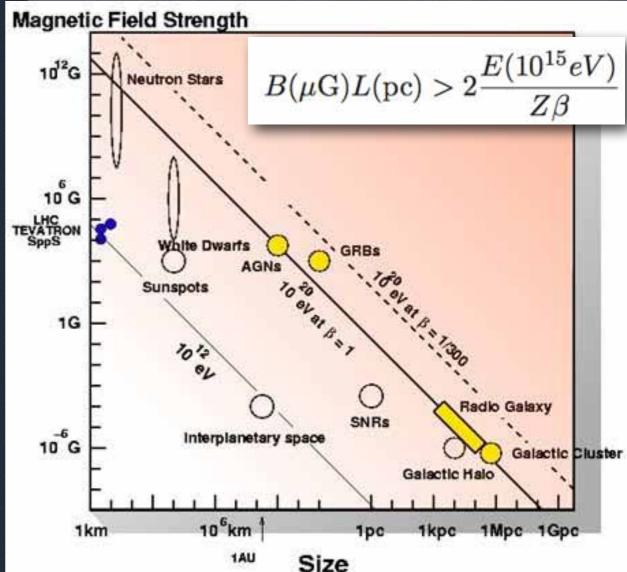


- 1912: V. Hess discovers an extraterrestrial source of ionization: Cosmic Rays
- o 1930-1932: A. Piccard reaches the stratosphere with a pressurized aluminum gondola attached to a ballon to measure CRs
- 1940: B. Rossi and P. Auger measure Extensive Air Showers:
 - CRs up to 10⁴-10⁵ GeV

The CR spectrum at Earth





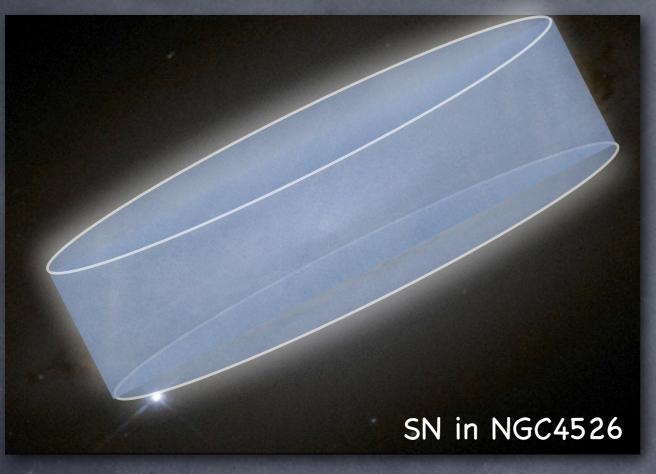


Hillas criterion: size of the system larger than the particle gyroradius

SNR paradigm: energetics



Baade-Zwicky (1934) energetic argument, updated



$$\left| \varepsilon_{\rm CR} = 0.5 eV cm^{-3} \right|$$

$$V_{conf} = \pi R^2 h = 2 \times 10^{67} cm^3$$

$$W_{CR} = \varepsilon_{CR} V_{conf} \approx 2 \times 10^{55} \text{ erg}$$

$$L_{\rm CR} \approx \frac{W_{\rm CR}}{\tau_{\rm conf}} \approx 5 \times 10^{40} \text{ erg s}^{-1}$$

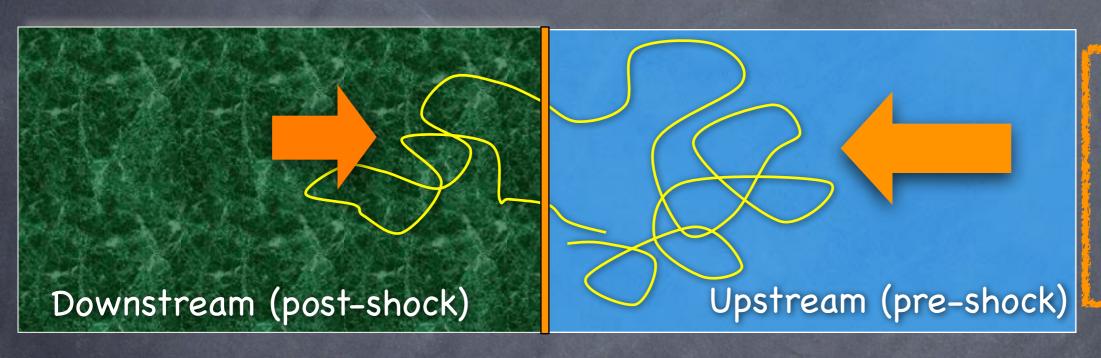
$$L_{SN} = R_{SN} E_{kin} \approx 3 \times 10^{41} \text{ erg s}^{-1}$$

10-20% of SN ejecta kinetic energy converted into CRs can account for the energetics

SNR paradigm: acceleration mechanism



- Fermi mechanism (Fermi, 1954): random scattering leads to energy gain
- In a shock a particle gains energy at any reflection (Blandford & Ostriker; Bell; Axford et al.; 1978): Diffusive Shock Acceleration (DSA)



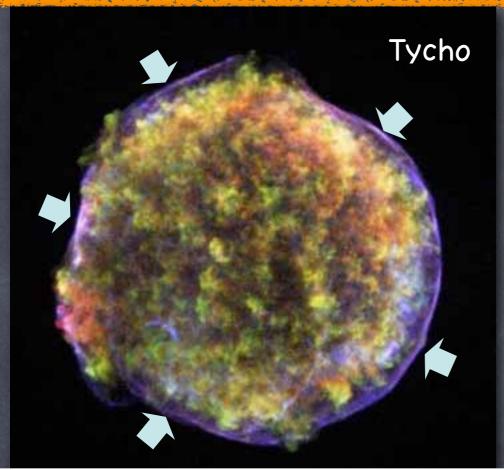
Test-particle squeezed between converging flows

OSA produces power-law $p^{-\alpha}$ in momentum, depending on the compression ratio $R = \rho_d/\rho_u$ only. For strong shocks: α =4

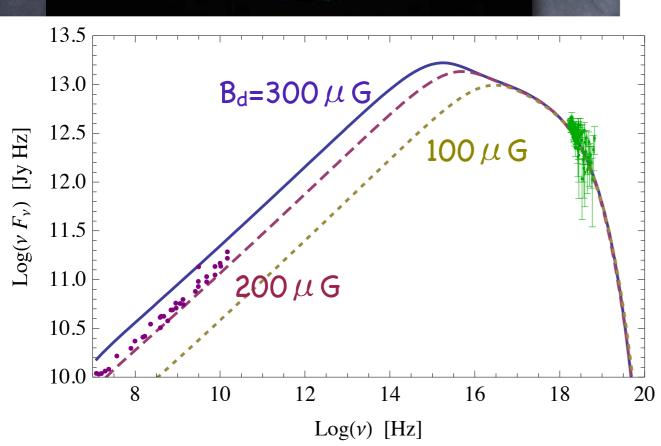
$$R = \frac{4M_s^2}{M_s^2 + 3}$$
 $\alpha = \frac{3R}{R - 1}$

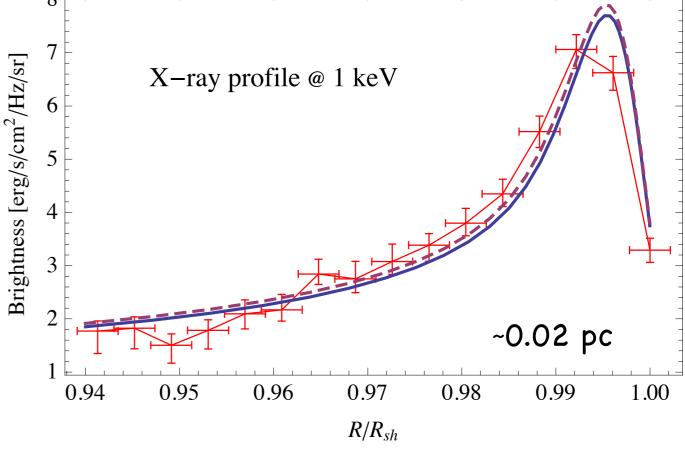
Evidence of magnetic field amplification





- Narrow (non-thermal) X-ray rims due to synchrotron losses of 10-100 TeV electrons...
- \odot ...in fields as large as B \sim 100-500 μ G





Morlino & DC, 2012

Conclusions?



Supernova Remnants

- Have the right energetics
- Diffusive shock acceleration produces power-laws
- B amplification may help reaching the knee



BUT

- Is acceleration at shocks efficient?
- How do CRs amplify the magnetic field?
- When is acceleration efficient?
- How are ions and electrons injected?



Acceleration from first principles



Full particle in cell approach

(..., Spitkovsky 2008, Niemiec+2008, Stroman+2009, Riquelme & Spitkovsky 2010, Sironi & Spitkovsky 2011, Park+2012,2015, Niemiec+2012, Guo+2014, DC+15...)

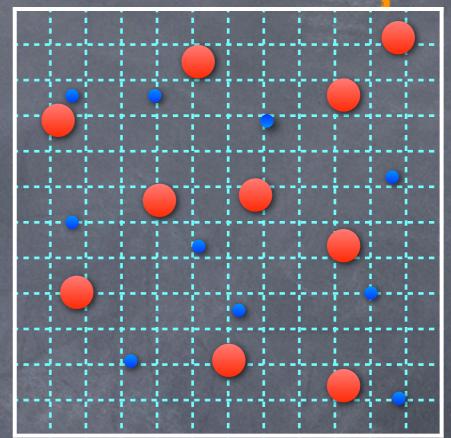
- Define electromagnetic field on a grid
- Move particles via Lorentz force
- Evolve fields via Maxwell equations
- Computationally very challenging!

