

Three-dimensional magnetic field line reconnection involving flux ropes and Alfvén waves

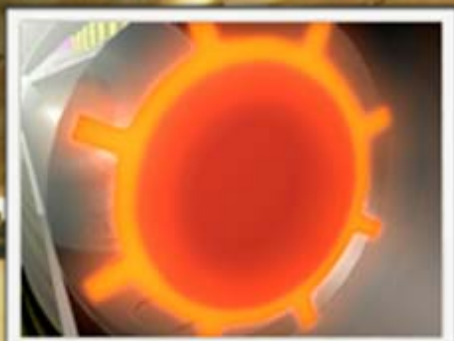
Walter Gekelman, E. Lawrence, A. Collette,
S. Vincena: Department of Physics, UCLA

Recurrent questions in Space and Astrophysics

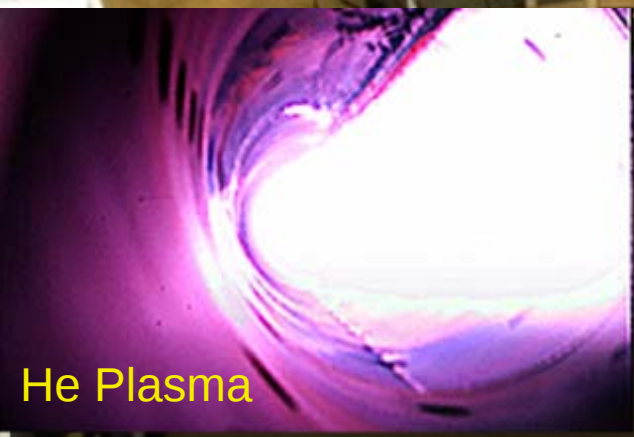
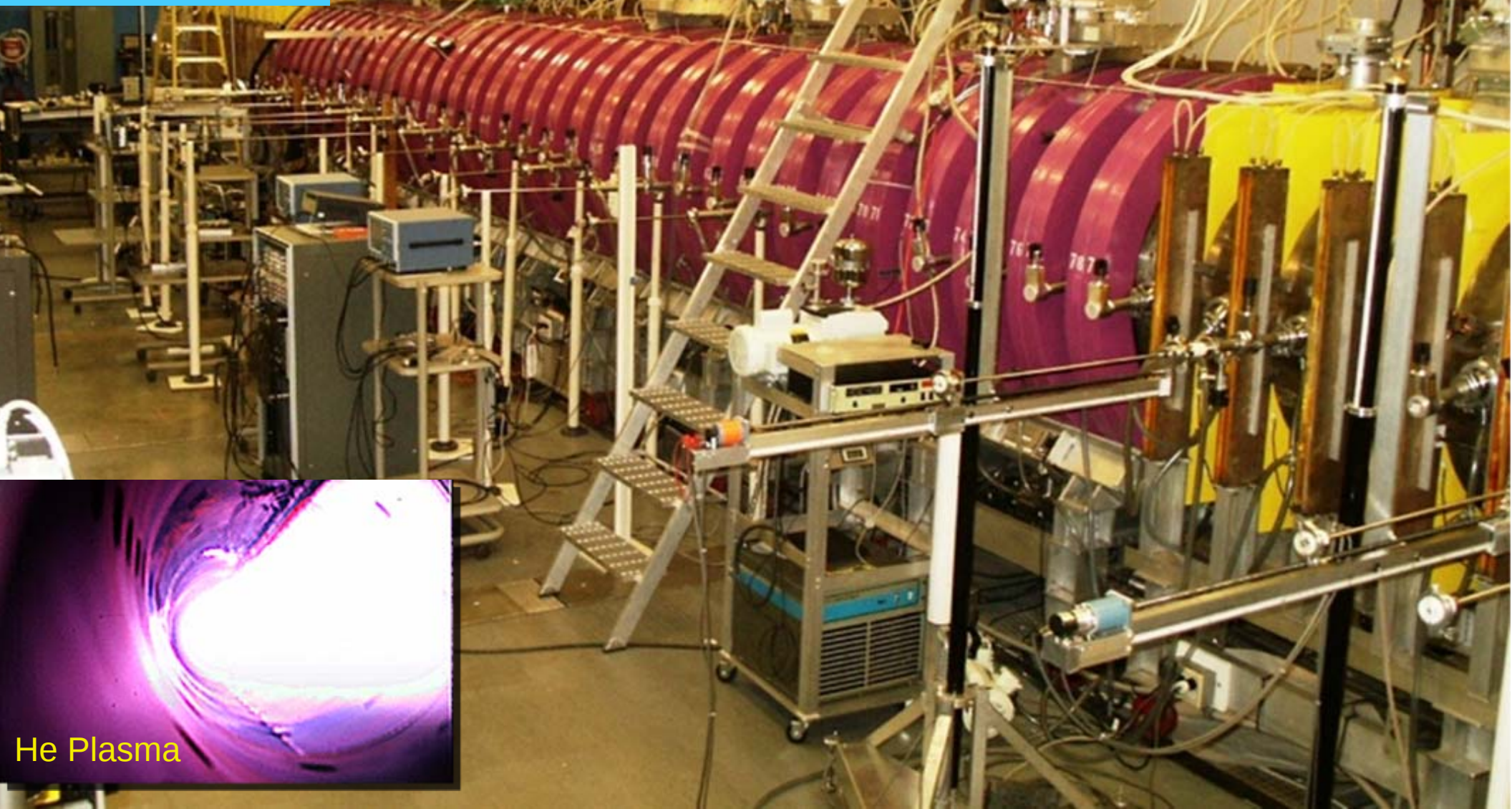
- Magnetic field generation (magnetic helicity)
- Relaxation of complex field geometries into simpler ones
- Magnetic Field Line Reconnection (2D, 3D?)
- role of guide field in reconnection
- Generation and interaction of magnetic flux ropes
- Role of waves (whistlers and Alfvén)
- Bursty verses steady phenomena (plasma instabilities)

Can we learn something from laboratory experiments?

LAPD device - UCLA



BaO Cathode



He Plasma

Plasma Production

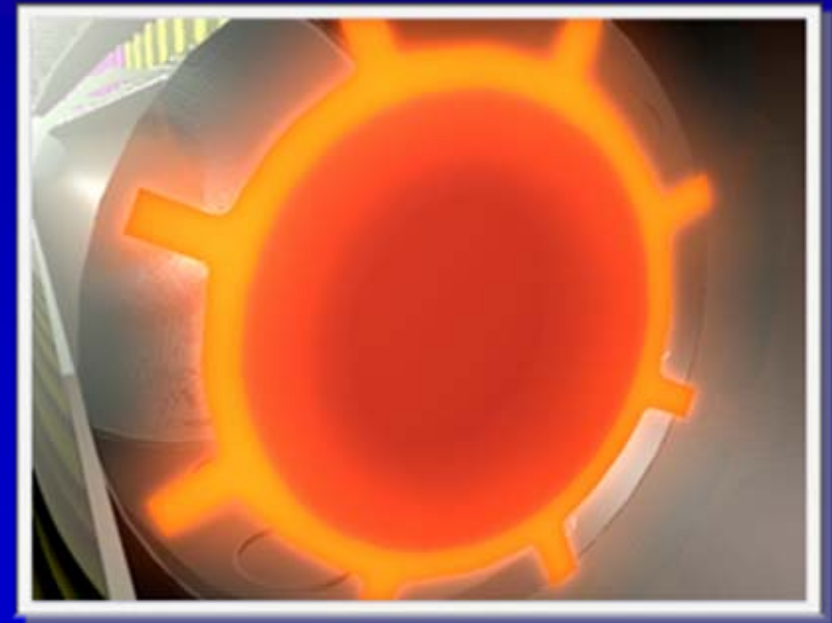
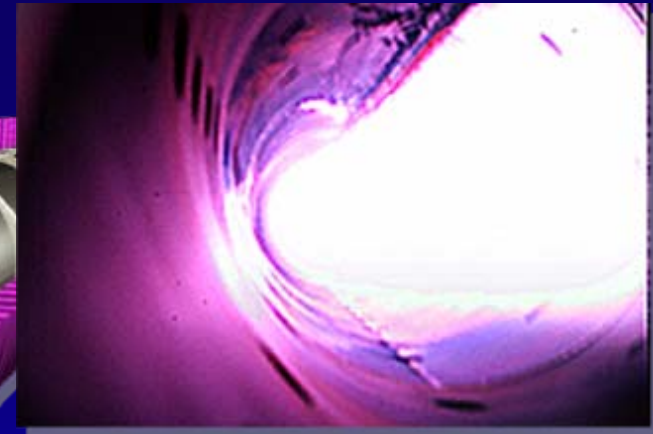
He Plasma

magnets

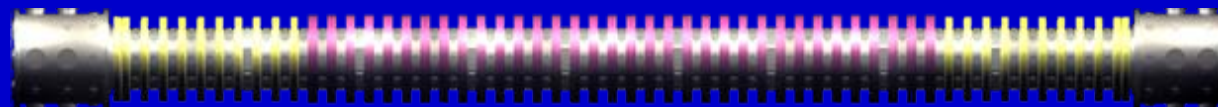
anode

cathode

Cathode - Anode

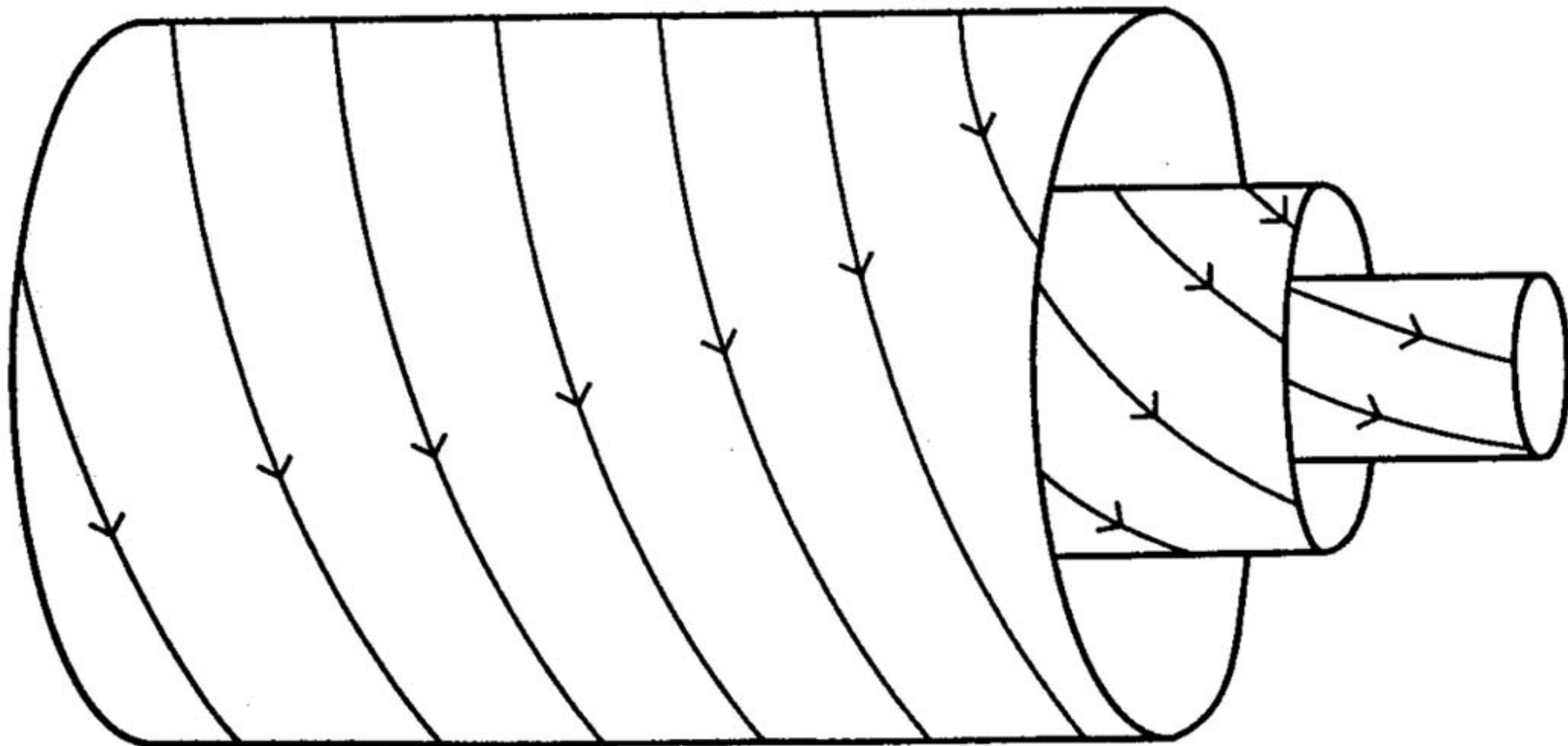


UCLA



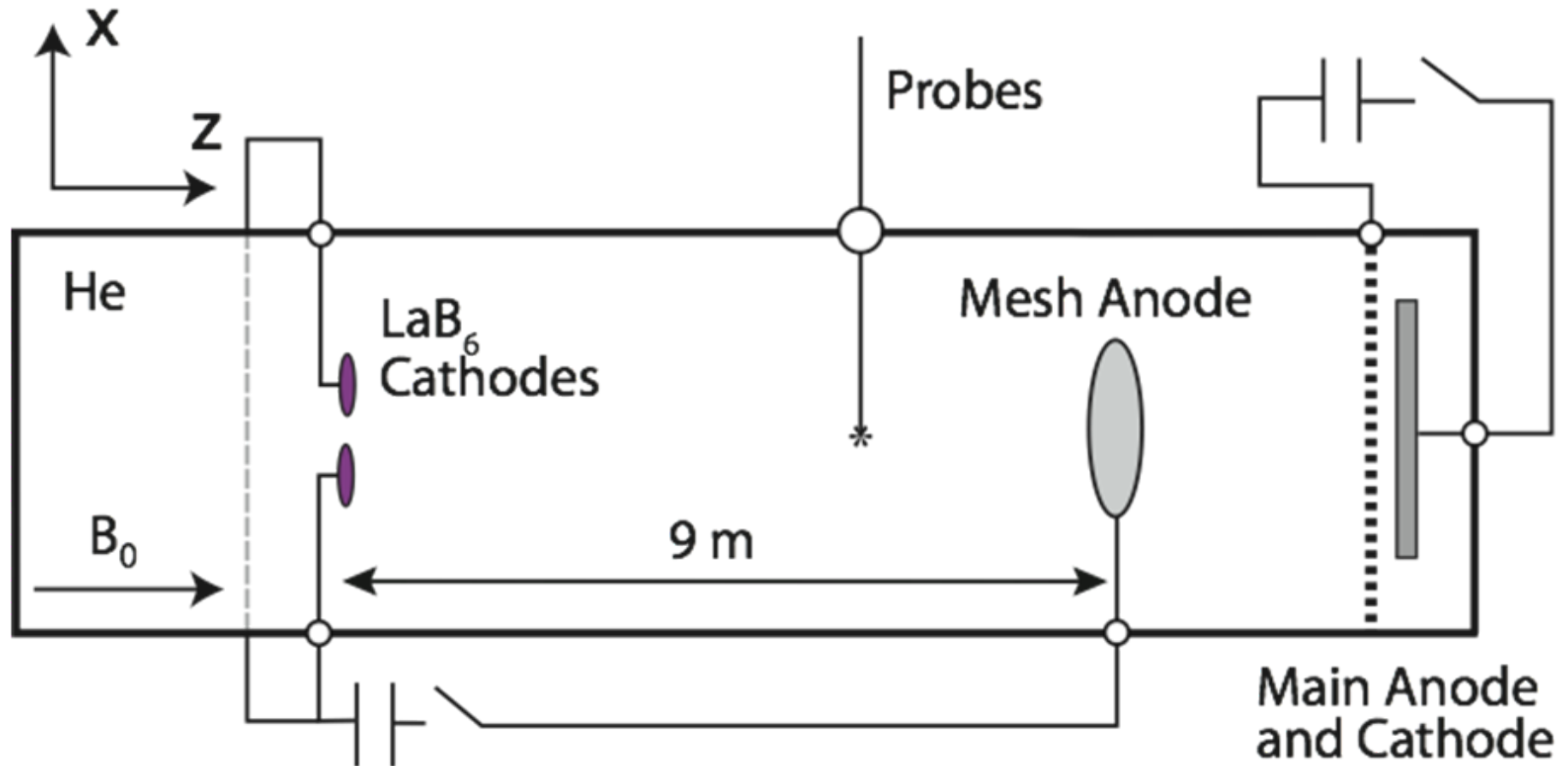
BAPSF

Magnetic Flux Ropes



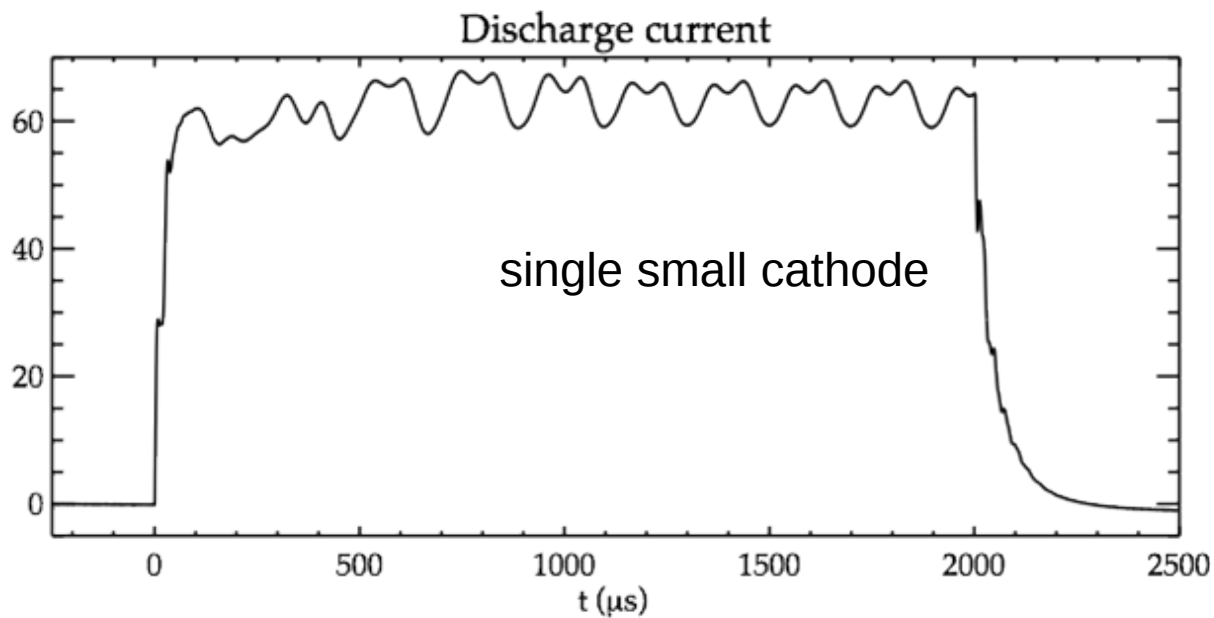
Making Flux Ropes

$$\delta x = \delta y = 3\text{mm} \quad , \quad \delta z = 64\text{ cm}, \quad \delta t = 40\text{ ns}$$



$$B_{0z} = 270\text{ G}$$

data acquired: 12 planes (20,000 spatial locations)



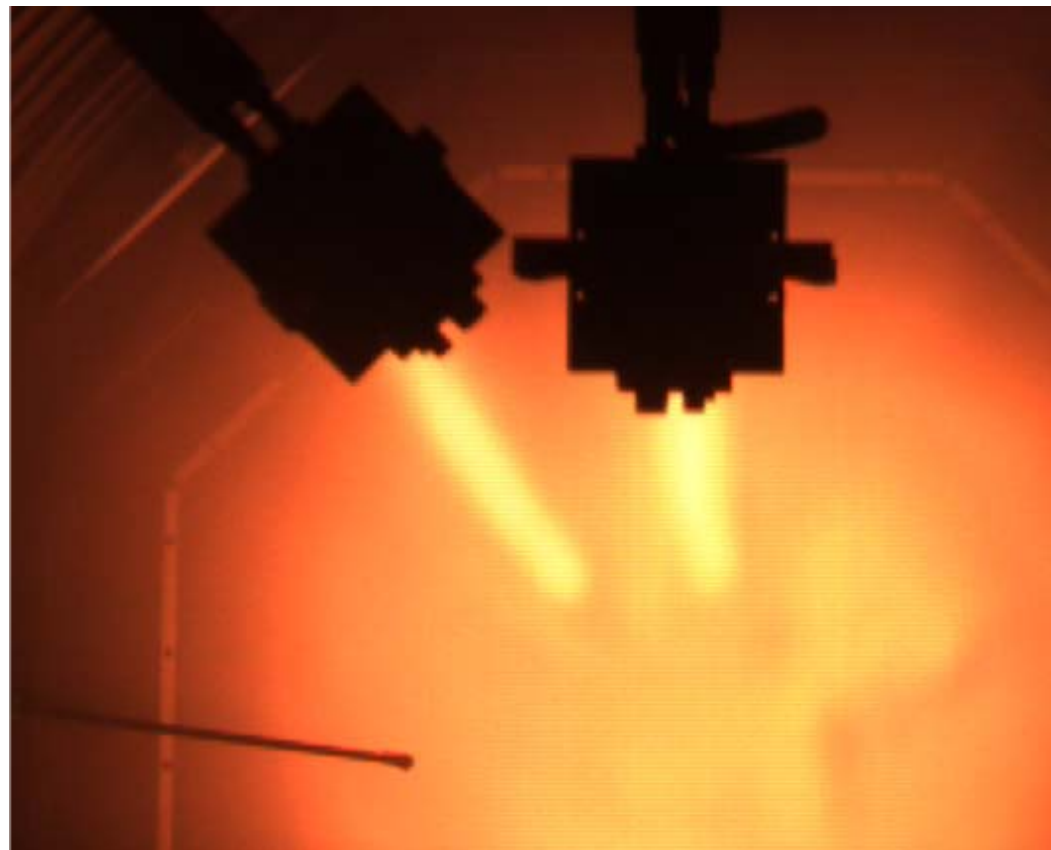
Discharge
 $V_D = 100V$

background plasma

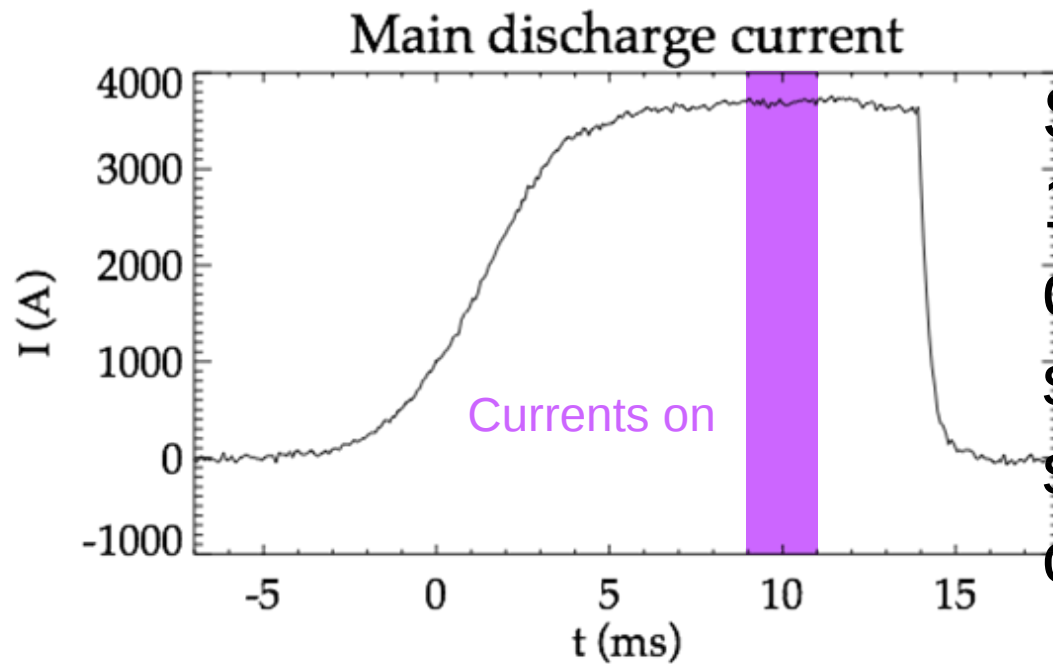
$$\gamma = \frac{c}{\omega_{pe}} = 3mm \quad S = \frac{\mu_0 V_A L}{\eta} \cong 3 \times 10^3 - 6 \times 10^4$$

$$\gamma_I = \frac{c}{\omega_{pi}} = 28cm \quad R_{ci} = 7.5mm$$

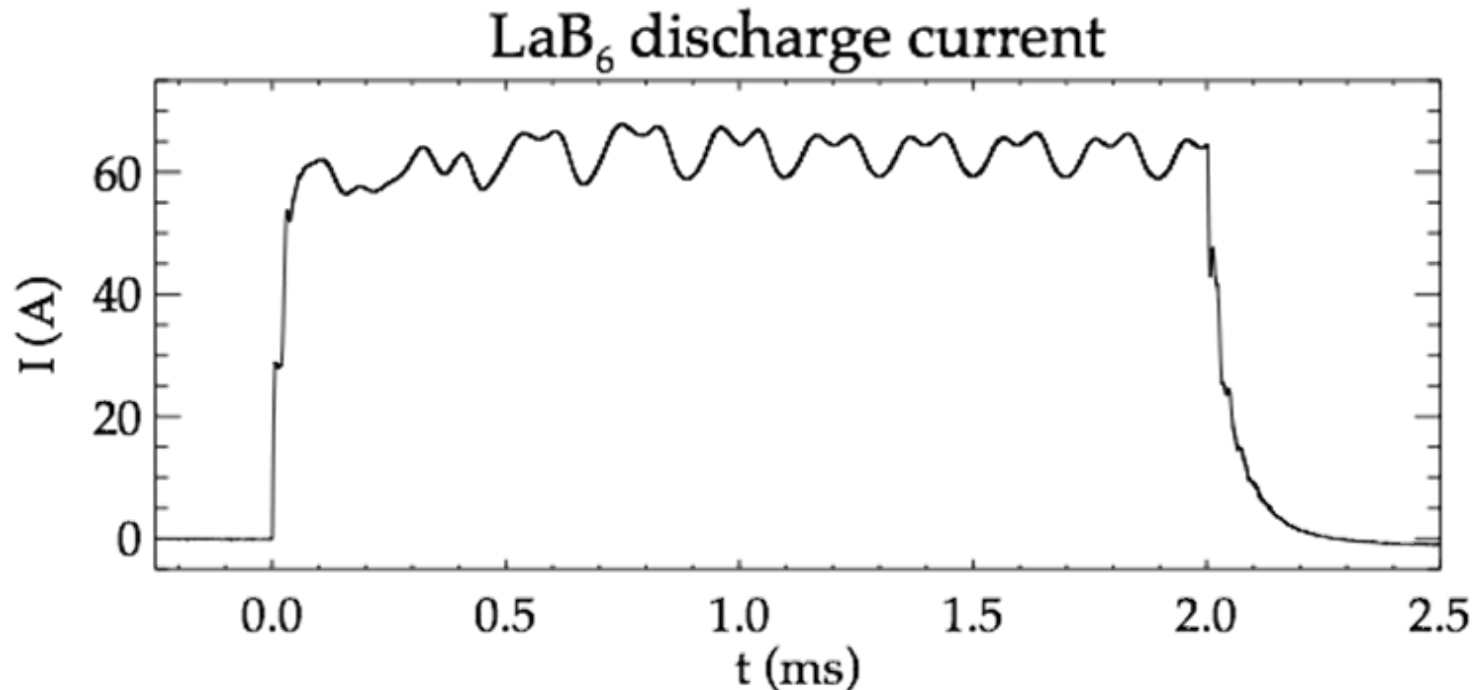
$$Q = \frac{2\pi a B_z}{LB_\theta} \cong 0.7 \quad B_z = 270G \quad B_\theta ; 6G$$



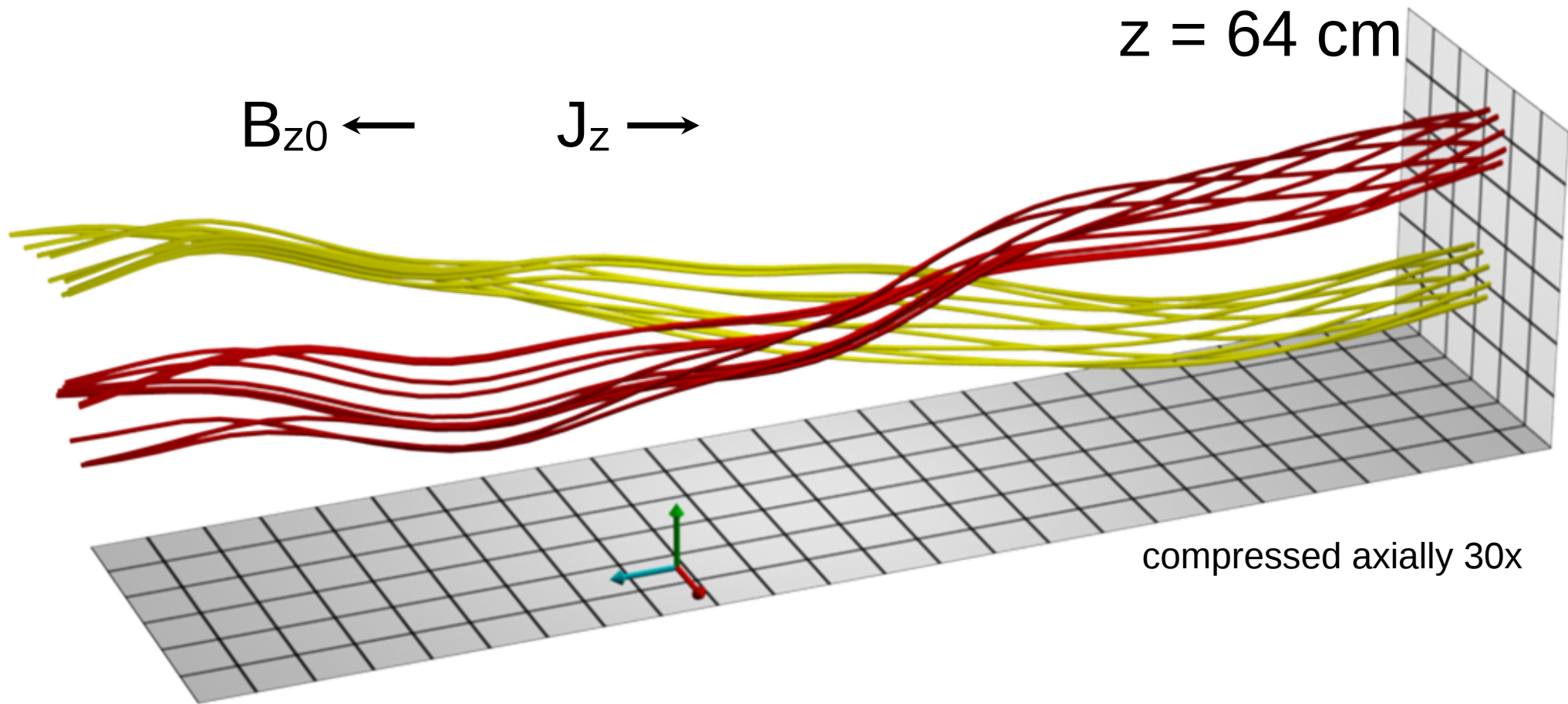
Discharge currents



Small cathodes are biased to 100V for 2 ms during the main discharge. After 300 μ s ($\sim 3\tau_A$), spontaneous oscillations are seen in the LaB_6 discharge current.

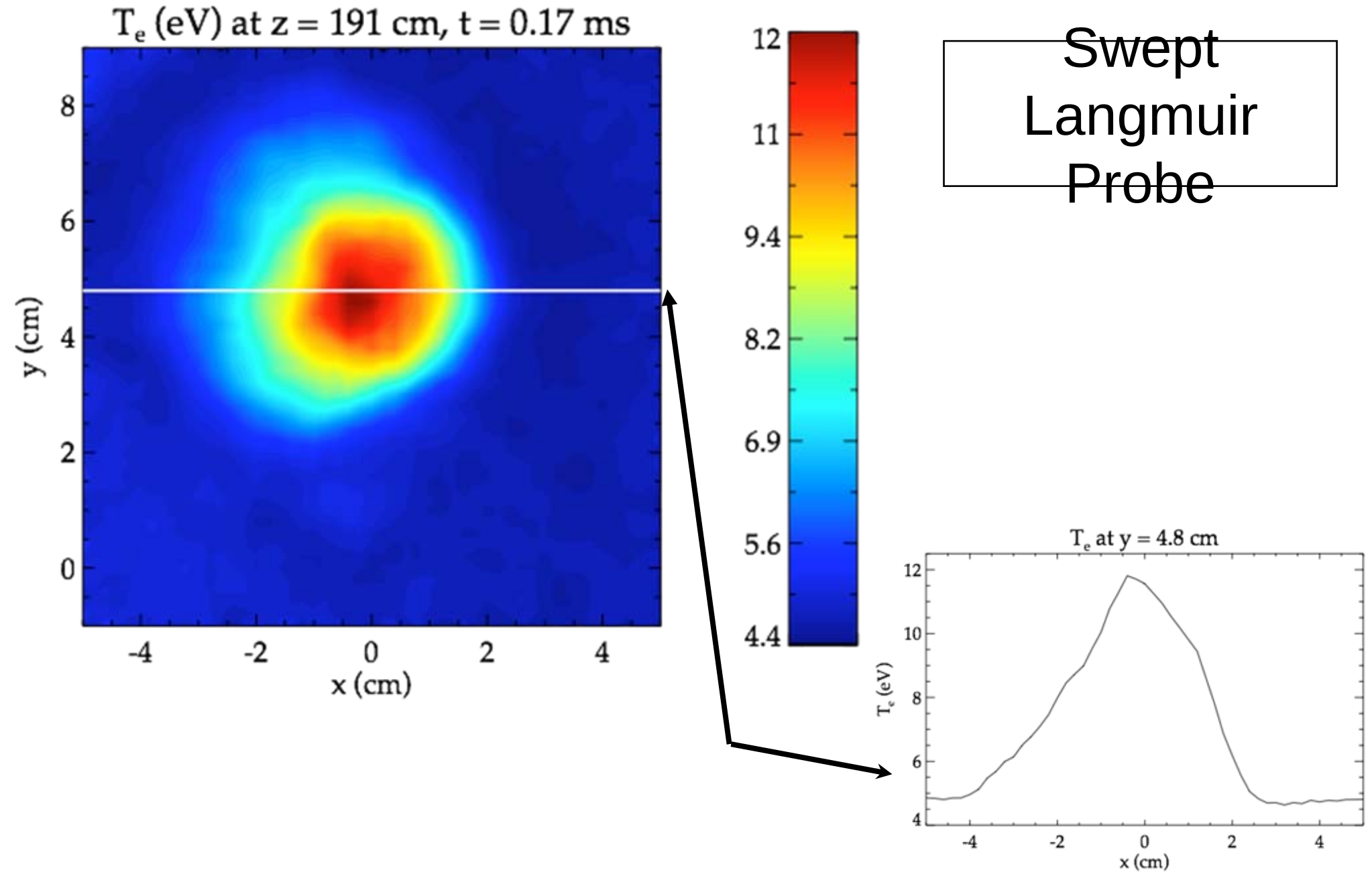


Magnetic field has flux rope structure



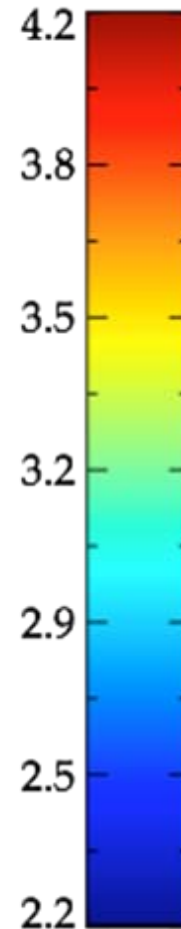
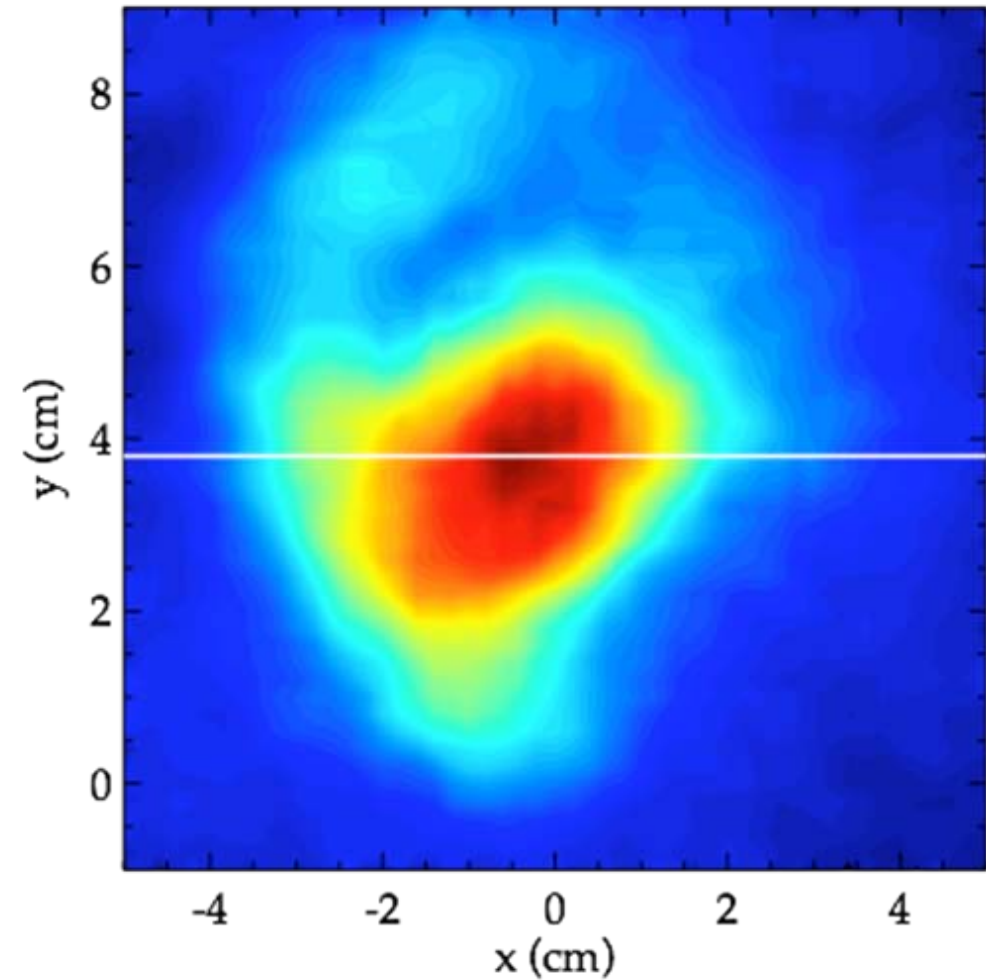
- ▶ Flux tube cross section is elliptical at the far end.
- ▶ Twist $\sim \pi - 3\pi/2$, writhe $\sim \pi$.

Electron temperature profile of a single channel



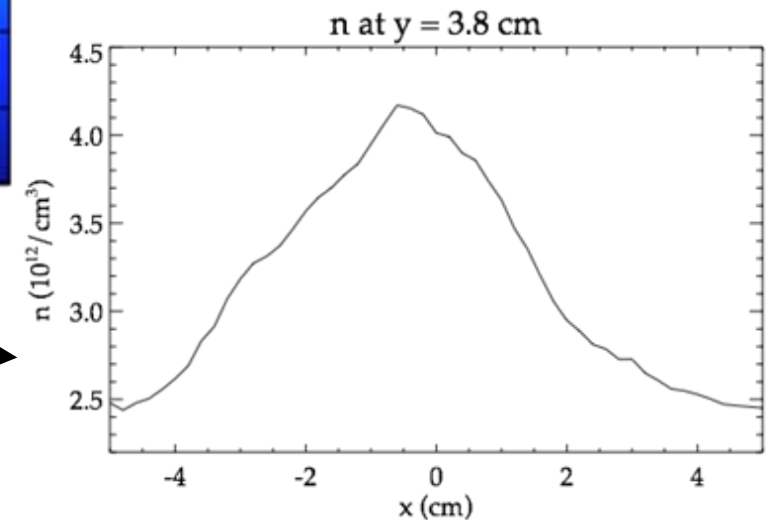
Density profile of a single current channel

n (10^{12} cm $^{-3}$) at $z = 191$ cm, $t = 0.17$ ms



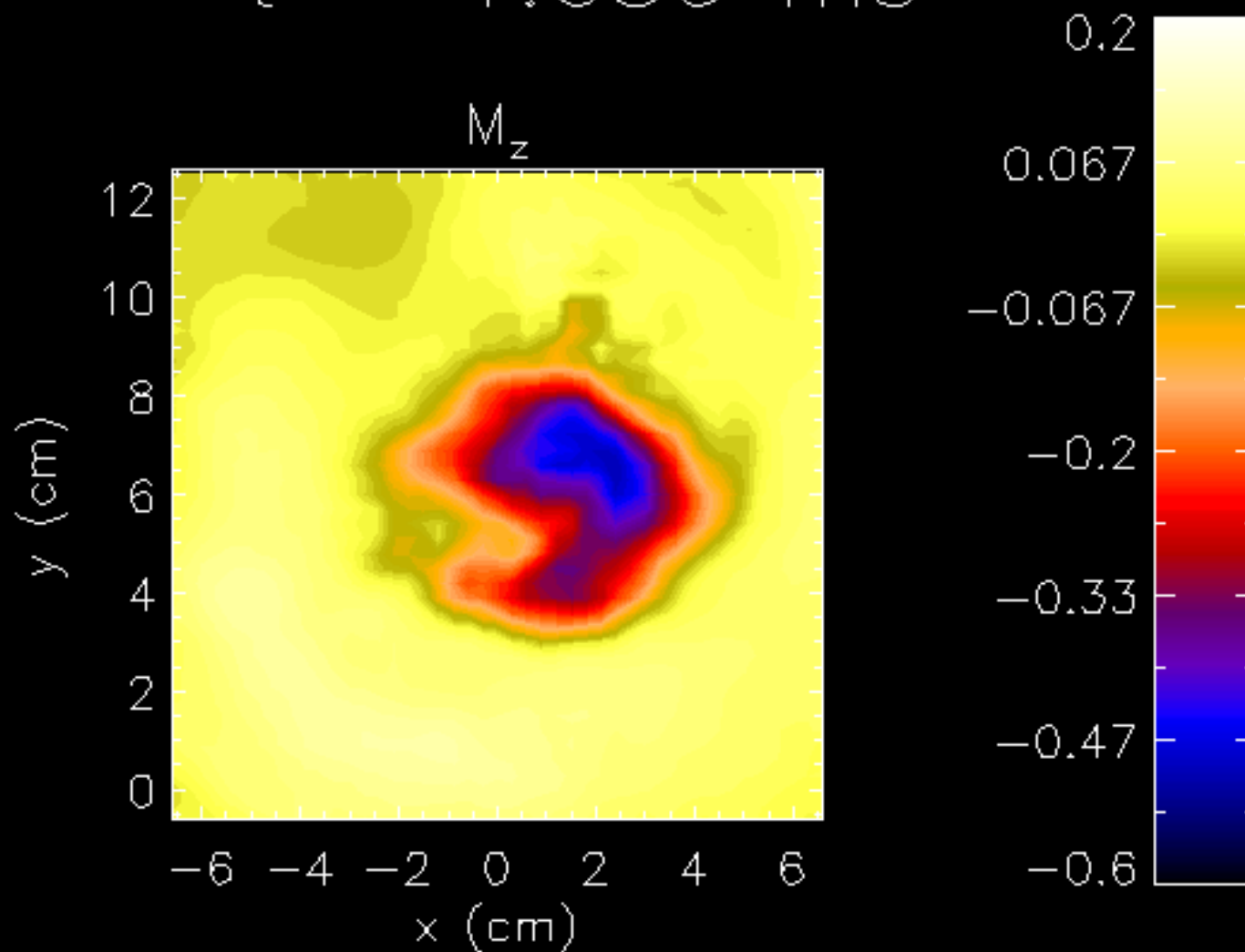
$n \times 10^{12}$ (cm $^{-3}$)