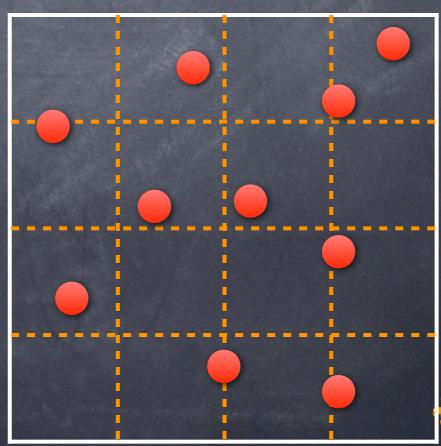
Hybrid approach:

Fluid electrons - Kinetic protons

(Winske & Omidi; Lipatov 2002; Giacalone et al.; Gargaté & Spitkovsky 2012, DC & Spitkovsky 2013-2015,...)

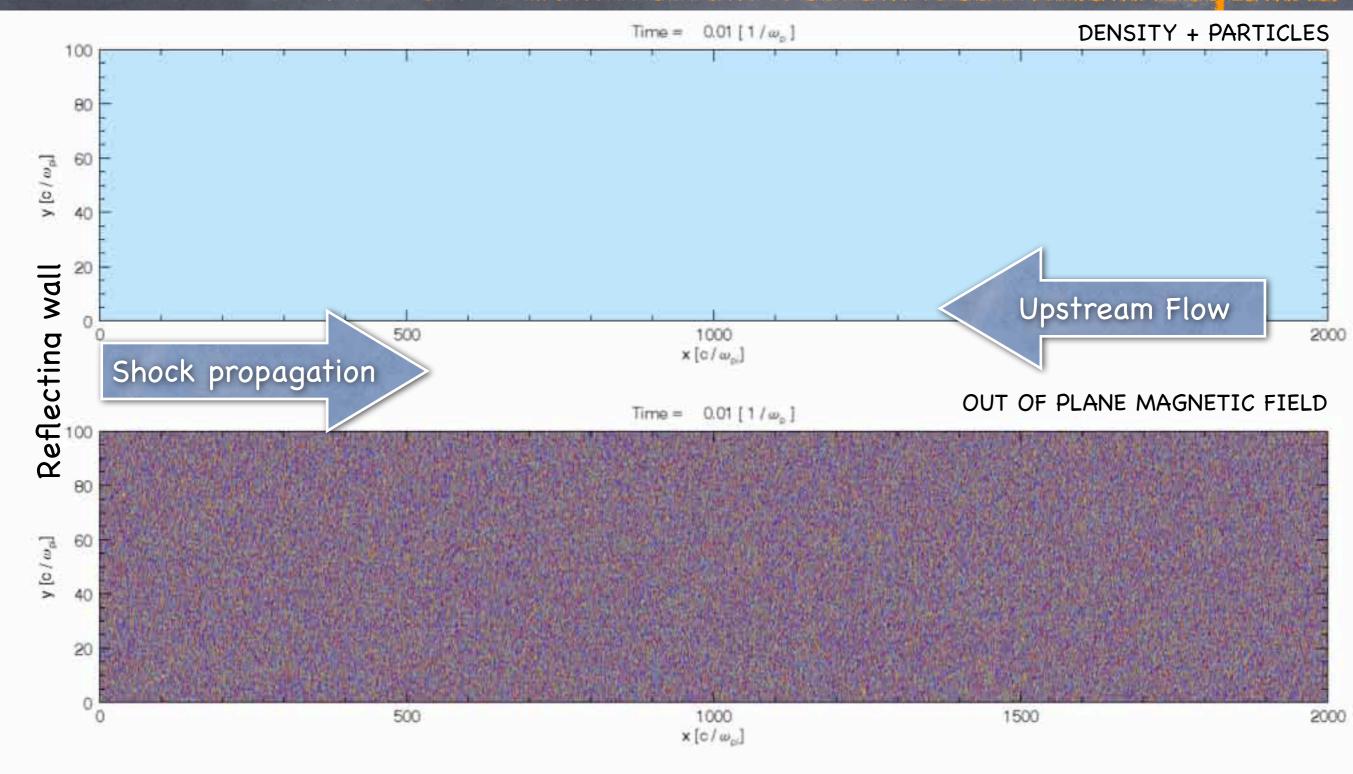
macroscopical time/length scales





Hybrid simulations of collisionless shocks



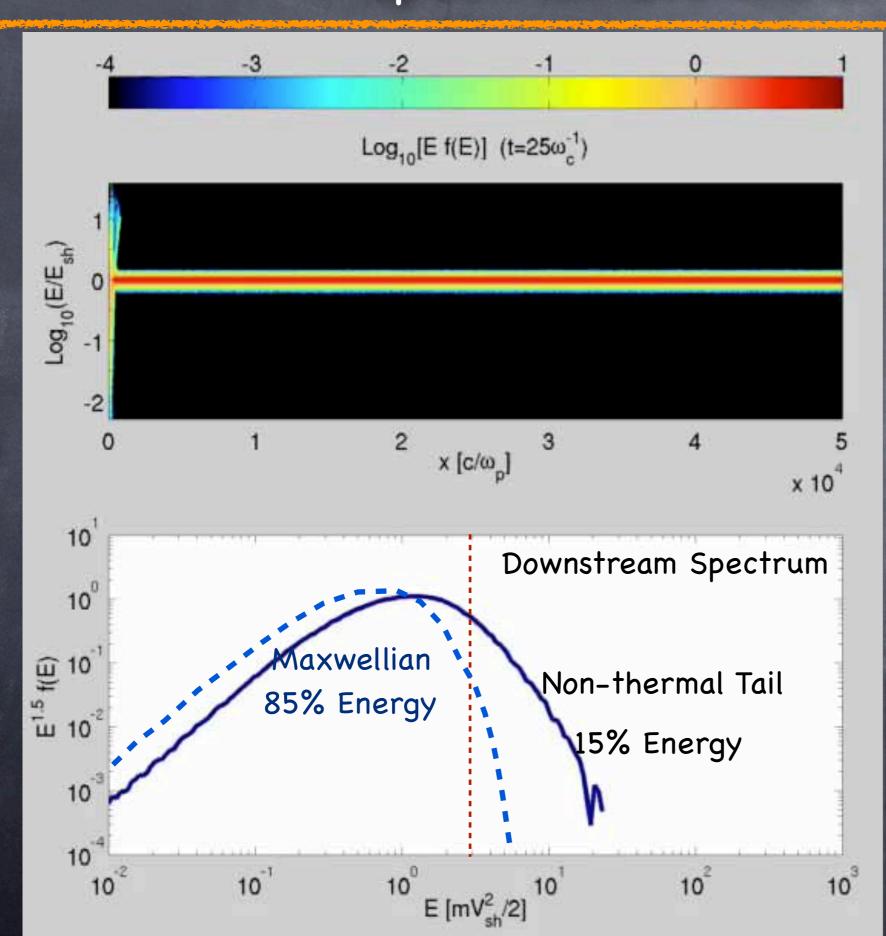






Spectrum evolution





First-order Fermi acceleration: $f(p) \propto p^{-4}$ $4\pi p^2 f(p) dp = f(E) dE$

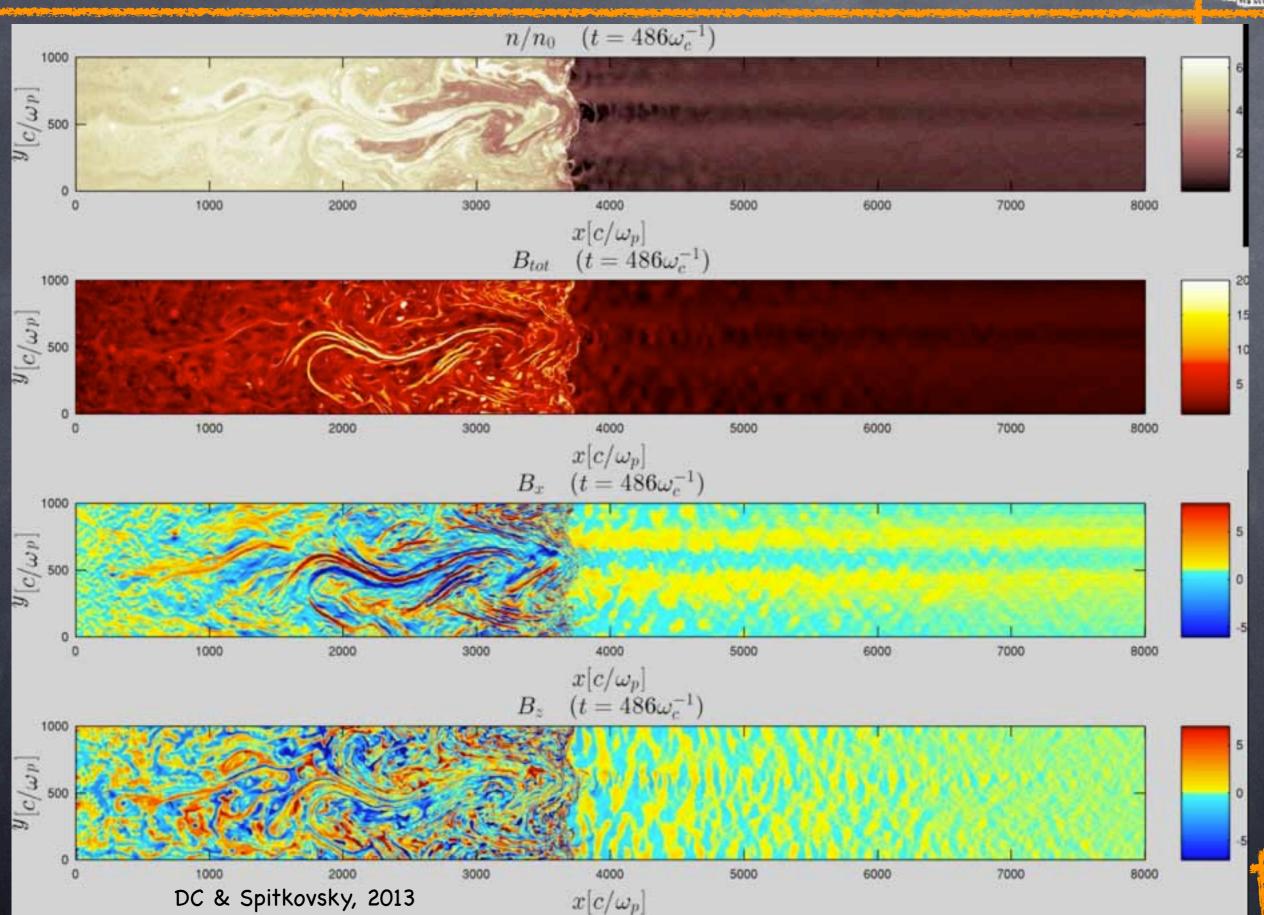


 $f(E) \propto E^{-2}$ (relativ.)

 $f(E) \propto E^{-1.5}$ (non rel.)

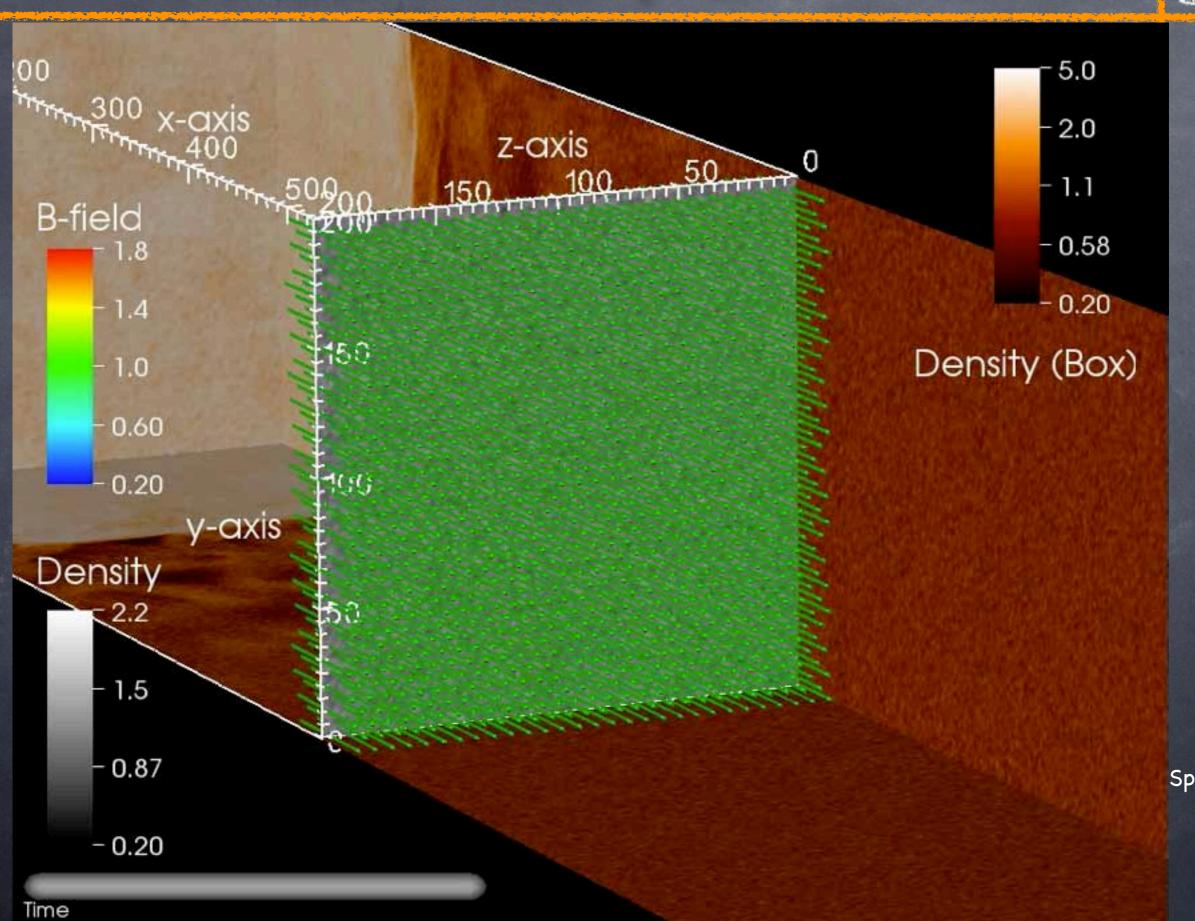
Filamentation instability





3D simulations of a parallel shock

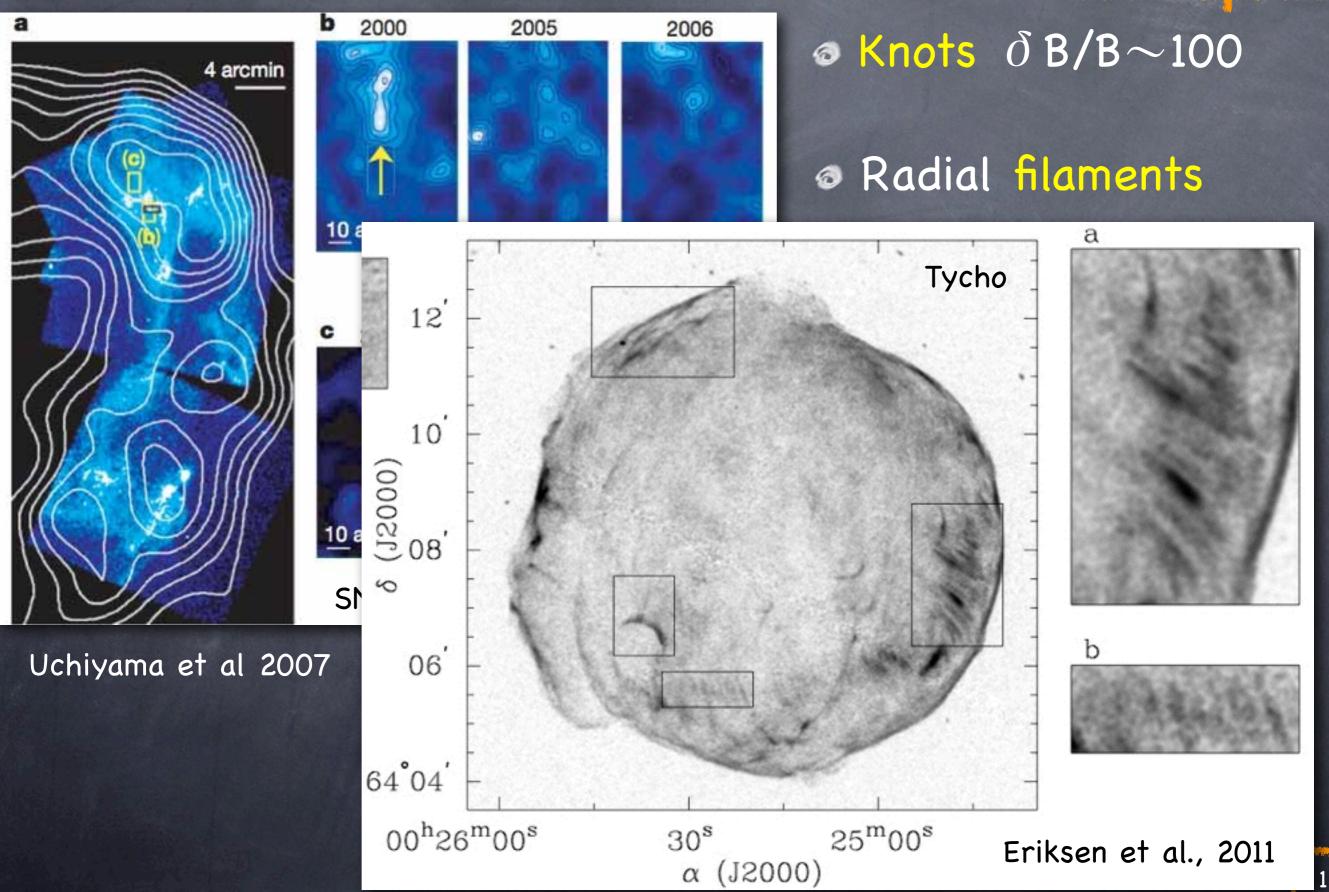




DC & Spitkovsky, 2014

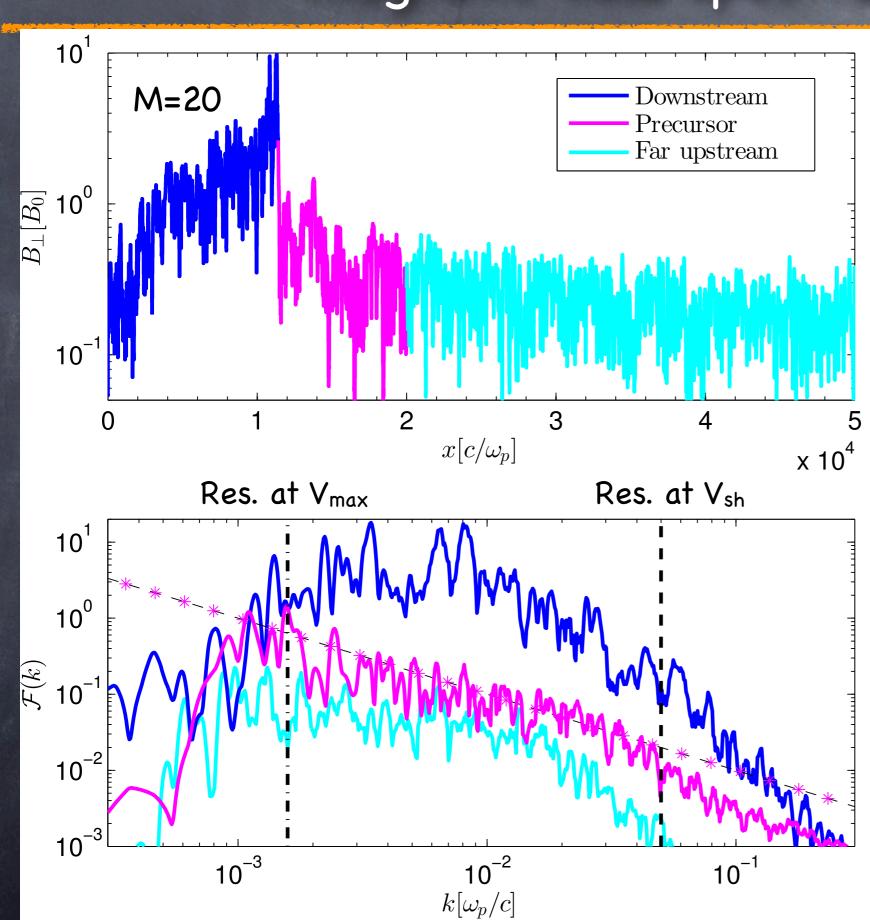
Knots and filaments





Magnetic field spectrum





Magnetic energy density per unit logarithmic band-width, F(k)