Swept
Langmuir
Probe

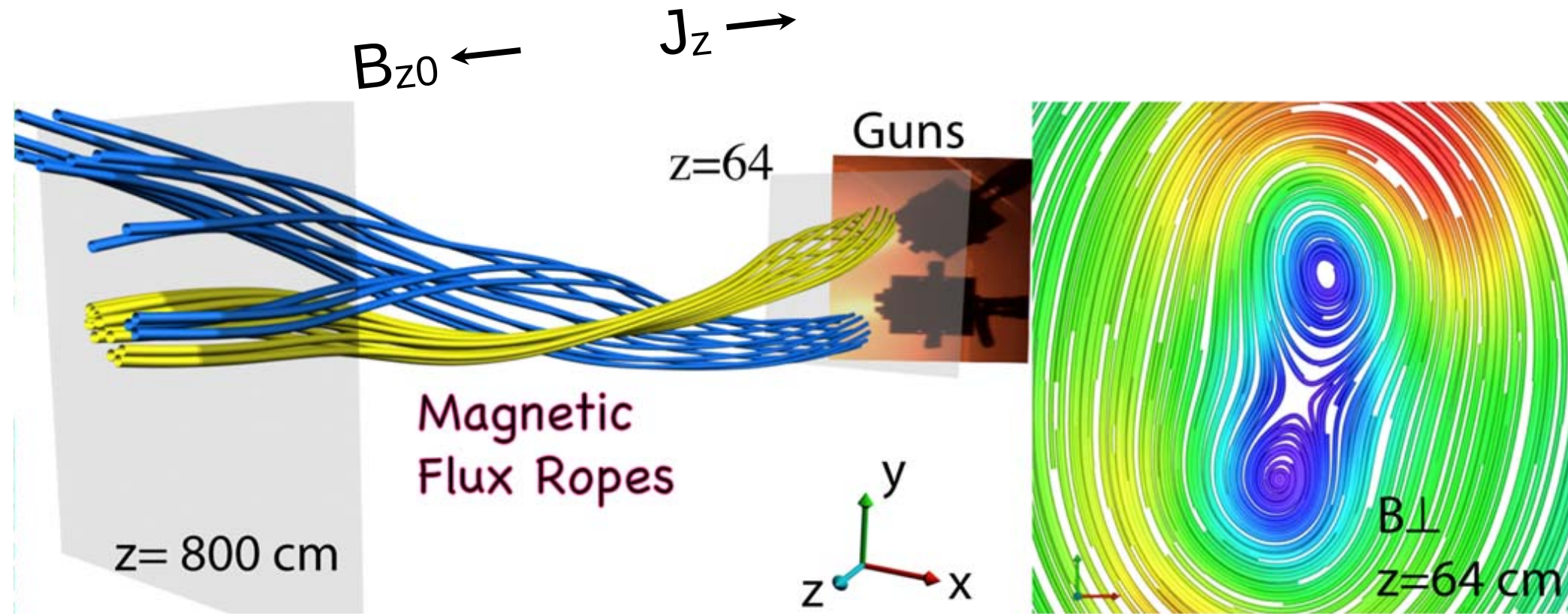


Plasma Flow

$t = 1.636 \text{ ms}$

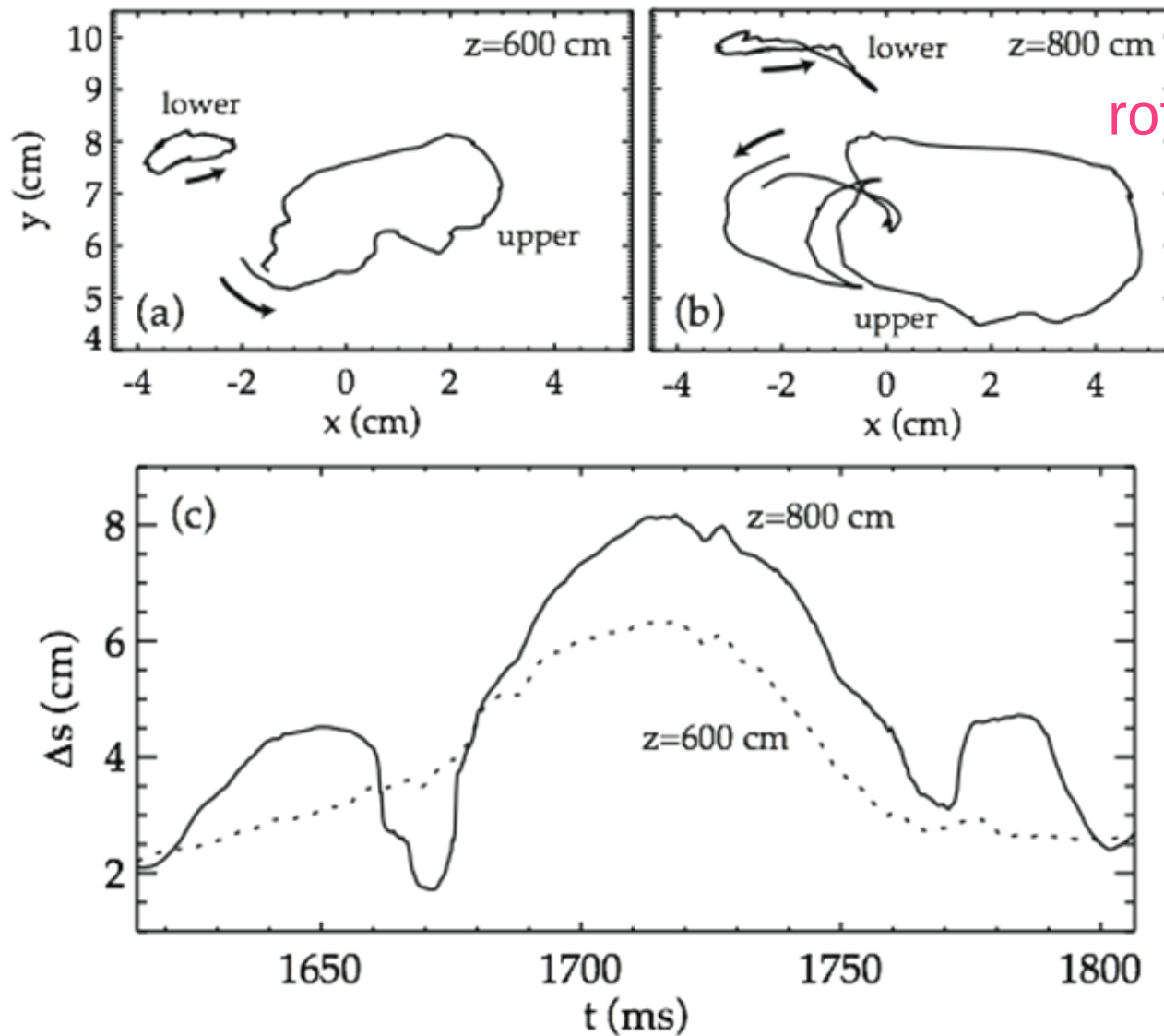


magnetic field is measured at 20,000 locations



- ▶ Flux tube cross section is elliptical at the far end.
- ▶ Twist $\sim \pi - 3\pi/2$, writhe $\sim \pi$.

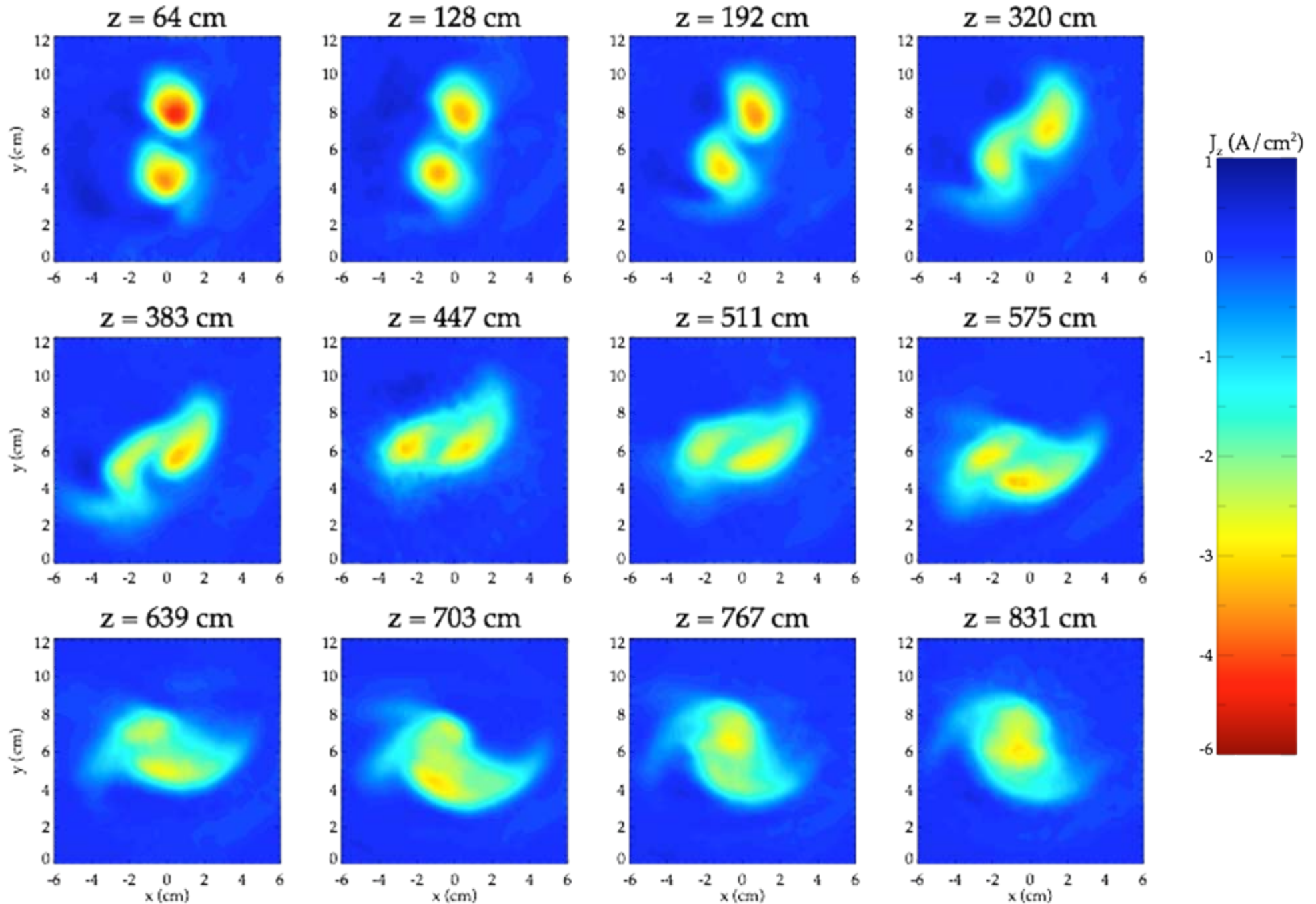
Hodogram of central field line in flux tubes



rotation due to parallel
flow?

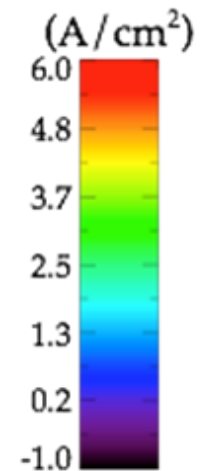
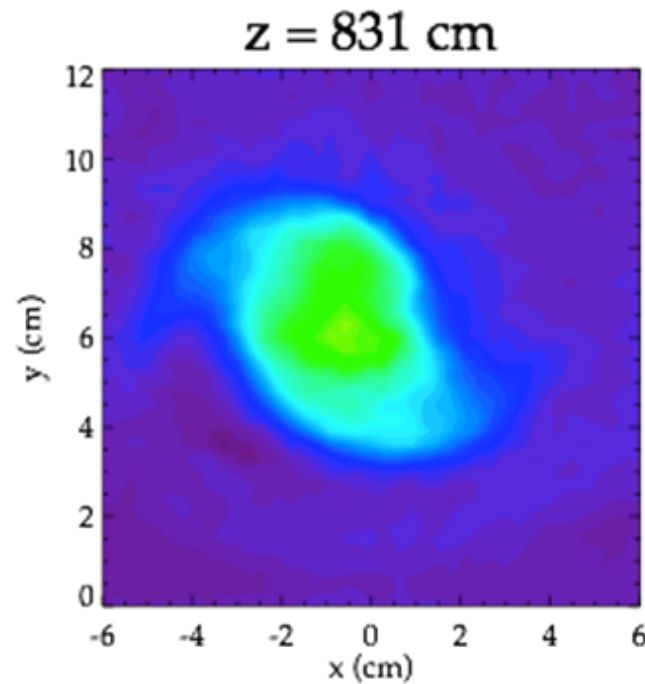
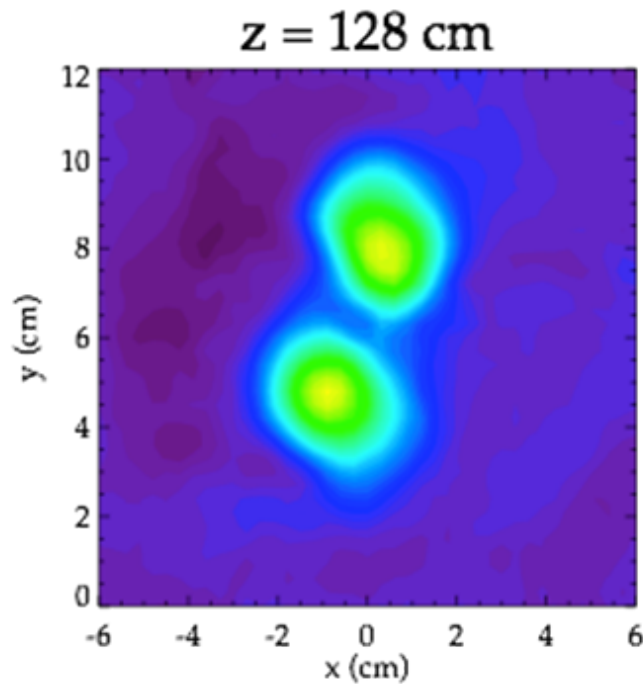
Distance between field lines (upper/lower) at two axial
positions

J_z slices at $t=190$ $\odot$ s (early in time)

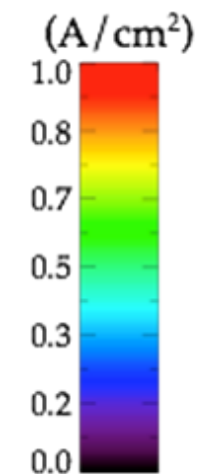
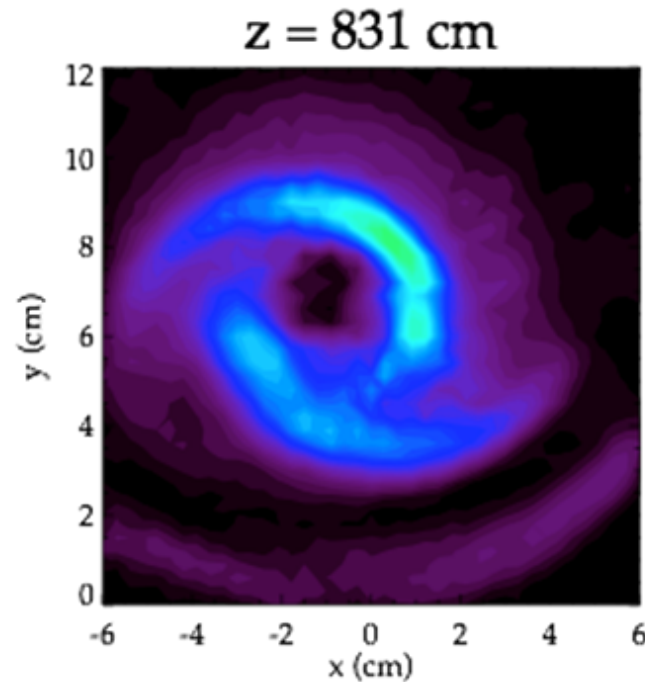
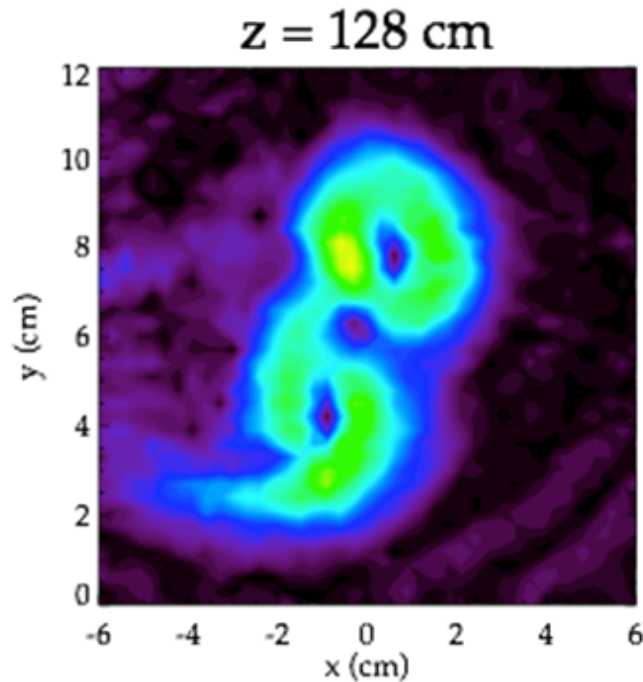


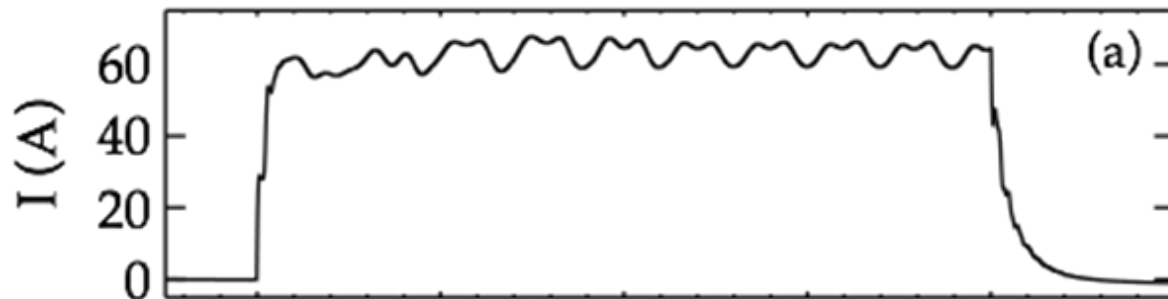
Parallel and Perpendicular Current

J_z

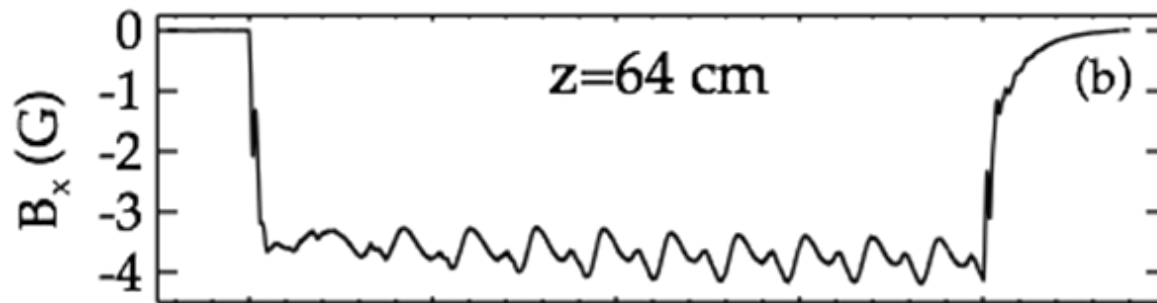


$J_{\perp}$

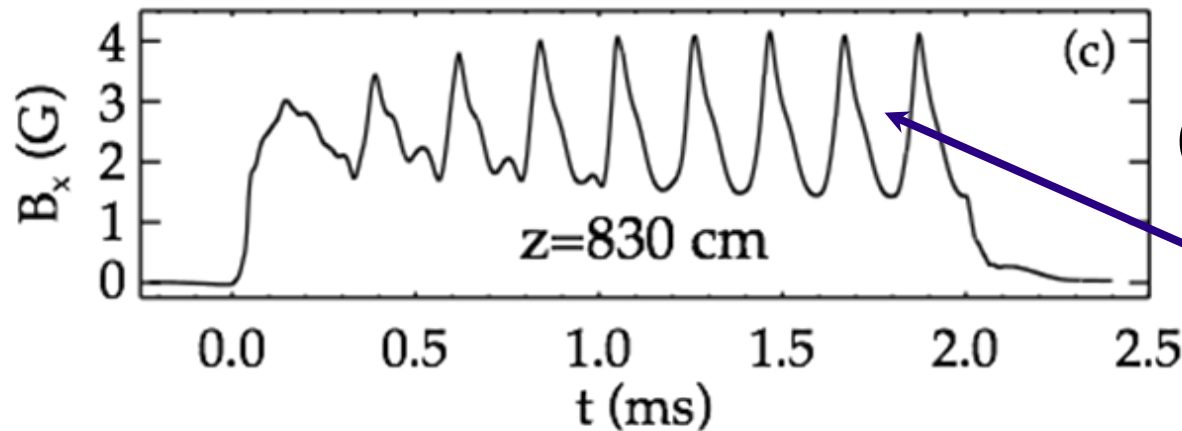




(a) emission current
from LaB6 cathode



(b) $B_x(t)$, (0.0, 2.7, $z=64$)