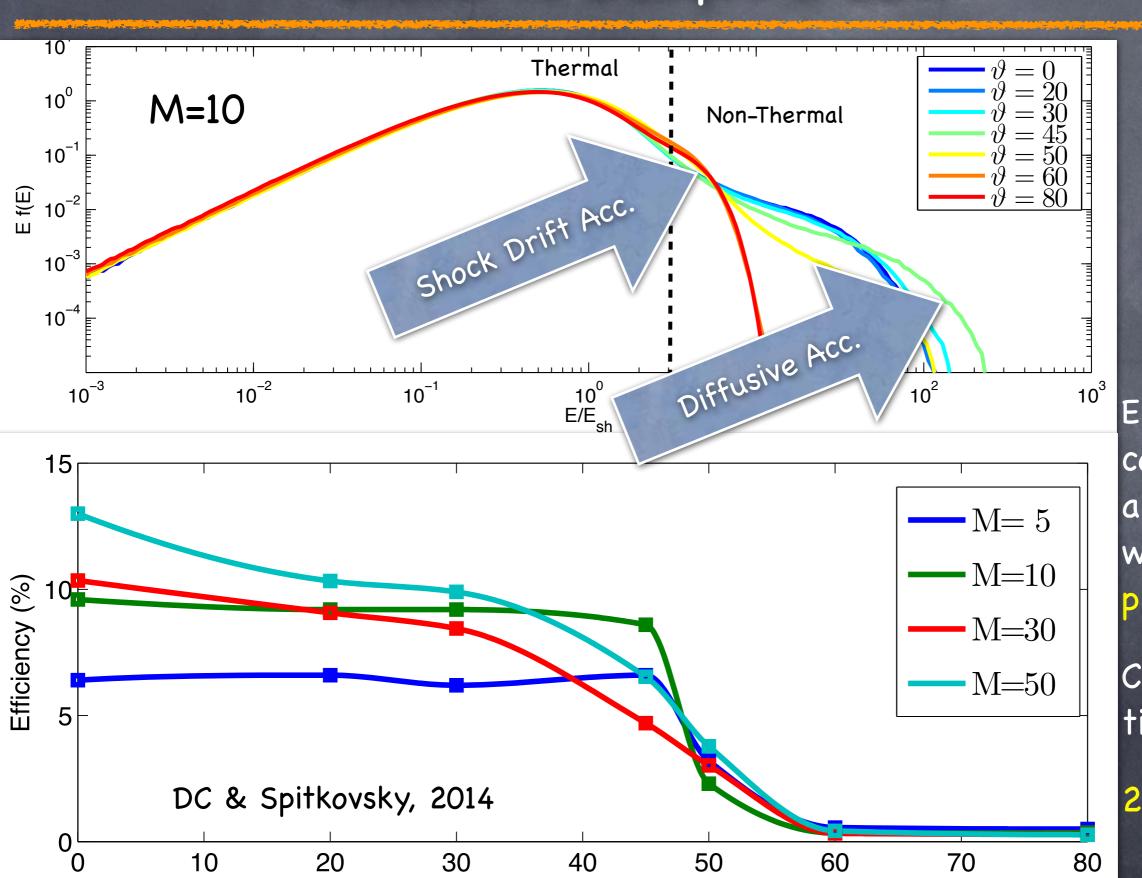
$$\frac{B_{\perp}^2}{8\pi} = \frac{B_0^2}{8\pi} \int_{k_{min}}^{k_{max}} \frac{dk}{k} \mathcal{F}(k)$$

- Turbulence selfgenerated by a spectrum ∝p-4

DC & Spitkovsky, 2014b

Parallel vs Oblique shocks





40

 θ (deg)

50

60

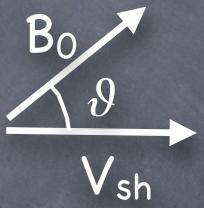
70

80

10

20

30

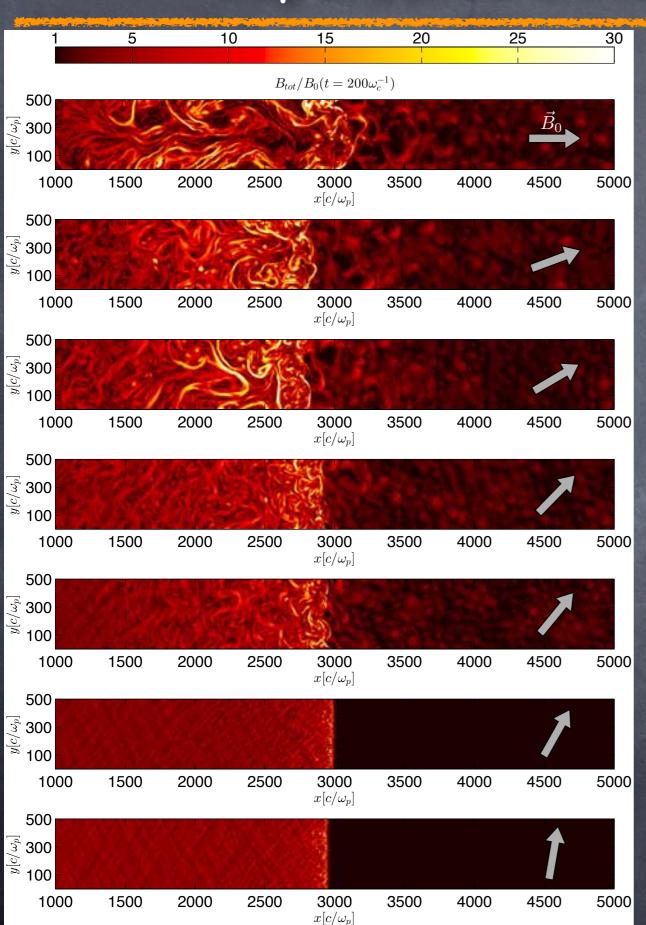


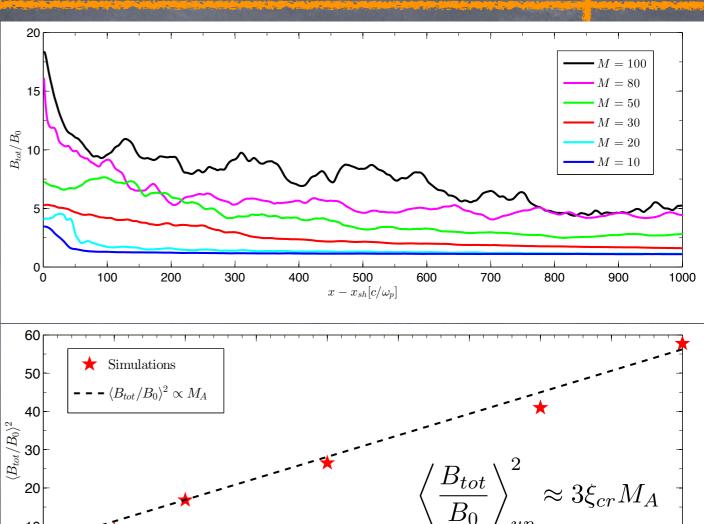
Each point corresponds to a simulation with about 109 particles

Computation time: almost 2x10⁶ cpu h

Dependence on inclination and M







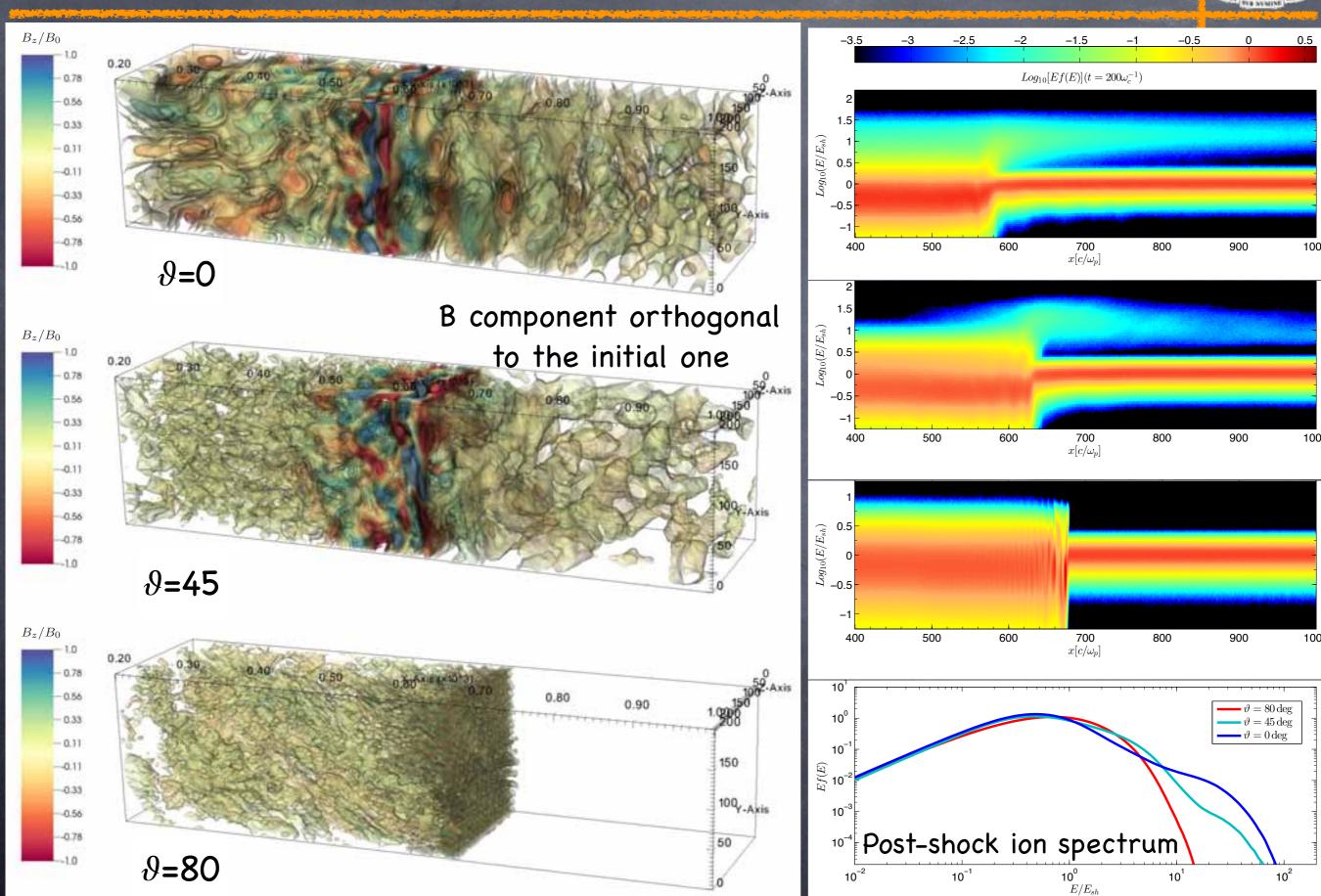
In agreement with the prediction of resonant streaming instability

 M_A

More B-field amplification for stronger shocks!

3D simulations





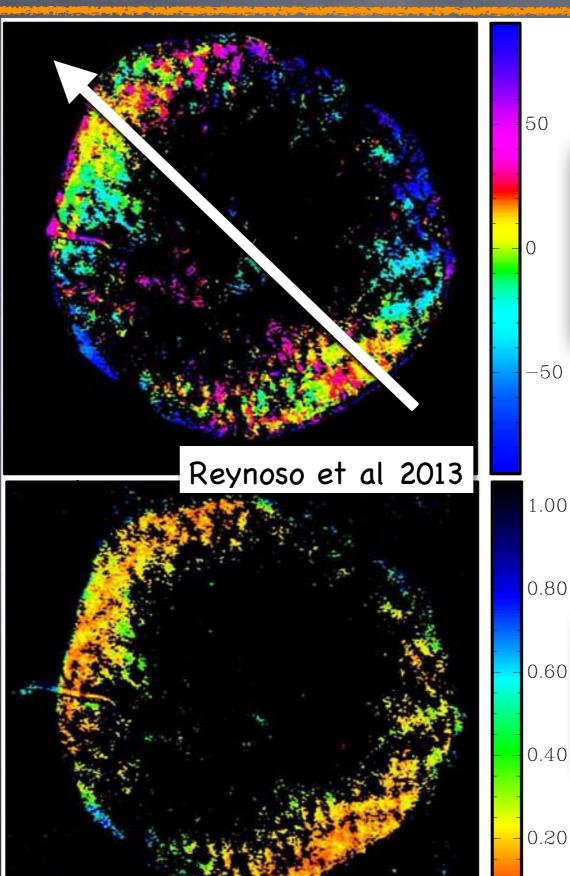
SN 1006: a parallel accelerator





X-ray emission (red=thermal white=synchrotron)

Magnetic field amplification and particle acceleration where the shock is parallel



Inclination of the B field wrt to the shock normal

1.00

0.60

0.40

Polarization (low=turbulent high=ordered)

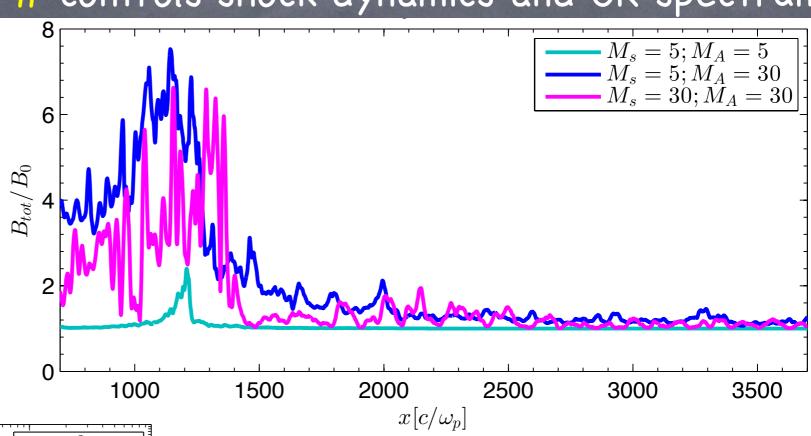
High-beta plasmas

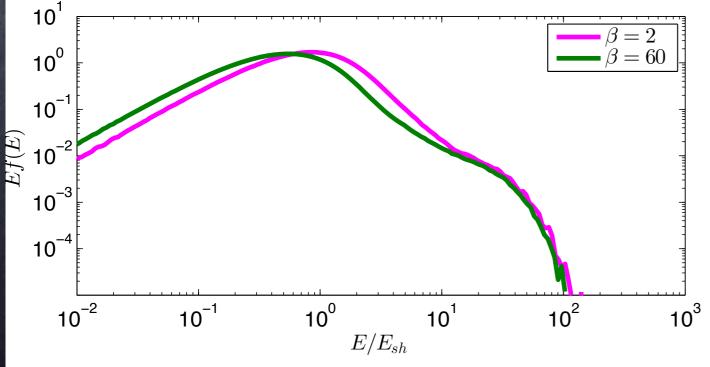


The Alfvènic Mach # controls magnetic field amplification

The (magneto-)sonic Mach # controls shock dynamics and CR spectrum

Magnetic fields are amplified also in high β plasmas!

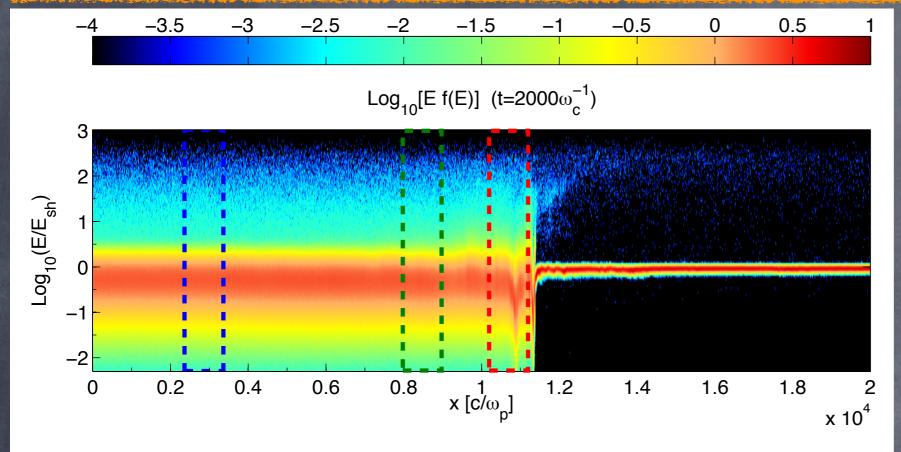


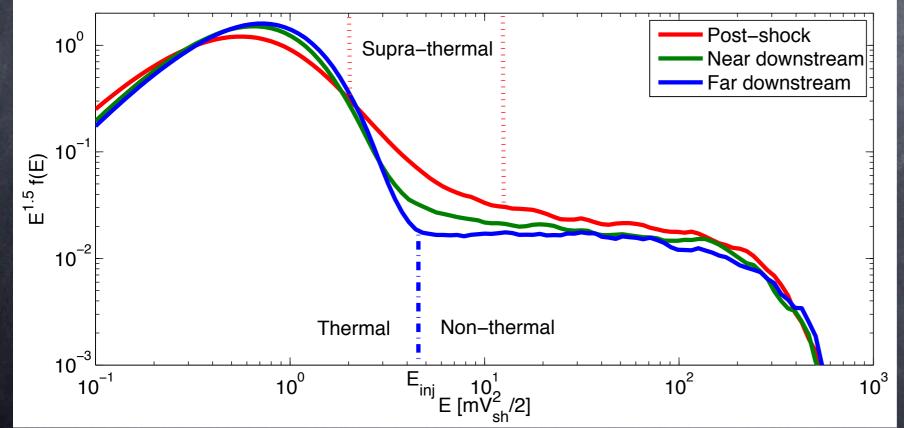


CR spectra agree with DSA prediction (steeper than p⁻⁴ for r<4)