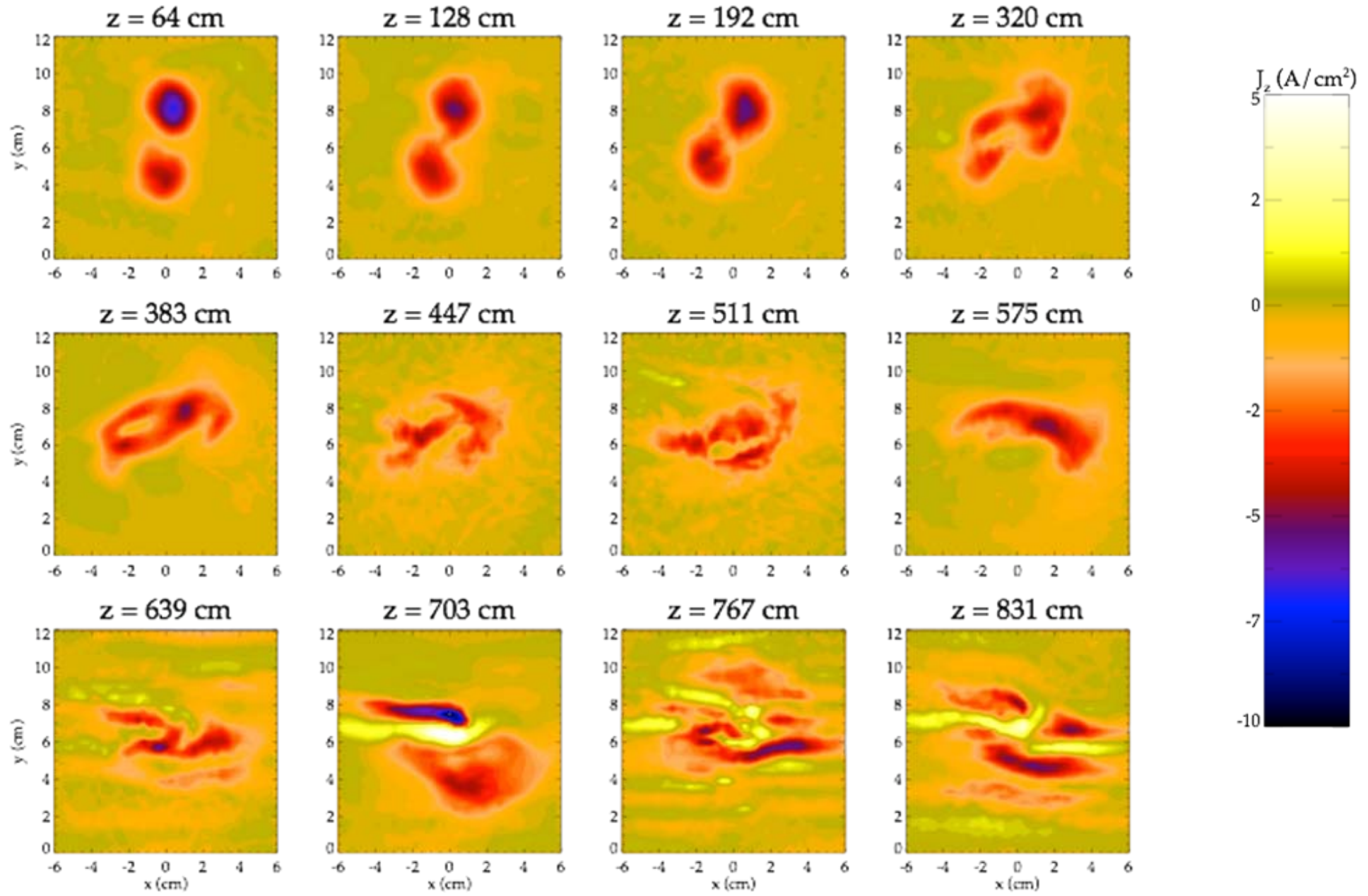
(c) $B_x(t)$, (-1.5, 9.9, $z=830$)

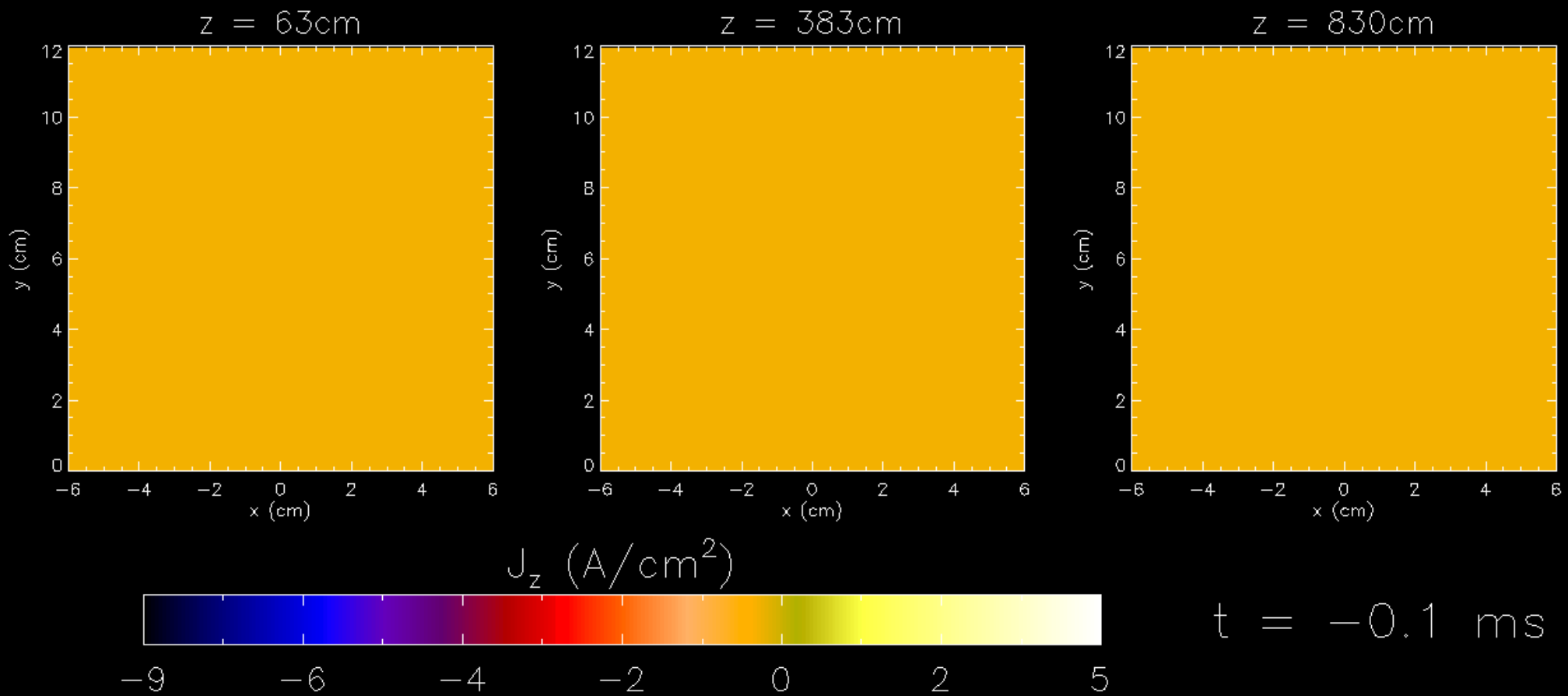
Period $t = 190 \text{ } \bigcirc \text{ s}$

(b) , (c) lie on same approximate flux
surface on the lower flux rope

Jz profiles during oscillations in discharge current

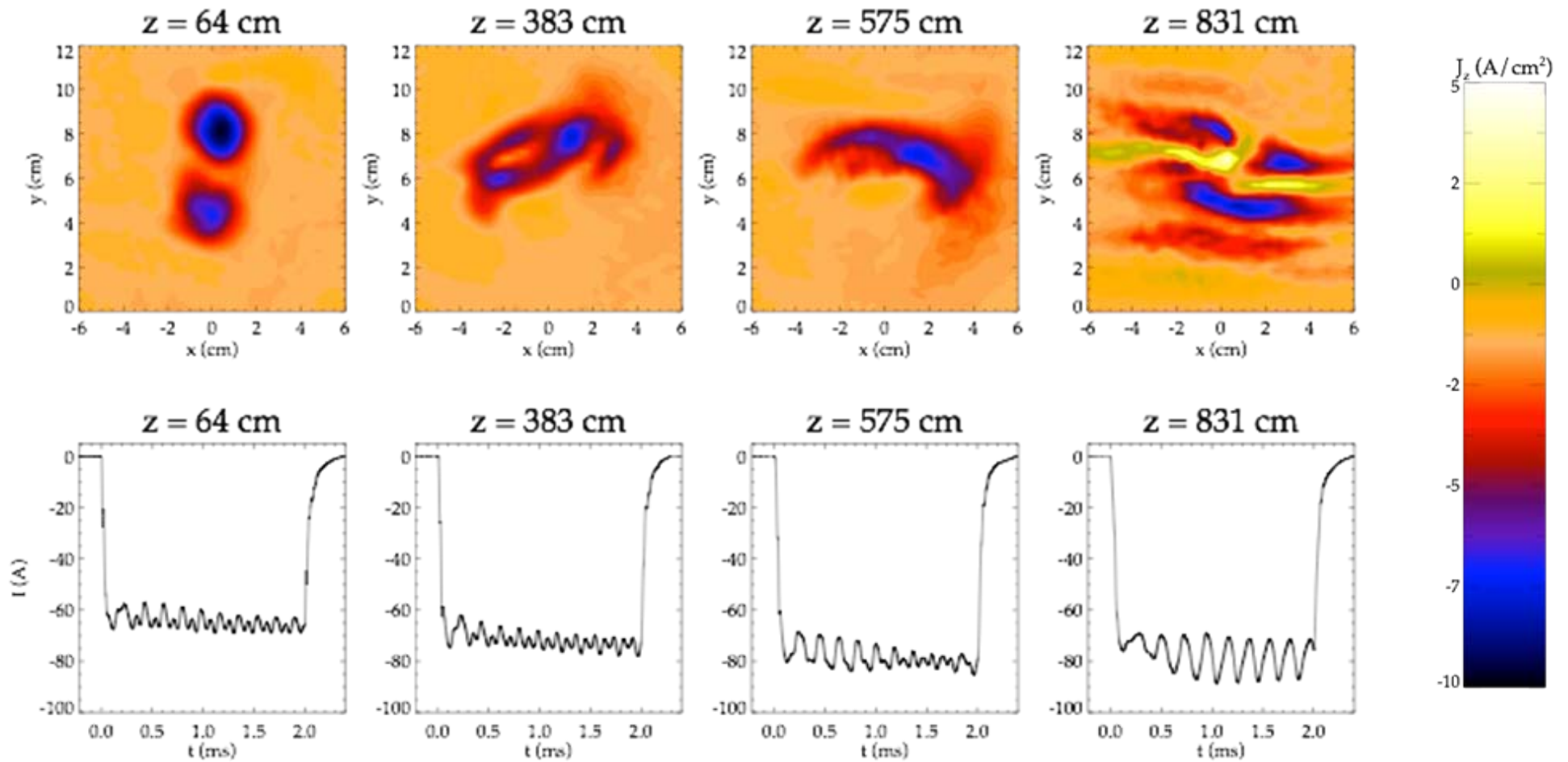
t=1.7ms



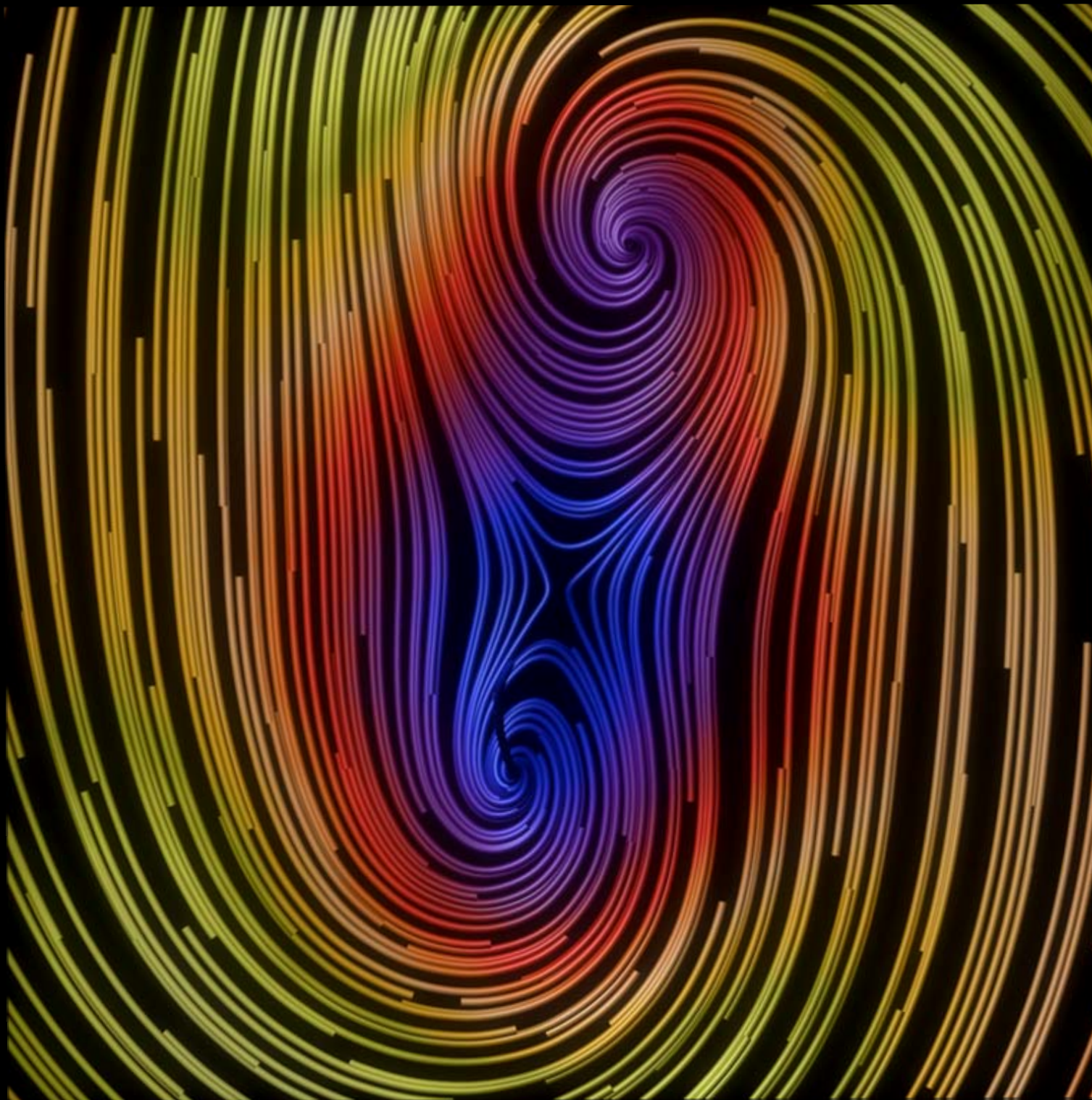


current density

Integrated current density



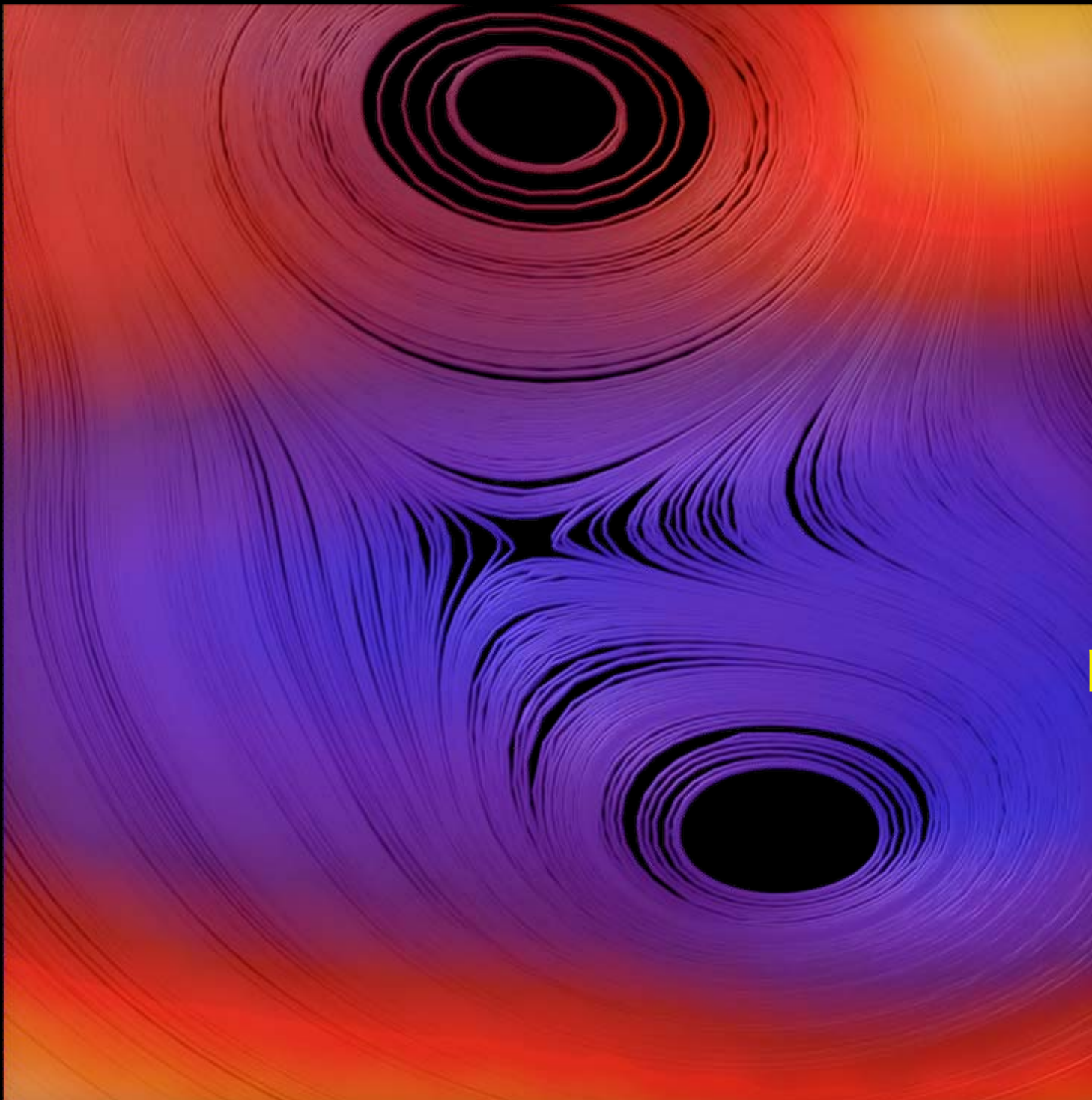
Bottom traces show $\int J_z dz$ at various z positions.