Supra-thermal ions







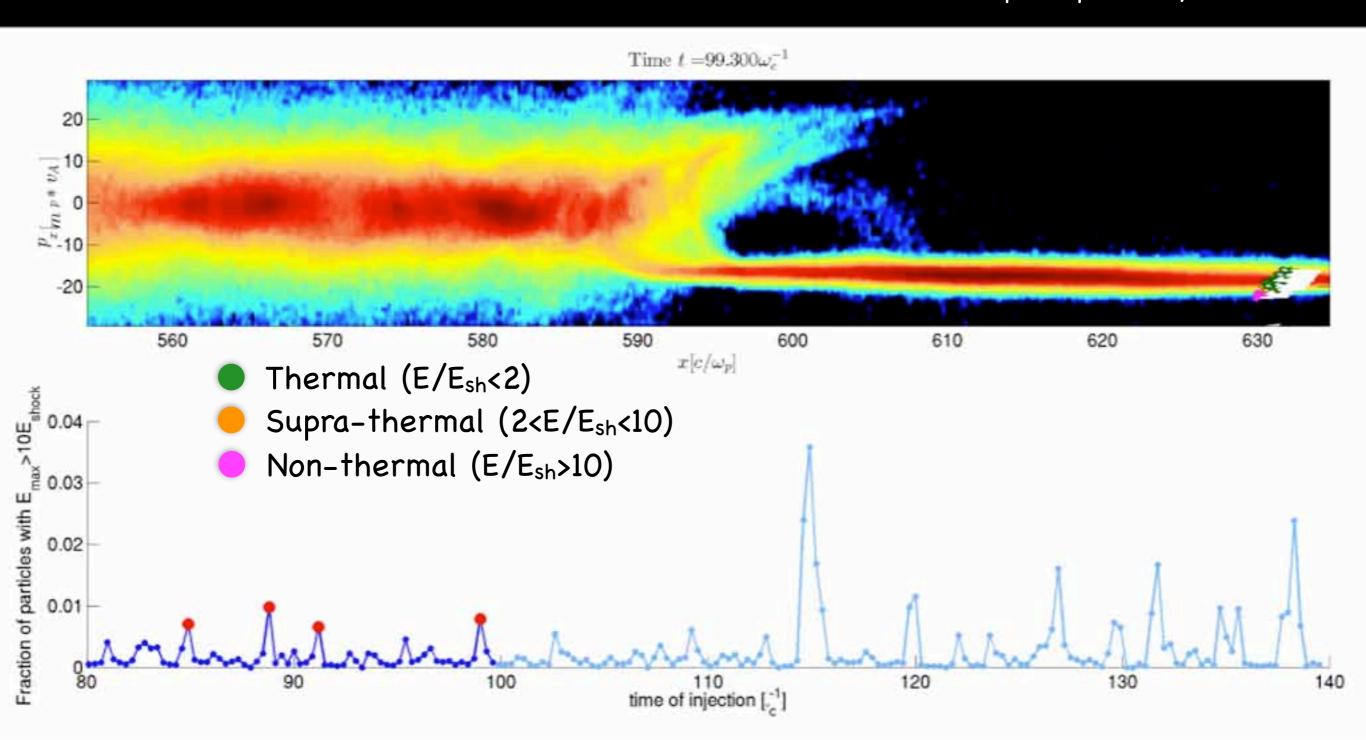
- Steep "bridge" immediately behind the shock
- Contains

 information on injection and thermalization
- The DSA powerlaw starts at Pinj∼3-4 Pth,d

Particle Injection - Simulations



DC, Pop & Spitkovsky, 2015

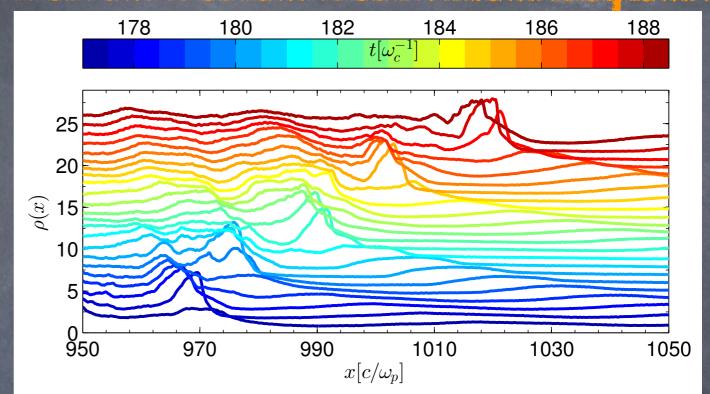


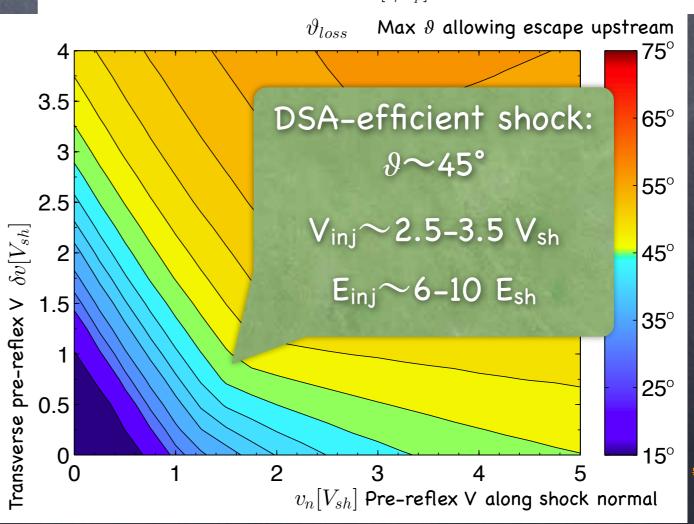
Ion Injection - Theory



Reflection off a reforming shock potential barrier (stationary in the downstream frame)

- Calculate ion trajectories (de Hoffmann-Teller frame)
 - Ion fate determined by pre-reflection velocity and shock inclination
 - \circ At given velocity, ion escape upstream of shocks with $\vartheta < \vartheta_{loss}$





Minimal Model for Ion Injection



- Time-varying potential barrier
 - High state (25% of the time)
 - -> Reflection
 - -> Shock Drift Acceleration
 - Low-state -> Thermalization
- Multiple cycles of SDA
- Spectrum à la Bell 1978

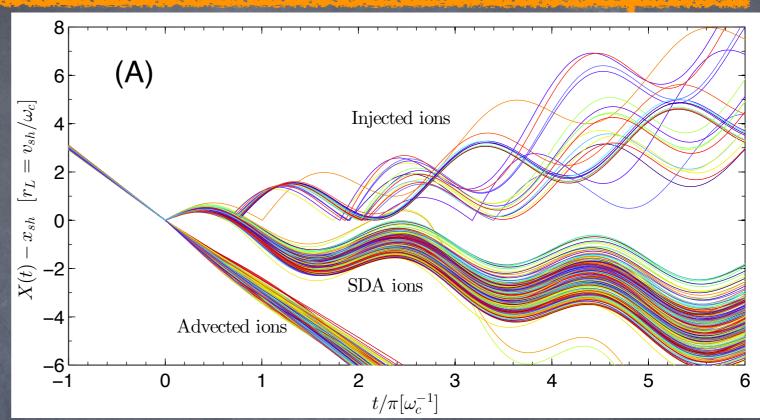
$$f(E) \propto E^{-1-\gamma}; \quad \gamma \equiv -\frac{\ln(1-\mathcal{P})}{\ln(1+\mathcal{E})}$$

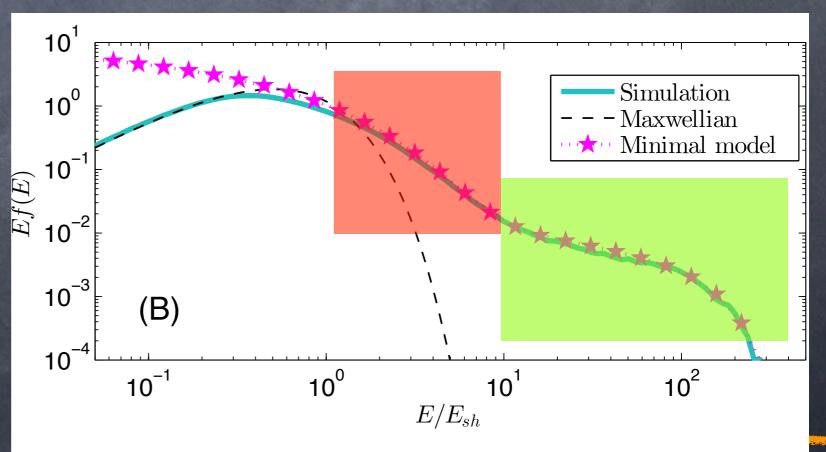
Supra-thermal

$$\bullet$$
 P=0.75; ε =2V_{sh}/v

Non-thermal

$$\bullet$$
 P=V_{sh}/v; ε =2V_{sh}/v

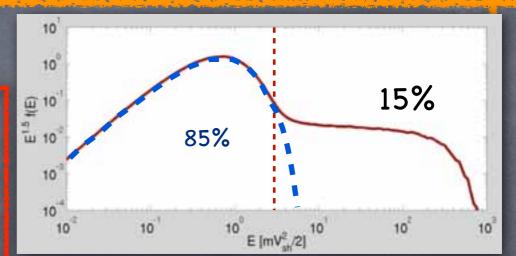


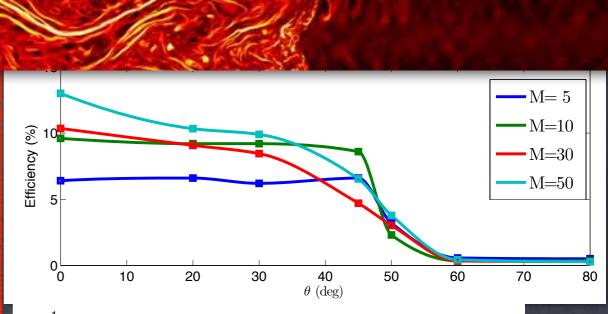


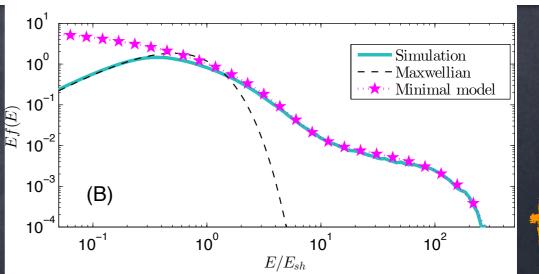
Conclusions!