$B_x - B_y$

◆ = 5 ms

$\Omega z = 64$ cm

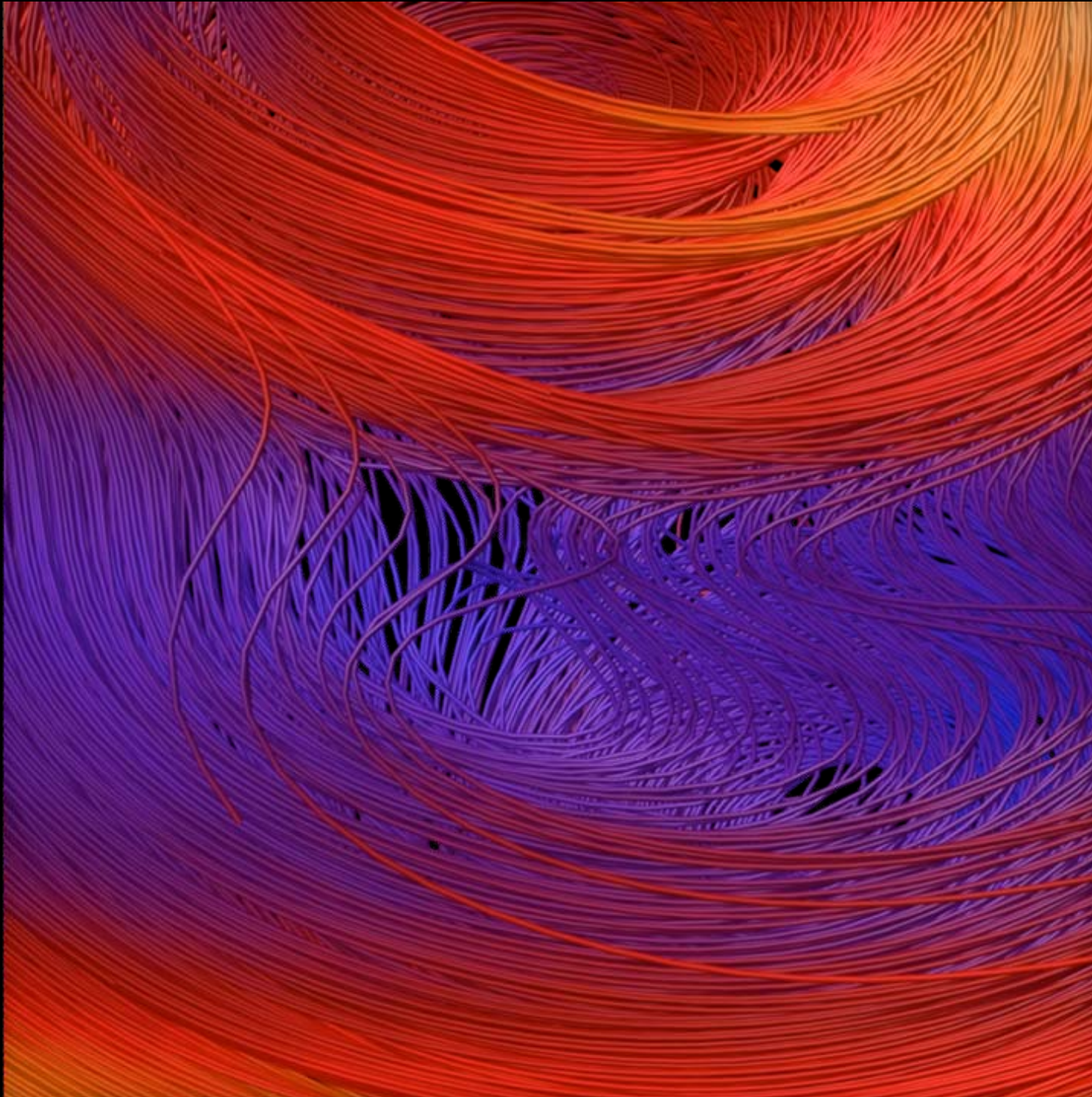


$B_{\perp}$

$z=6.6$ m

$t=88.8$ μ s

no guide field

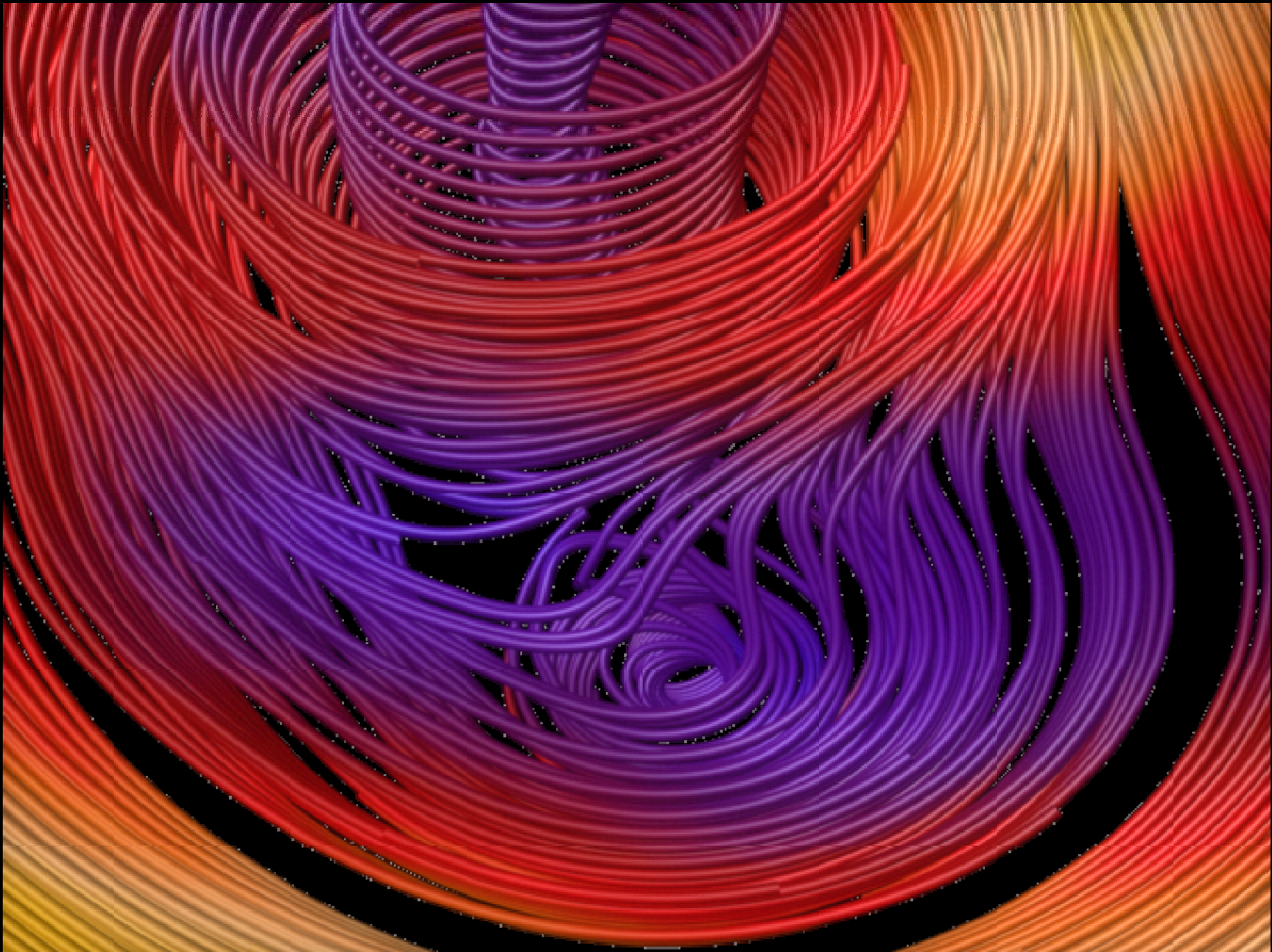


B

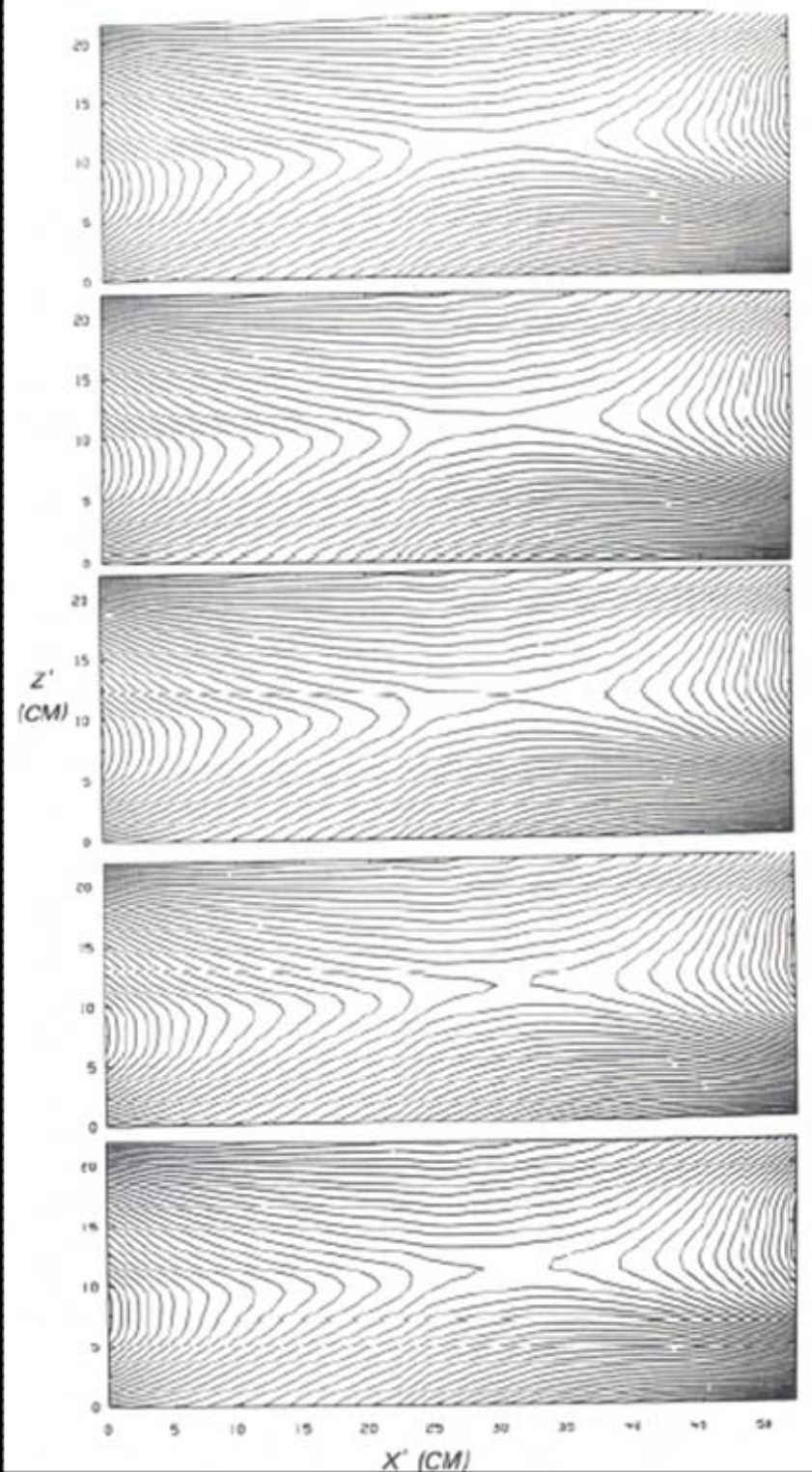
$z=6.6$ m

$t=88.8$ μ s

no guide field,
 B_{zropes}



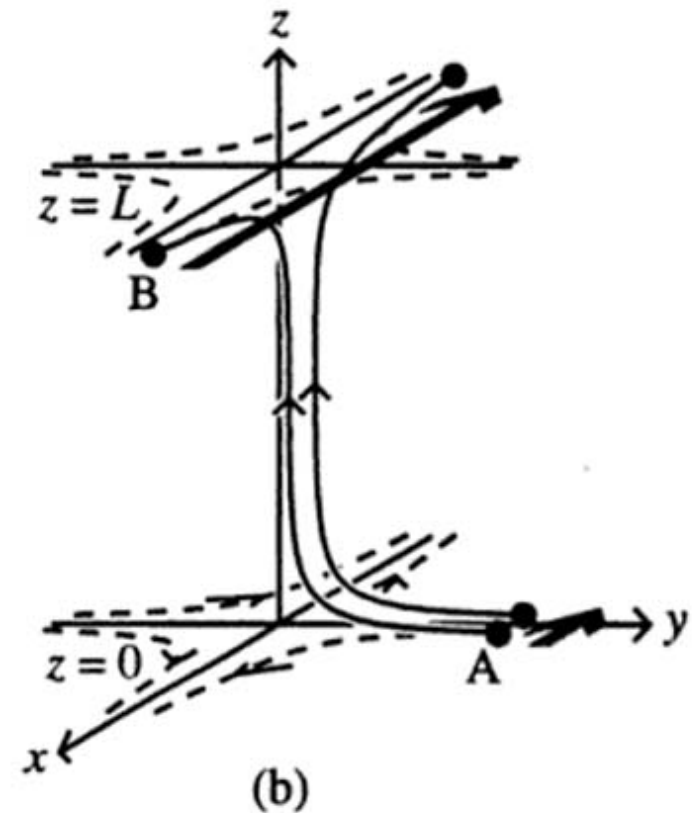
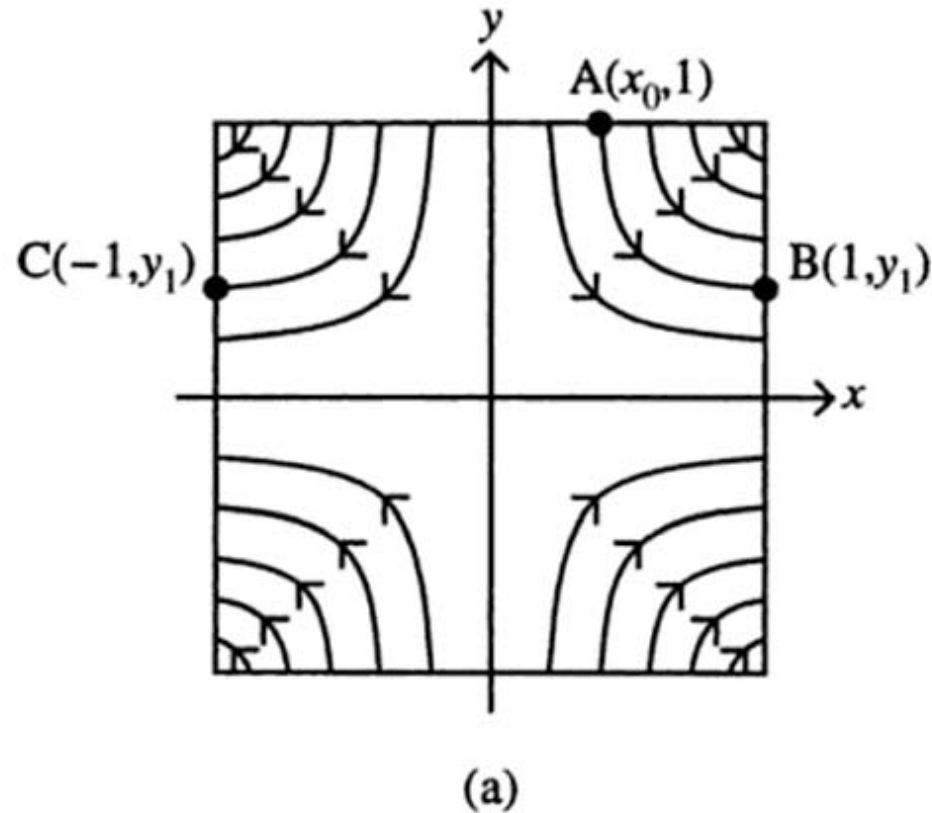
sound generated by R. Fanelli from Alfvén wave - waveform



2D reconnection

Gekelman, Stenzel, Wild, Phys. Scripta, T2/2 (1982)

Simple sheared X-point model



Small footpoint motions at point A would create a drastic shift at point B. In some cases point B can shift discontinuously [Priest and Forbes, 2000, Demoulin, 2006]

Definition of a quasi-separatrix layer

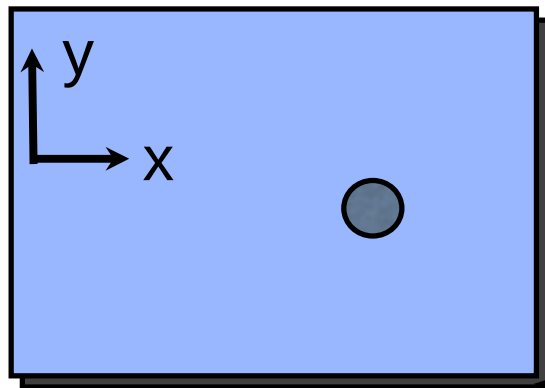
$$Q = \frac{N^2}{\left| \frac{B_z(z_0)}{B_z(z_1)} \right|}$$

$$N = \sqrt{\left(\frac{\partial X}{\partial x} \right)^2 + \left(\frac{\partial X}{\partial y} \right)^2 + \left(\frac{\partial Y}{\partial x} \right)^2 + \left(\frac{\partial Y}{\partial y} \right)^2}$$

in our experiment

$$Q \approx N^2$$

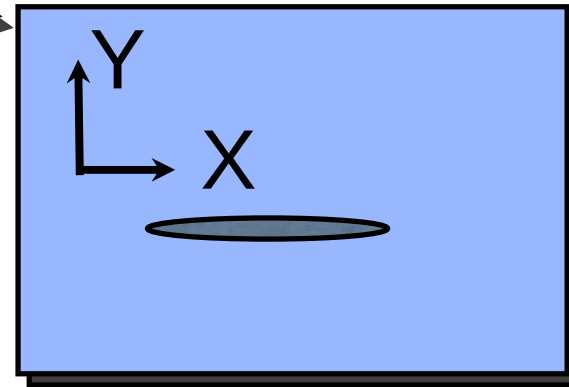
“slip squash factor”



$z = z_0$

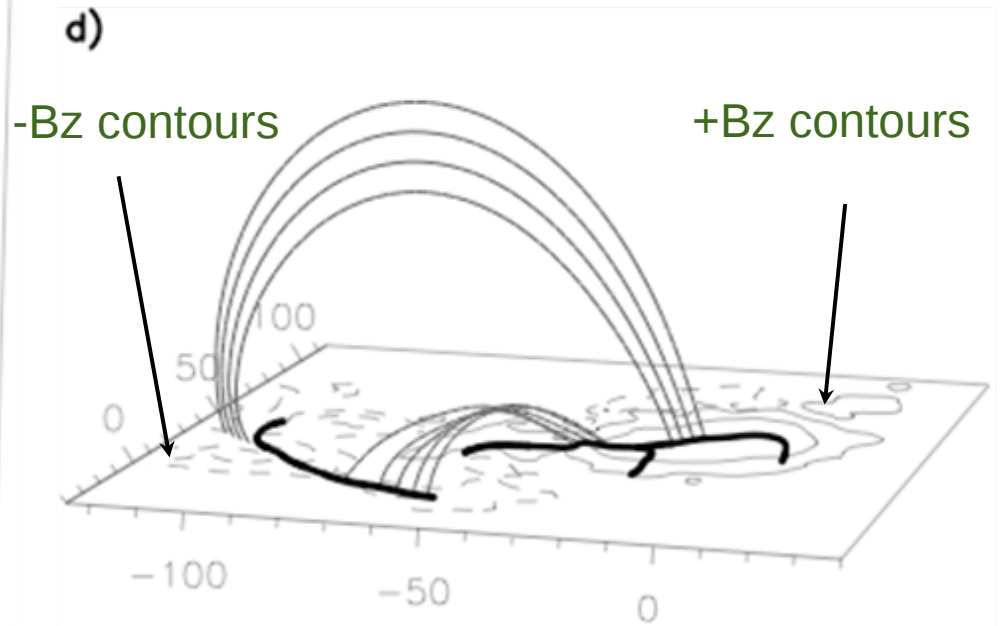
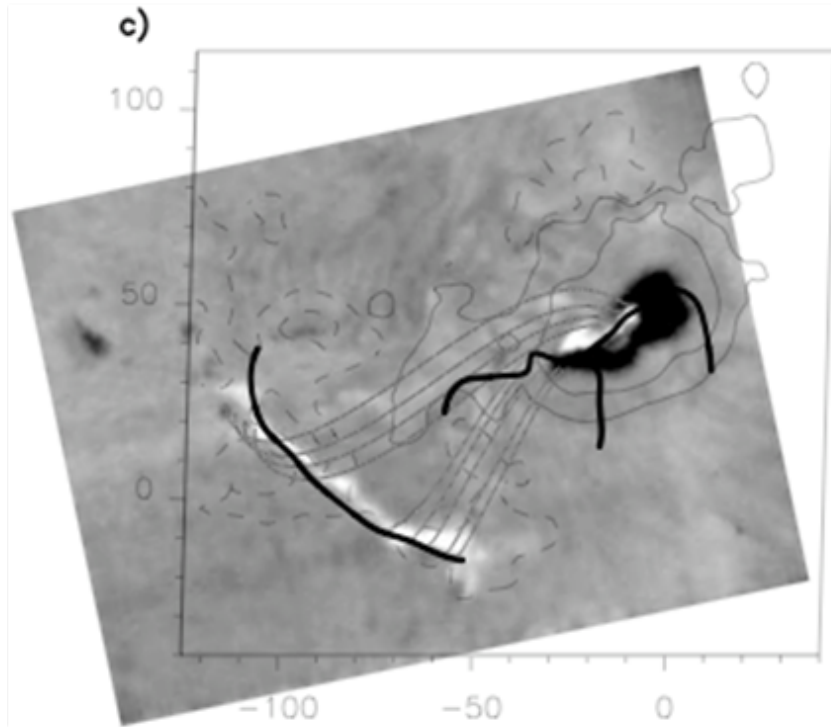
flux tube

$z = z_1$



flux tube squashed: aspect ratio = Q

QSLs in solar flare observations



computed field lines

QSLs at boundaries (thick lines) coincident with H_α brightening in solar flares. [Bagalá et al. A&A 2000]

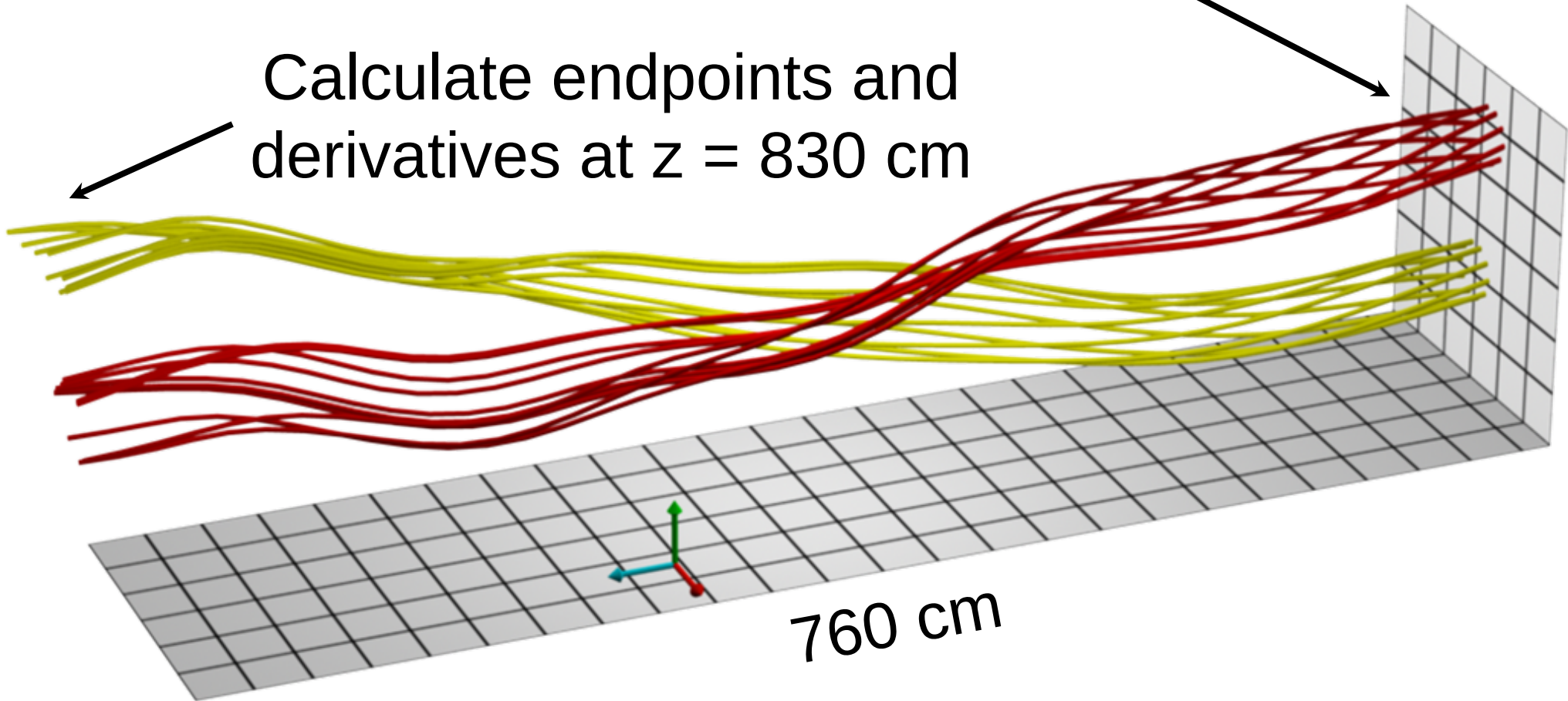
QSL calculation

QSL region if $N \gg 1$

in this experiment $N=2000$

Seed field lines
at $z = 64$ cm

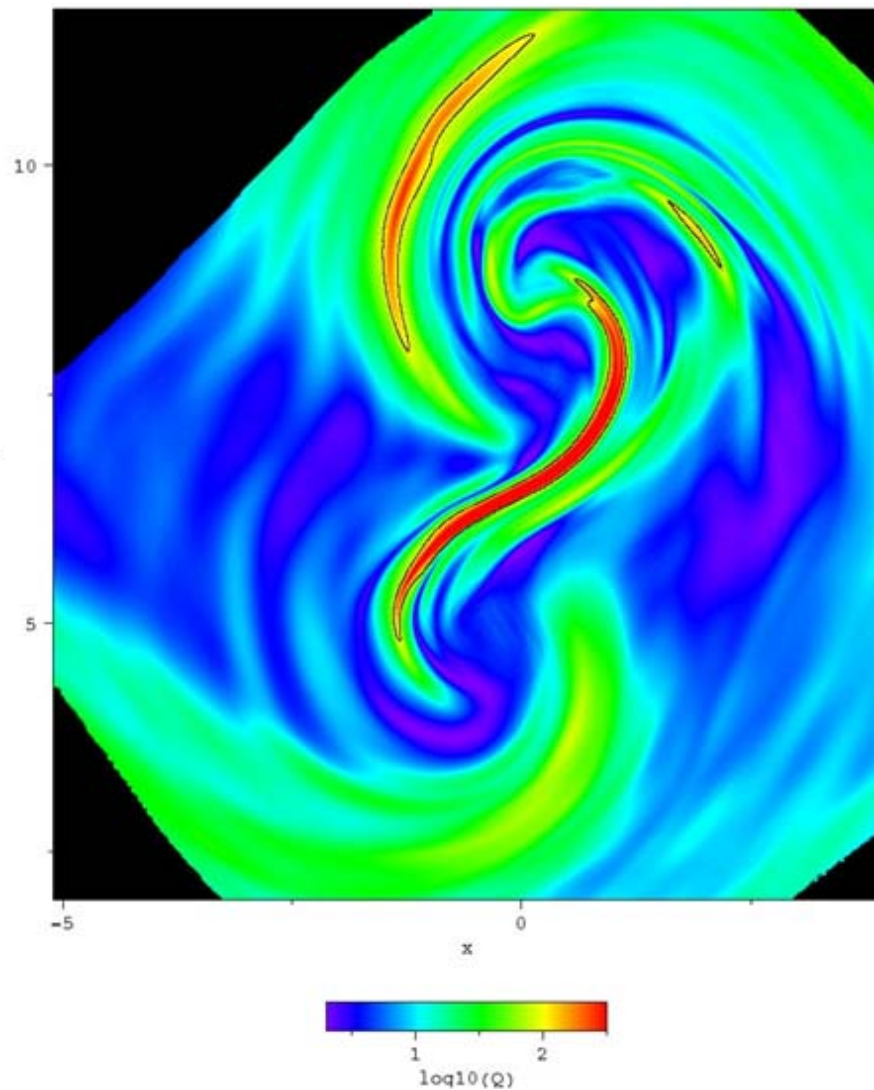
Calculate endpoints and
derivatives at $z = 830$ cm



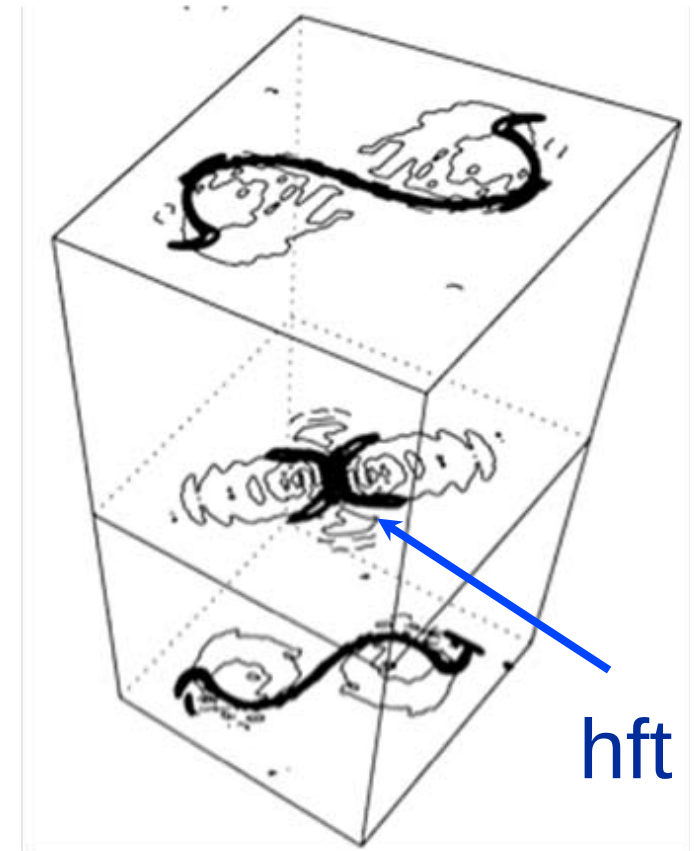
In the following slides, Q is shown during
the merging phases

QSL forms between flux ropes

Experiment

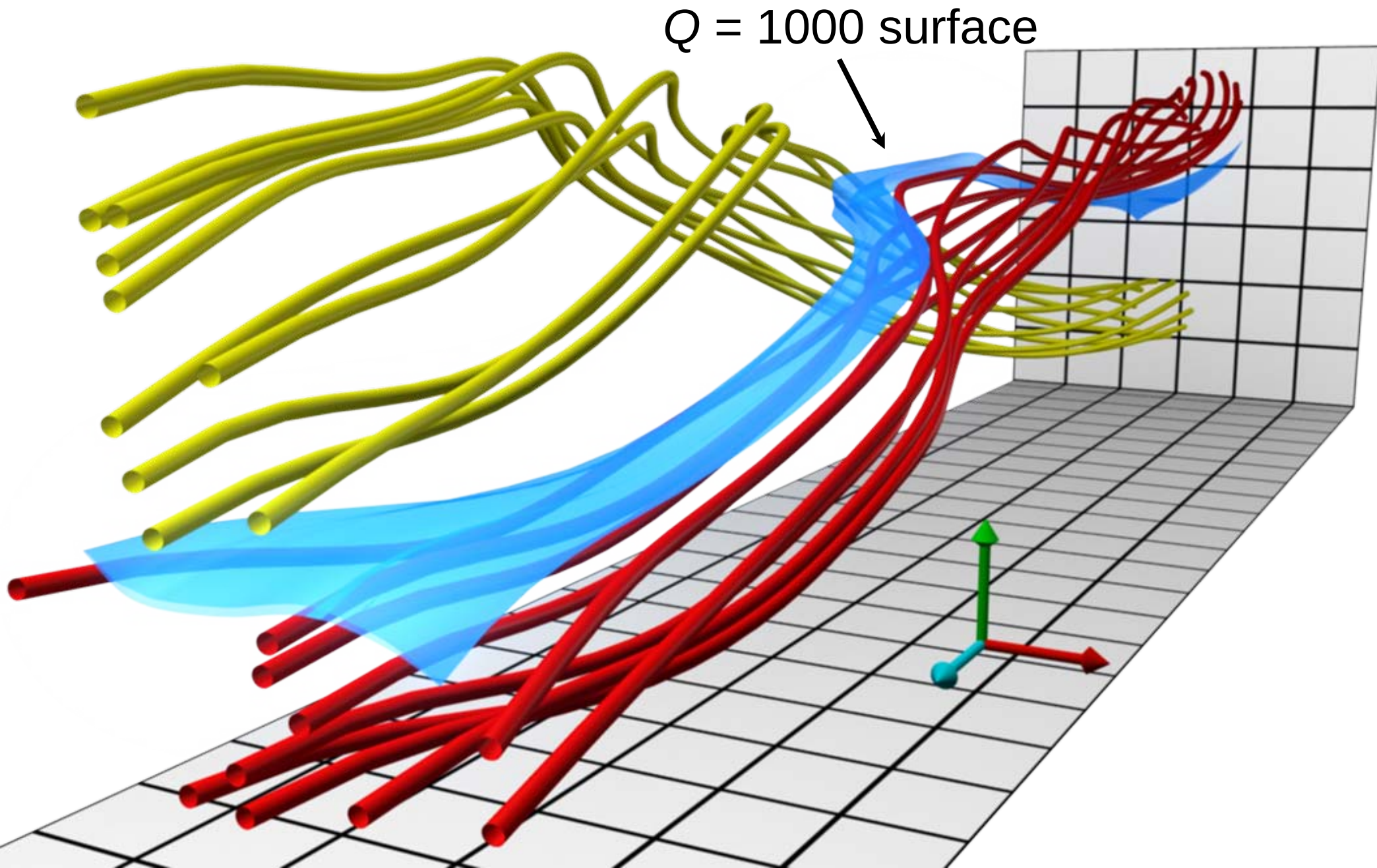


Q calculated between $z = 64$ cm and $z = 830$ cm planes at $t = 1.7$ ms.
 $J_z = -5.5, -3$ A/cm² contours overplotted.



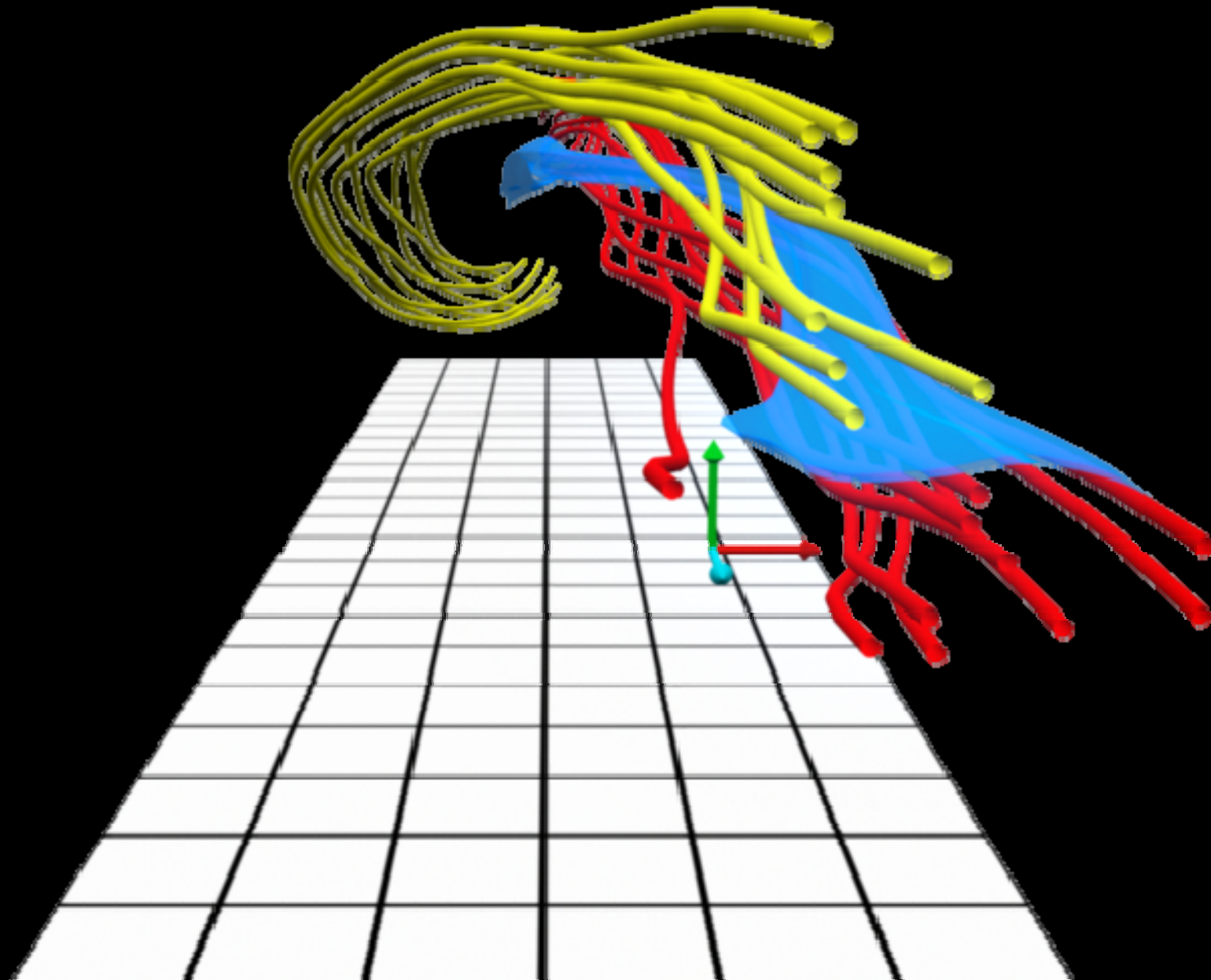
Numerical simulations of merging twisted flux tubes
[Milano, et al. ApJ 1999]