- Acceleration at shocks can be efficient: >15%
- CRs amplify the B field via streaming and filamentation instabilities
- DSA efficient at parallel, strong shocks
- Ions are injected via reflection and shock drift acceleration

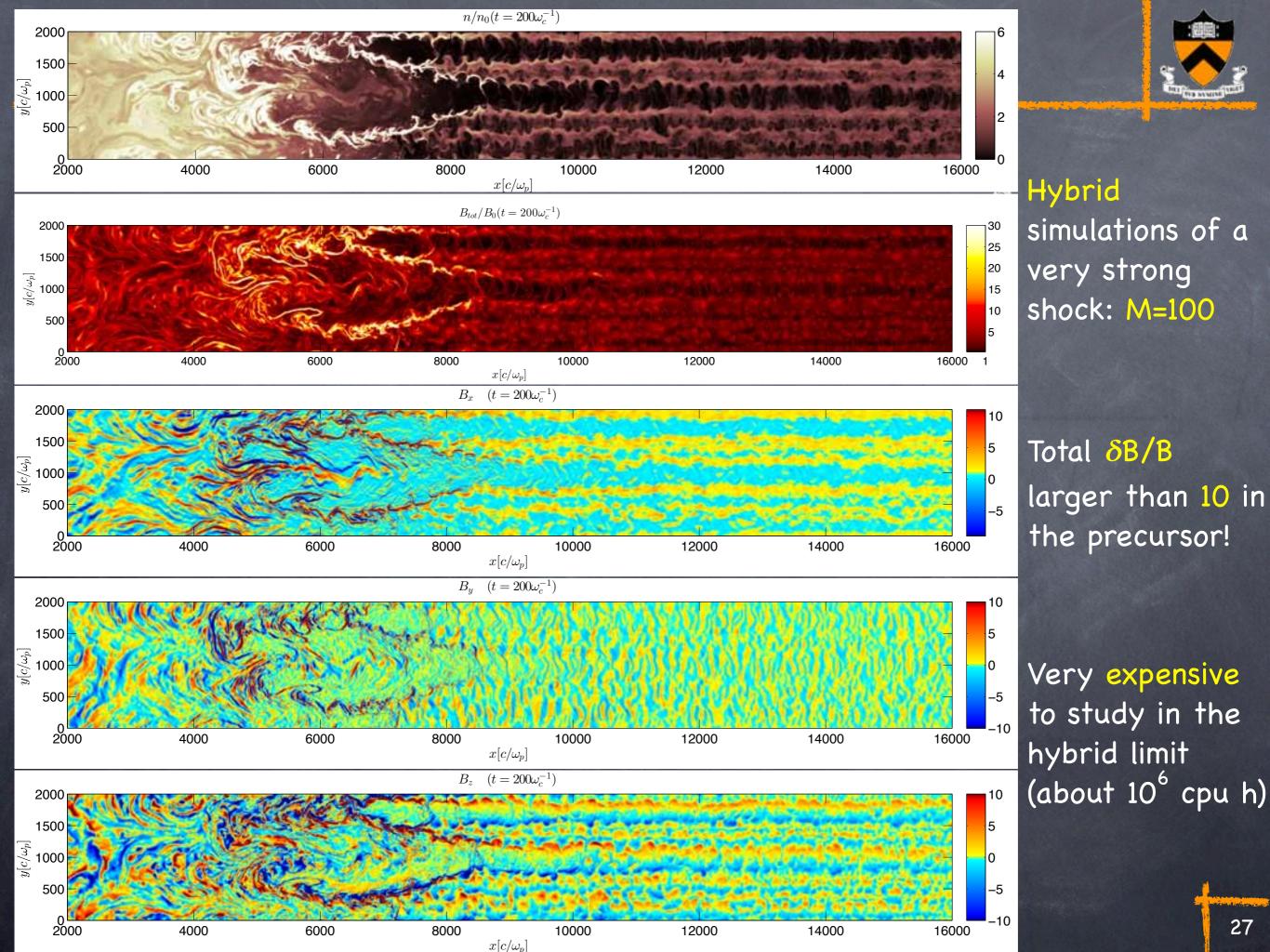






Perspectives

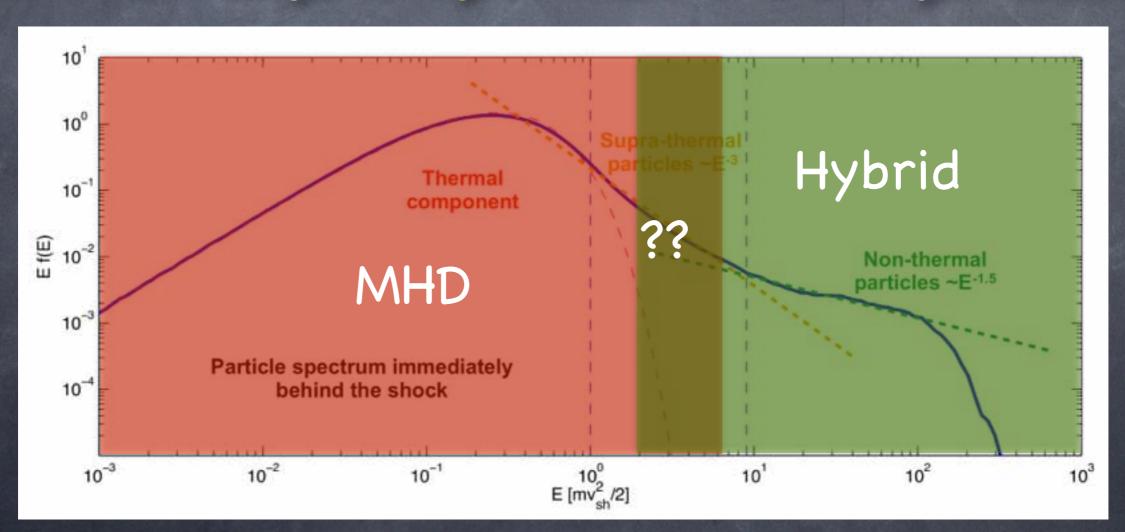
- Towards real shocks: going bigger and faster
 - Super-Hybrid (Bai, DC, Sironi, Spitkovsky 2014)
- Electron physics with full PIC (Park, DC, Spitkovsky 2015)
- Embedding microphysics in hydro/MHD simulations
 - © CRAFT: CR Analytic Fast Tool (DC et al., in prep)
- Relativistic shocks (GRB, AGN jets, pulsars, radio-SNe,...)
 - Partitioning energy into ions, electrons, and magnetic fields
- Almost any problem in collisionless astro and lab plasmas!



Going bigger: Super-Hybrid

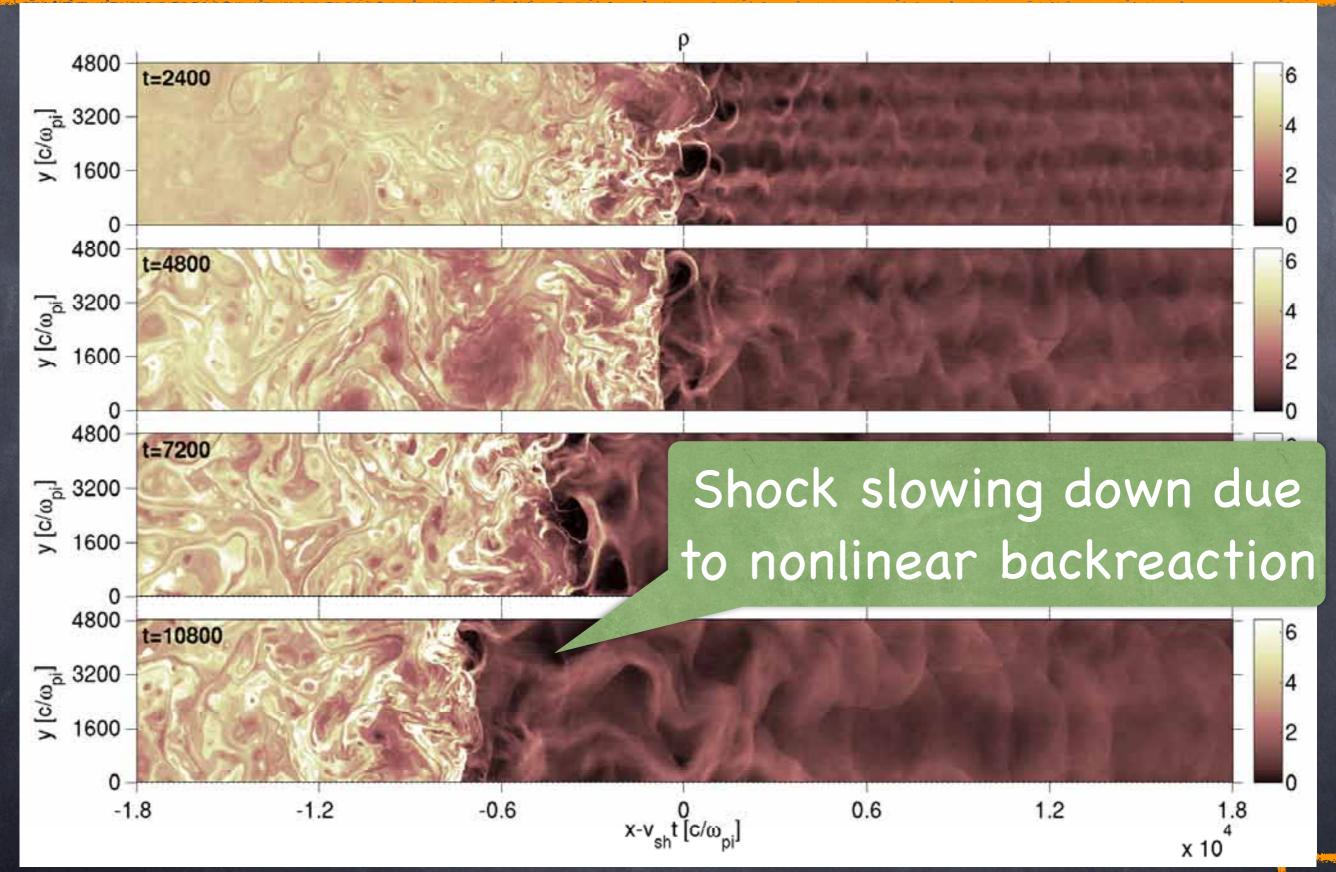


- MHD (Athena) + kinetic ions (also relativistic)
- Needs injection (tuned form hybrid)
- Allows to go to higher Mach # and larger scales