QSL forms between flux ropes

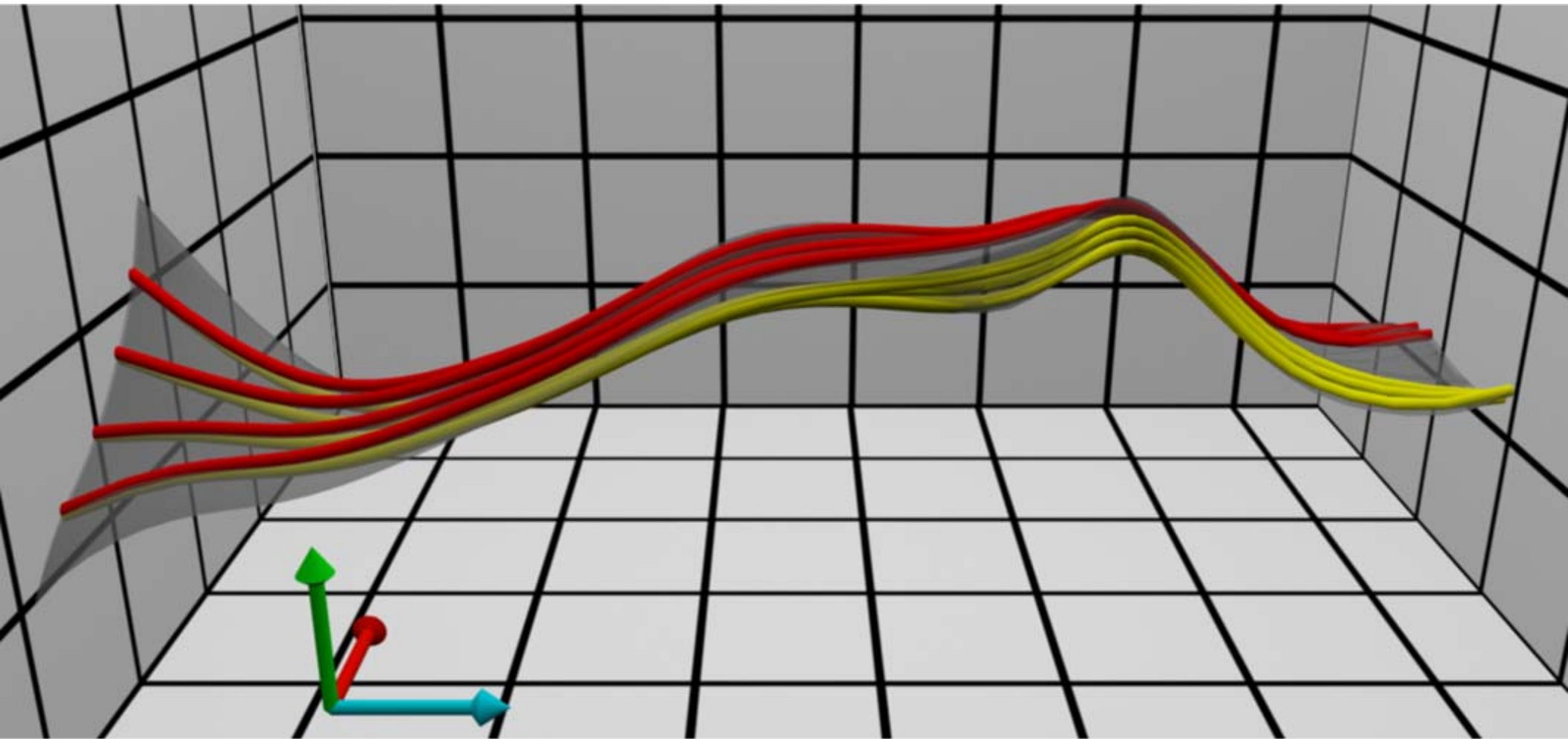


QSL forms between flux ropes

$Q = 1000$ surface



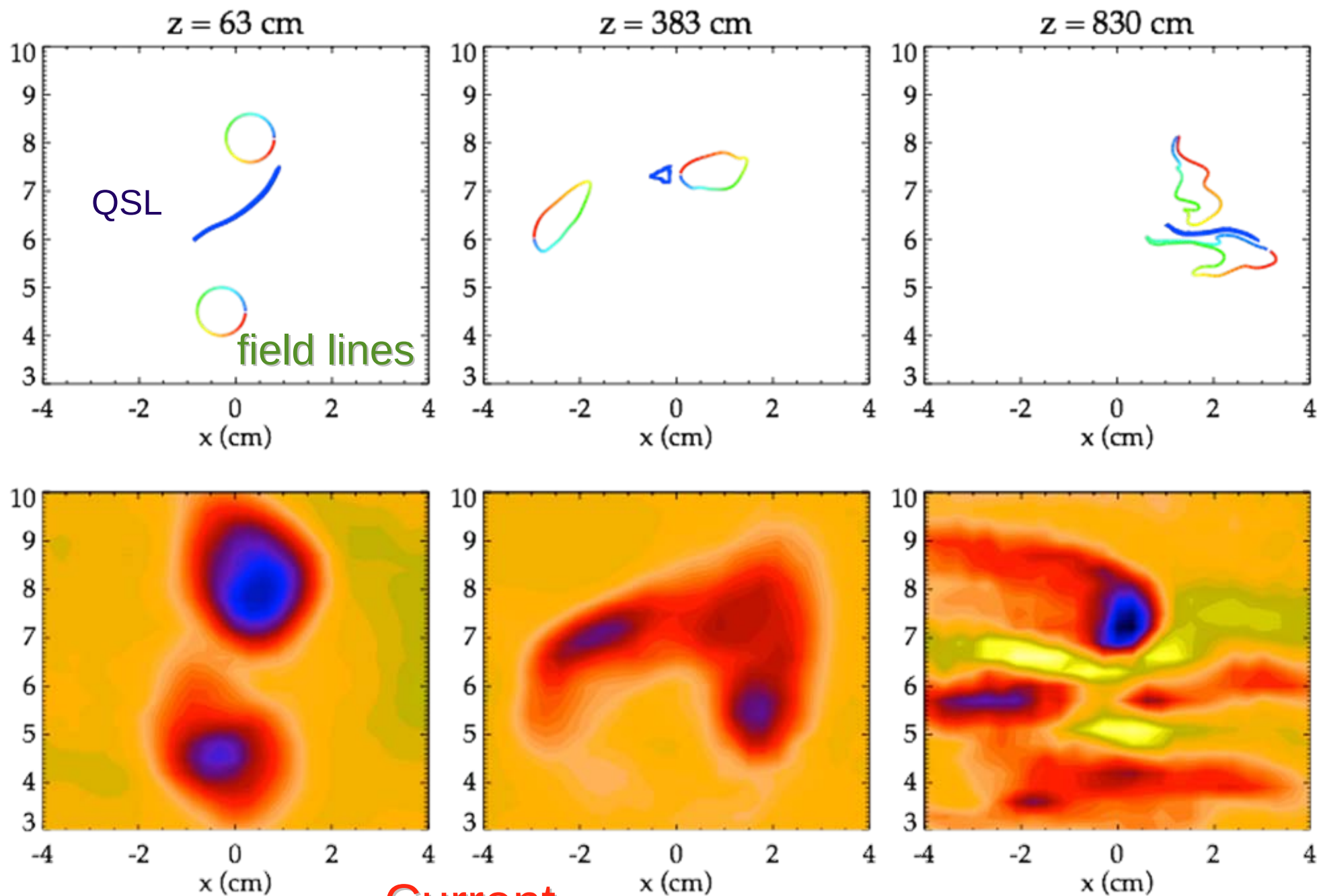
QSL has hyperbolic flux tube geometry



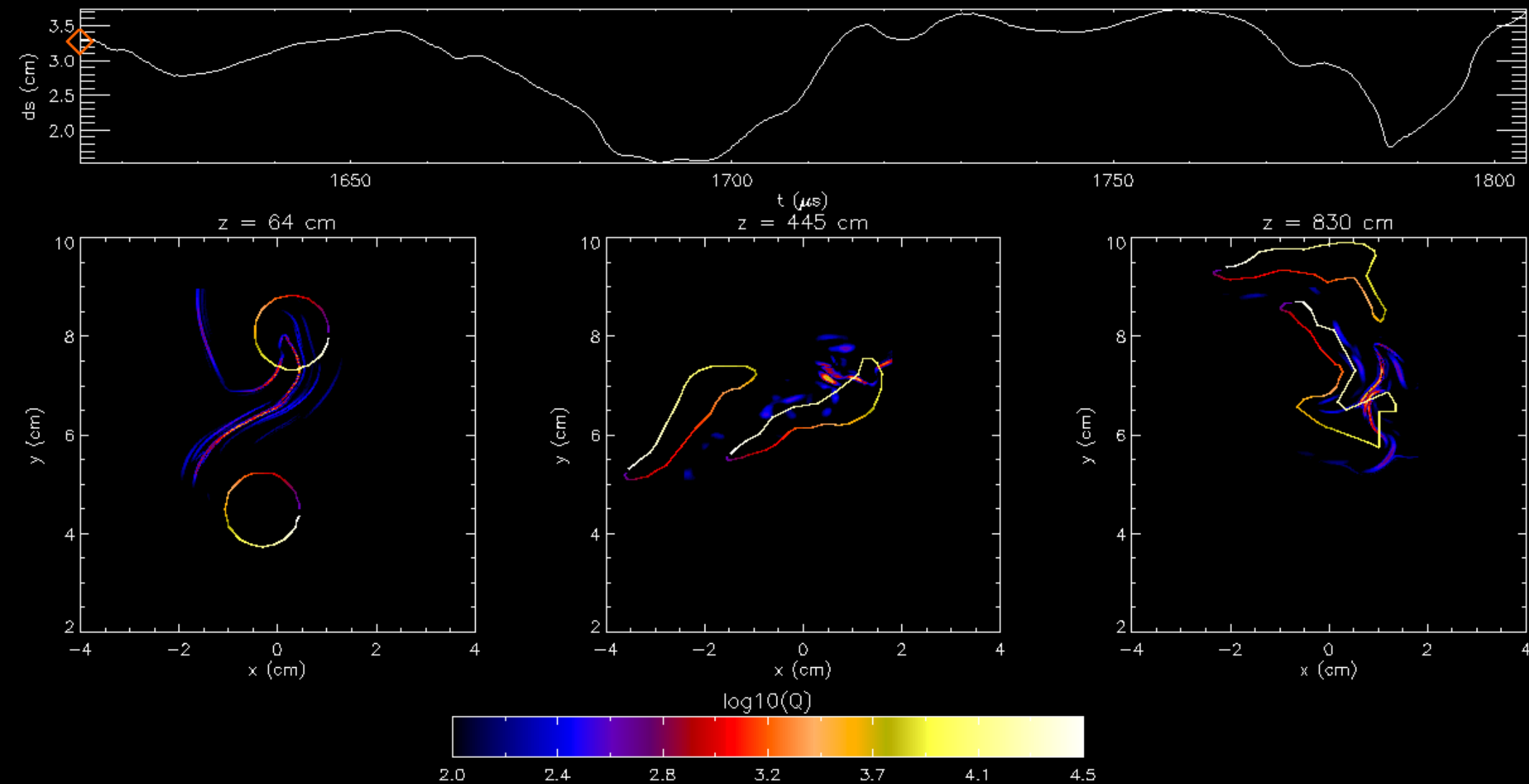
compressed 100x axially

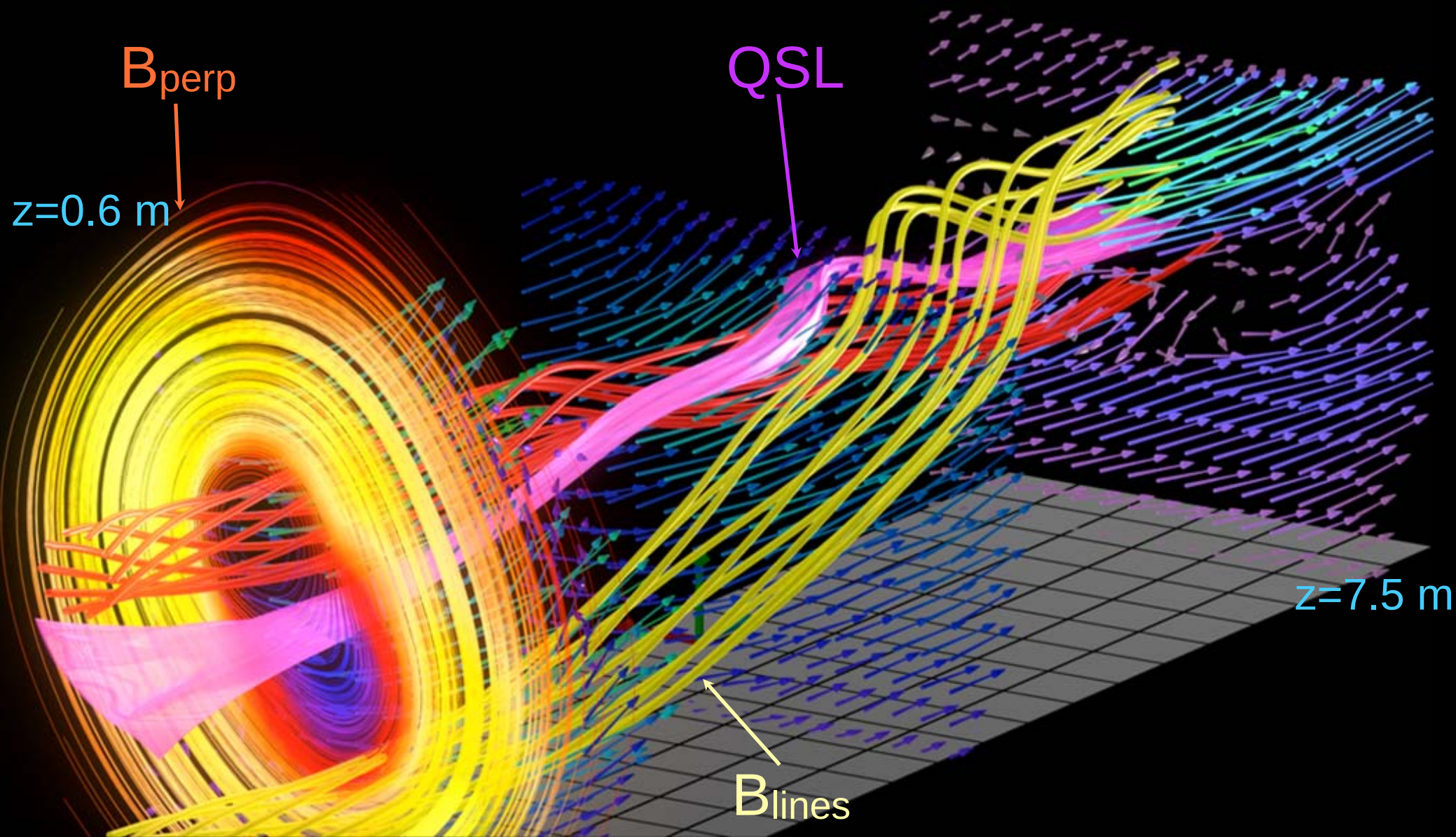
- ▶ Initial field line separation is ~ 0.05 cm, but diverges to ~ 3 cm.

Axial slices show HFT (hyperbolic flux tube) structure

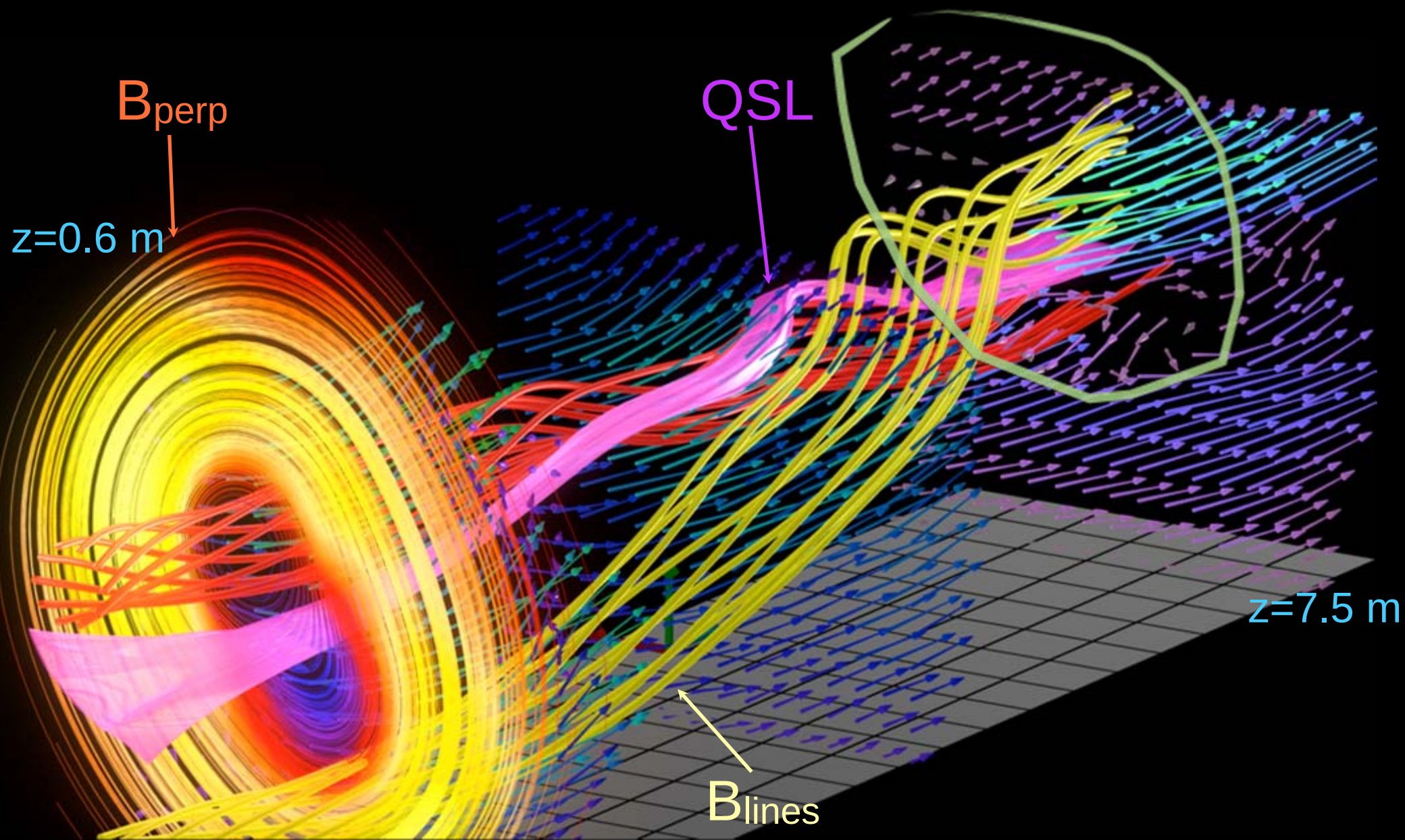


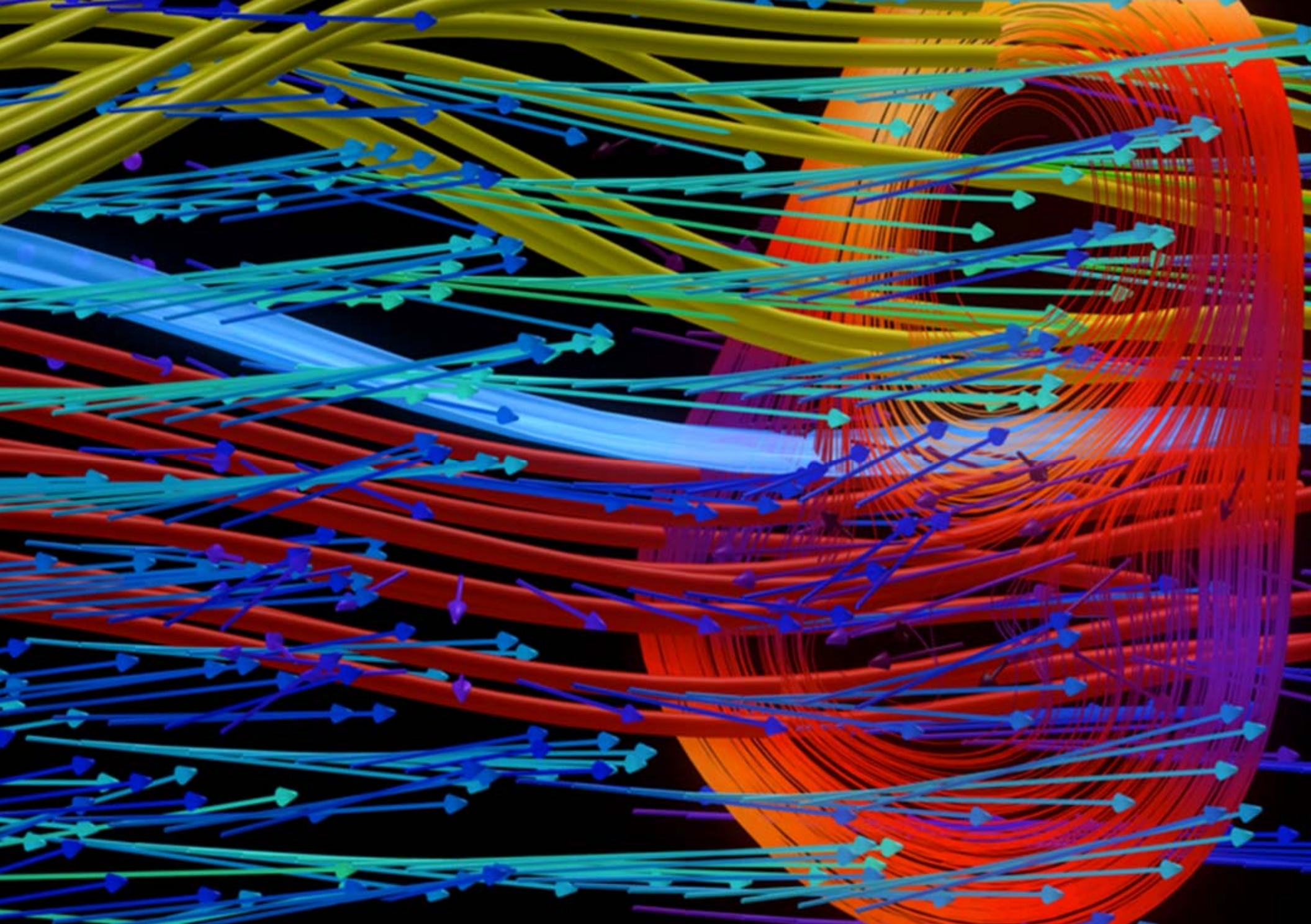
Current





$\Omega t = 1.69 \text{ ms}$



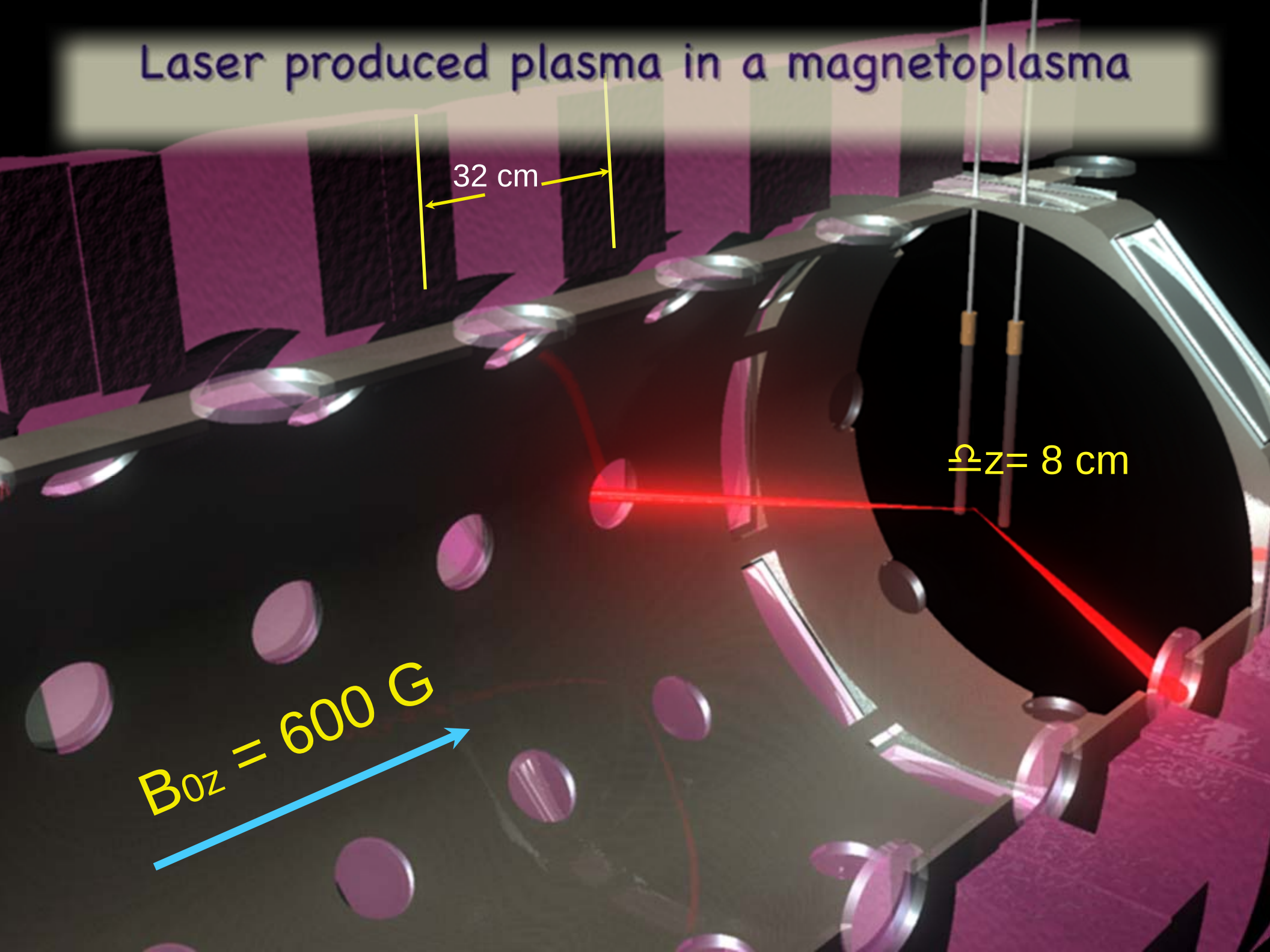


$z=7.5$ m

QSL's Quasi-Separatrix Layers

- a) Can now be used to identify regions in which field lines become stochastic
- b) A methodology is under development to relate QSL's to current layers and reconnection regions
- c) Reconnection can occur when no null points exist.
In 3D this is usually the case
- d) QSL's may soon be used to clarify the problem of the role of a guide field in reconnection