Long-term evolution

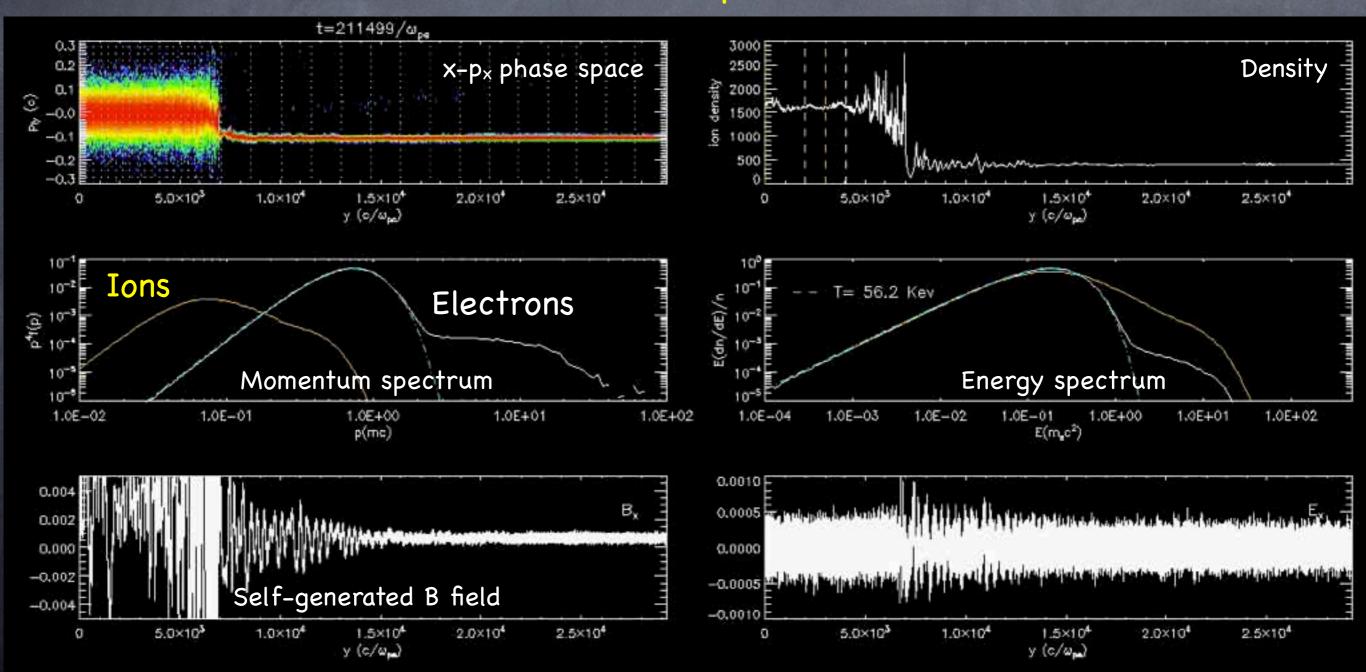




Electron/ion acceleration



- Full PIC simulations: Tristan-MP (Park, DC, Spitkovsky 2015, acc. to PRL)
 - M=20, v_{sh}=0.1c, quasi-parallel shock
 - Electrons are accelerated, but ele/proton ratio is a few %



CRAFT: a Cosmic-Ray Fast Analytic Tool



(Caprioli et al. 2009-2015, to be publicly released soon)

Iterative solution of the CR transport equation:

$$\tilde{u}(x) \frac{\partial f(x,p)}{\partial x} = \frac{\partial}{\partial x} \left[D(x,p) \frac{\partial f(x,p)}{\partial x} \right] + \frac{p}{3} \frac{\mathrm{d} \tilde{u}(x)}{\mathrm{d} x} \frac{\partial f(x,p)}{\partial p} + Q(x,p)$$
 $Q(x,p) = \eta \frac{\rho_1 u_1}{4\pi m_p p_{inj}^2} \delta(p - p_{inj}) \delta(x)$

$$Q(x,p) = \eta \frac{\rho_1 u_1}{4\pi m_p p_{inj}^2} \delta(p - p_{inj}) \delta(x)$$

Injection

conservation eqs. $\frac{p(x)}{\rho(x)^{\gamma}} = \frac{p_0}{\rho_0^{\gamma}};$ Mass+momentum

$$rac{p(x)}{
ho(x)^{\gamma}} = rac{p_0}{
ho_0^{\gamma}};$$

$$ho(x)u(x) =
ho_0 u_0$$

$$ho(x)u(x)^2 + p(x) + p_{cr}(x) + p_B(x) =
ho_0 u_0^2 + p_{g,0} + p_{B,0}$$

PB + Pcr

$$2\tilde{u}(x)\frac{dp_B(x)}{dx} = v_A(x)\frac{dp_{cr}(x)}{dx} - 3p_B(x)\frac{d\tilde{u}(x)}{dx}$$

Magnetic turbulence transport eq.

$f(x,p) = f_2(p) \exp \left[-\int_x^0 dx' \frac{\tilde{u}(x')}{D(x',p)} \right] \left[1 - \frac{W(x,p)}{W_0(p)} \right]$

$$\Phi_{esc}(p) = -D(x_0, p) \left. \frac{\partial f}{\partial x} \right|_{x_0} = -\frac{u_0 f_2(p)}{W_0(p)};$$

$$W(x,p) = \int_{x}^{0} dx' \frac{u_0}{D(x',p)} \exp \left[\int_{x'}^{0} dx'' \frac{\tilde{u}(x'')}{D(x'',p)} \right].$$

$$f_2(p) = \frac{\eta n_0 q_p(p)}{4\pi p_{inj}^3} \exp\left\{-\int_{p_{inj}}^p \frac{dp'}{p'} q_p(p') \left[U_p(p') + \frac{1}{W_0(p')}\right]\right\}$$

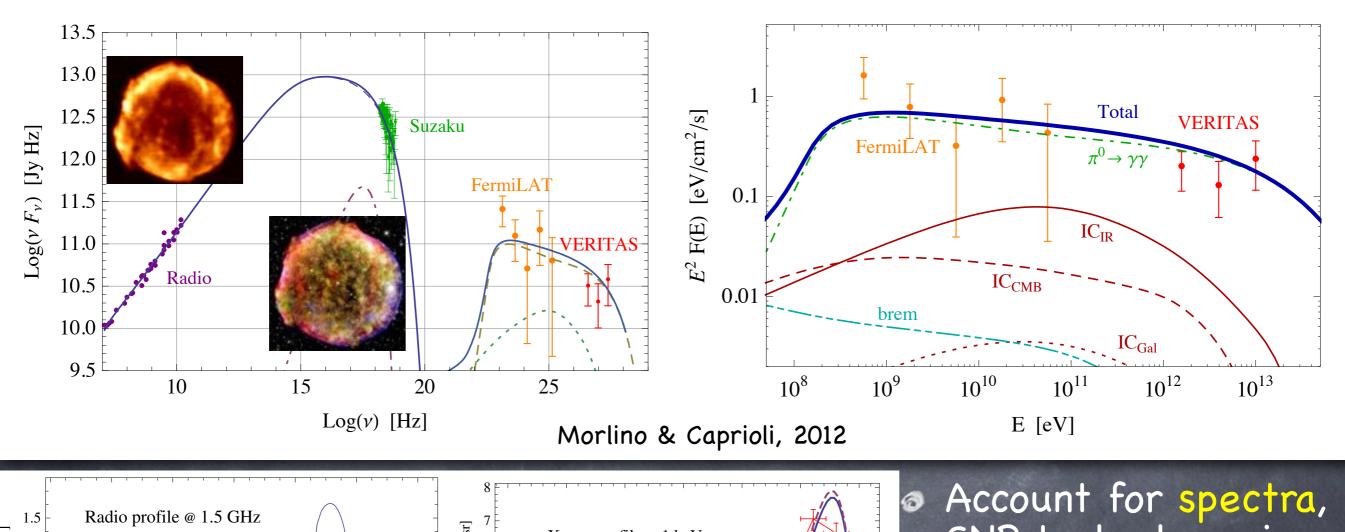
$$U_p(p) = \frac{\tilde{u}_1}{u_0} - \int_{x_0}^0 \frac{dx}{u_0} \left\{ \frac{\partial \tilde{u}(x)}{\partial x} \exp\left[-\int_x^0 dx' \frac{\tilde{u}(x')}{D(x',p)} \right] \left[1 - \frac{W(x,p)}{W_0(p)} \right] \right\}$$

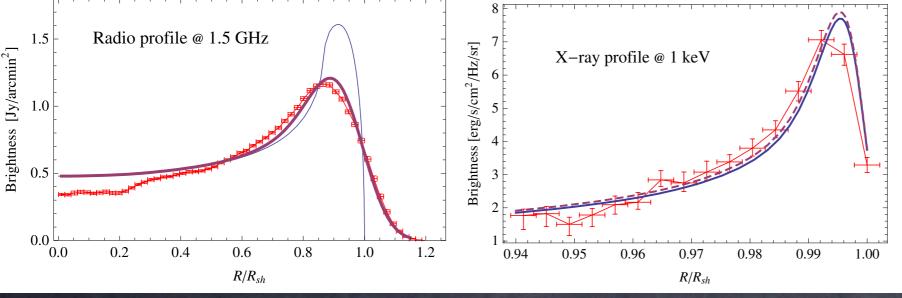
CR distribution function

- Very fast: a few seconds on a laptop (vs days on clusters)
- Embeds microphysics from kinetic simulations into (M)HD

Tycho: a clear-cut hadronic accelerator







- Account for spectra, SNR hydrodynamics, and morphology
- Hadron acc. eff. $\sim 10\%$
- Protons up to 0.5 PeV

Only two free parameters: injection efficiency and electron/proton ratio ...

Thank you!