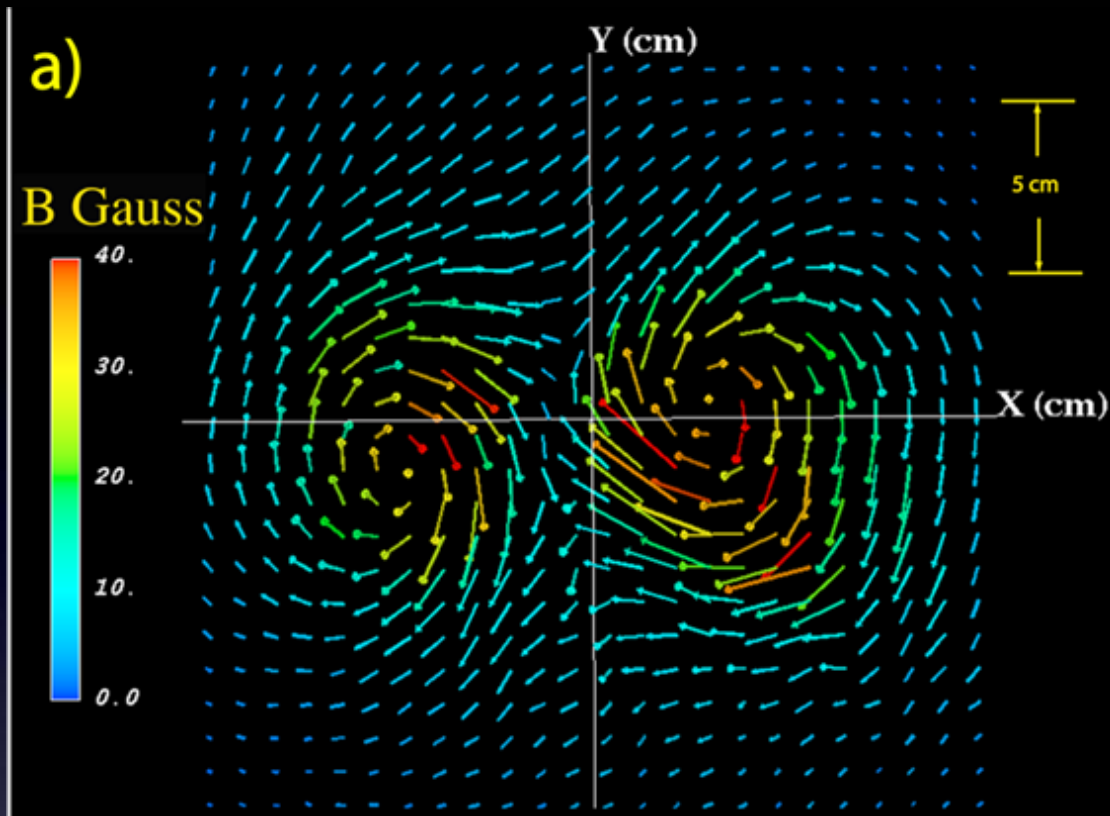
Laser produced plasma in a magnetoplasma



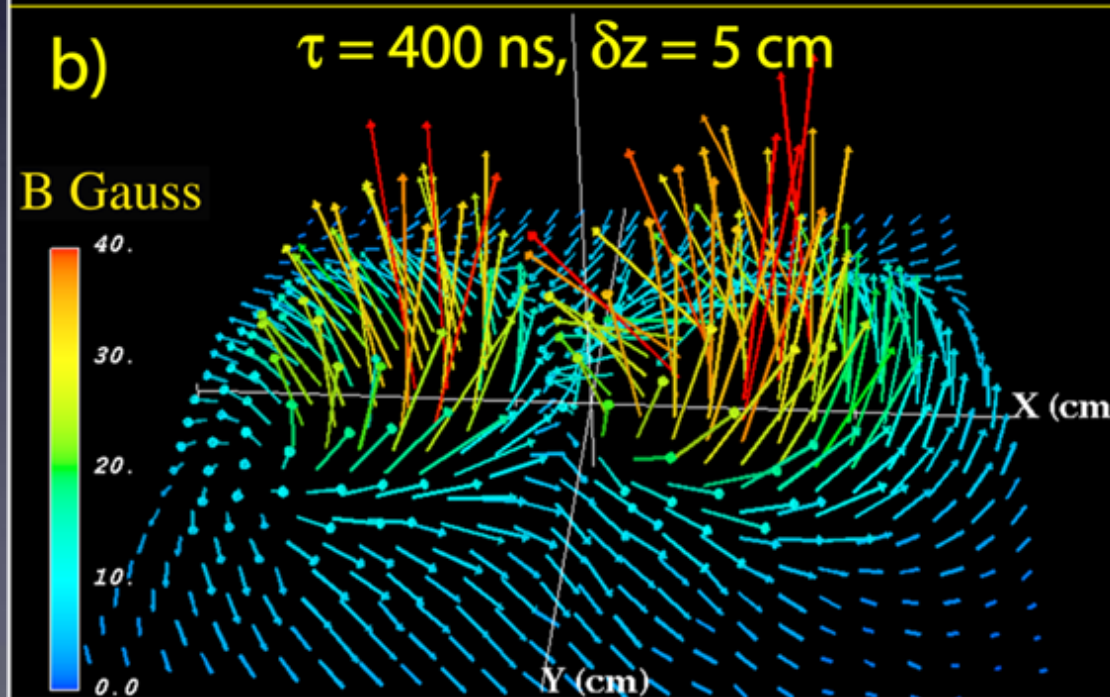
32 cm

$\Omega z = 8 \text{ cm}$

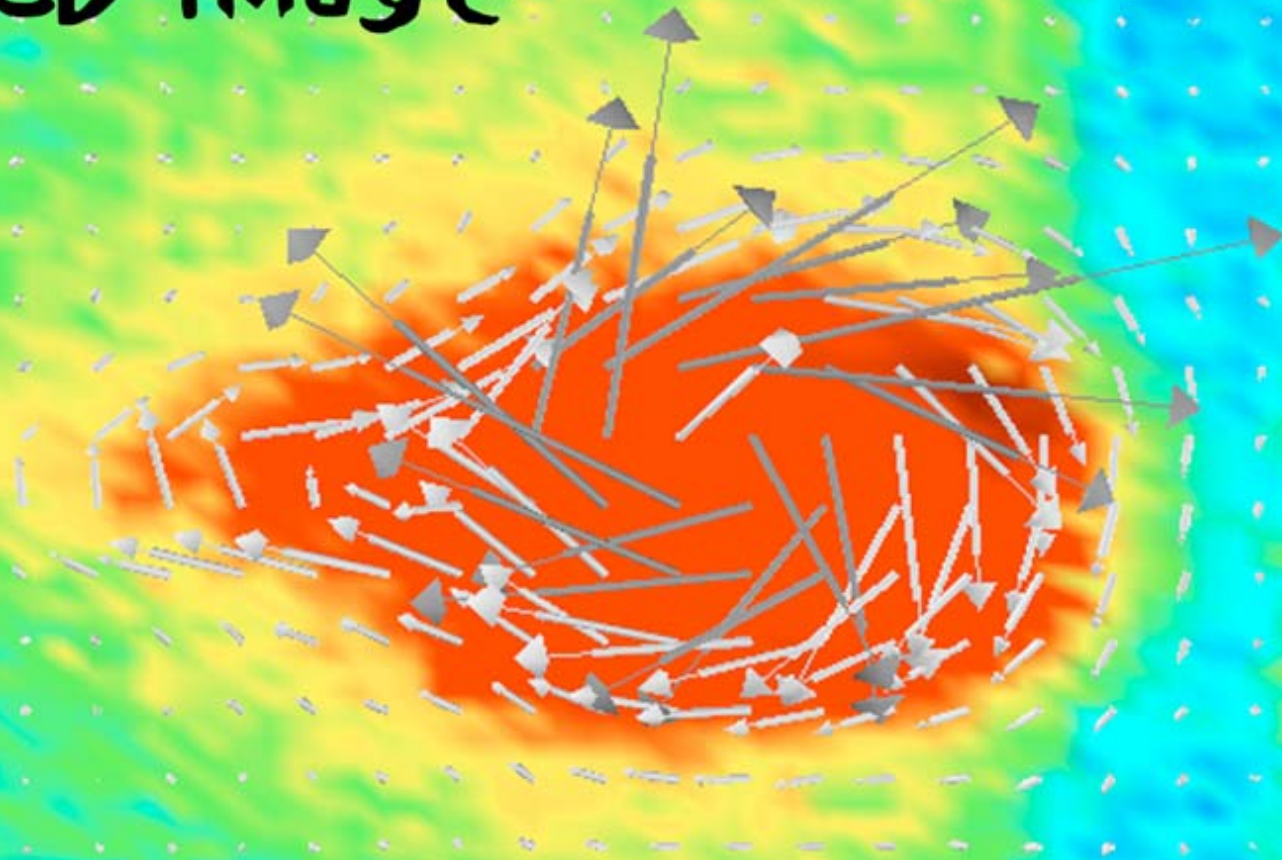
$B_{0z} = 600 \text{ G}$



Three
dimensional
neutral
sheet



$J = \text{curl}(H)$ **arrows**
10 ns **CCD image**

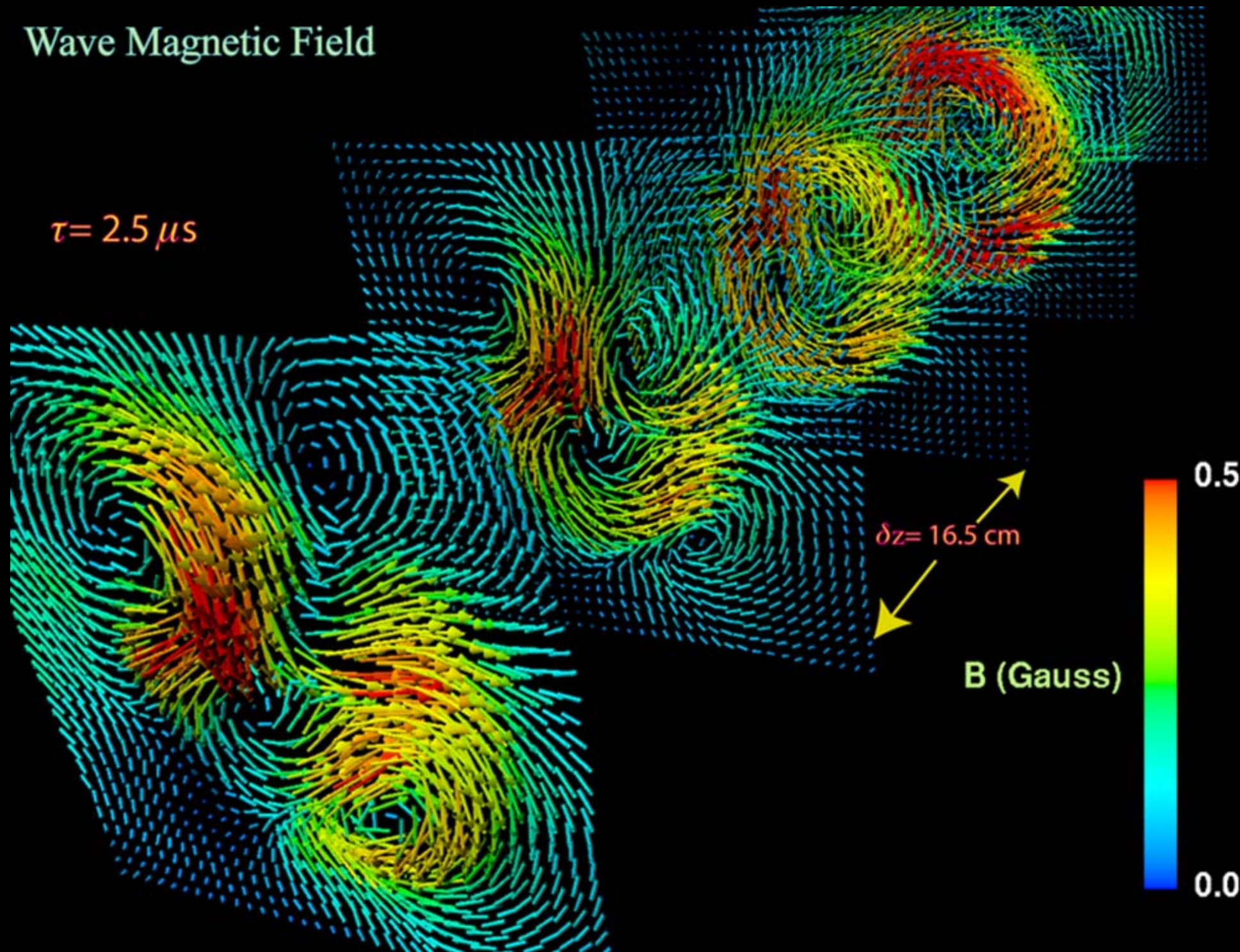


$t_{\text{delay}} = 1.36 \mu\text{s}$
 $dz = 1.0 \text{ cm}$

$\rightarrow 1 \text{ cm} \leftarrow$

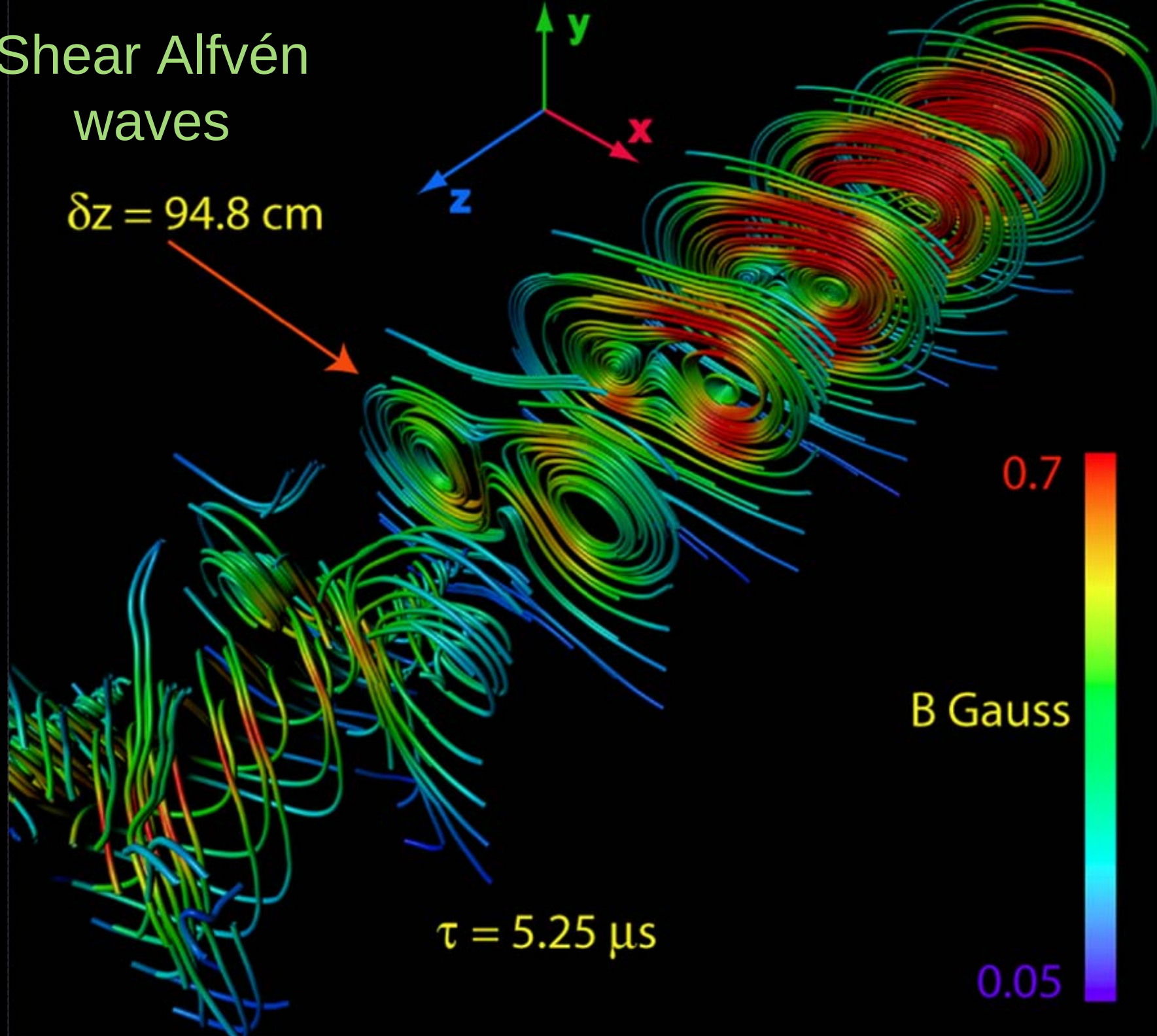
Wave Magnetic Field

$\tau = 2.5 \mu s$



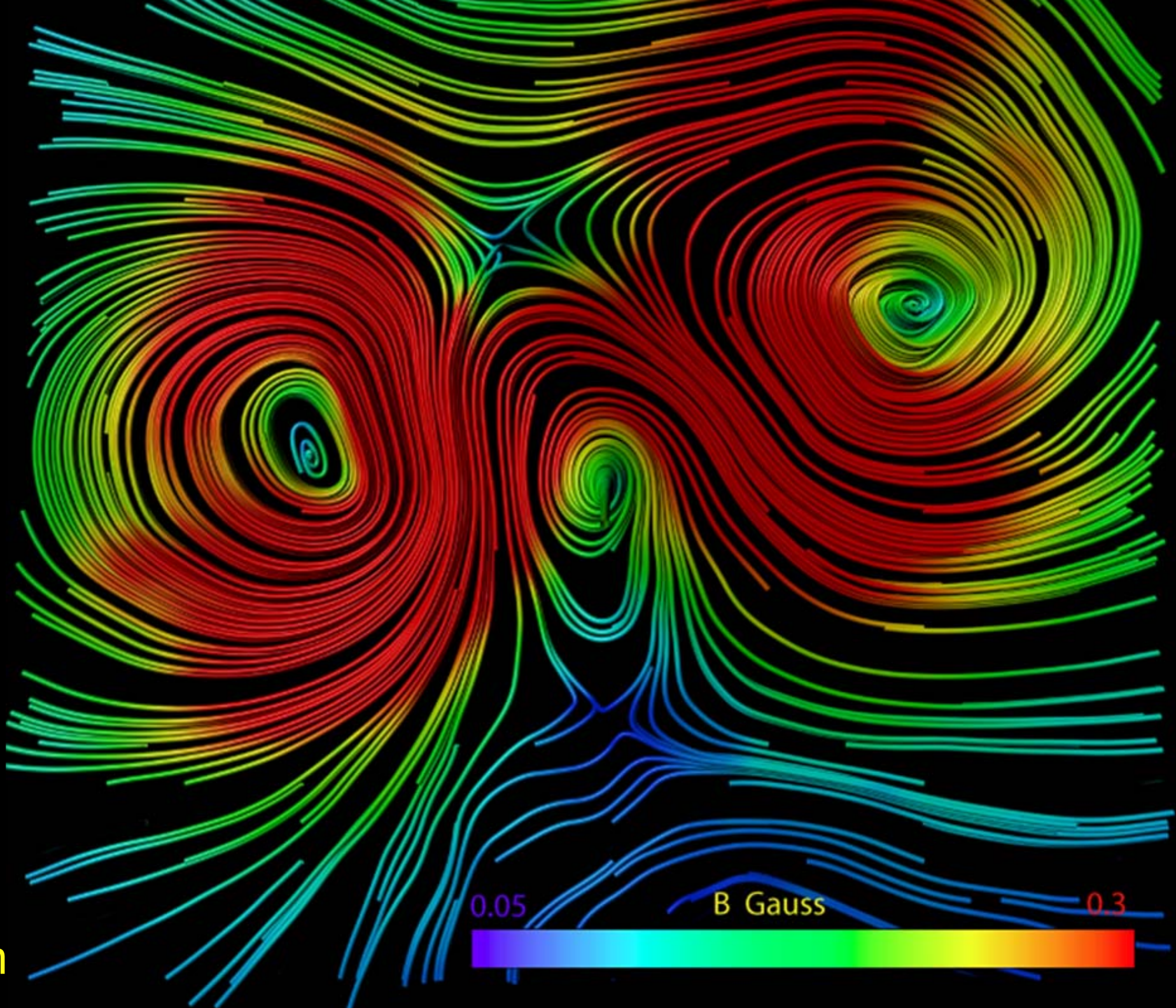
Shear Alfvén waves

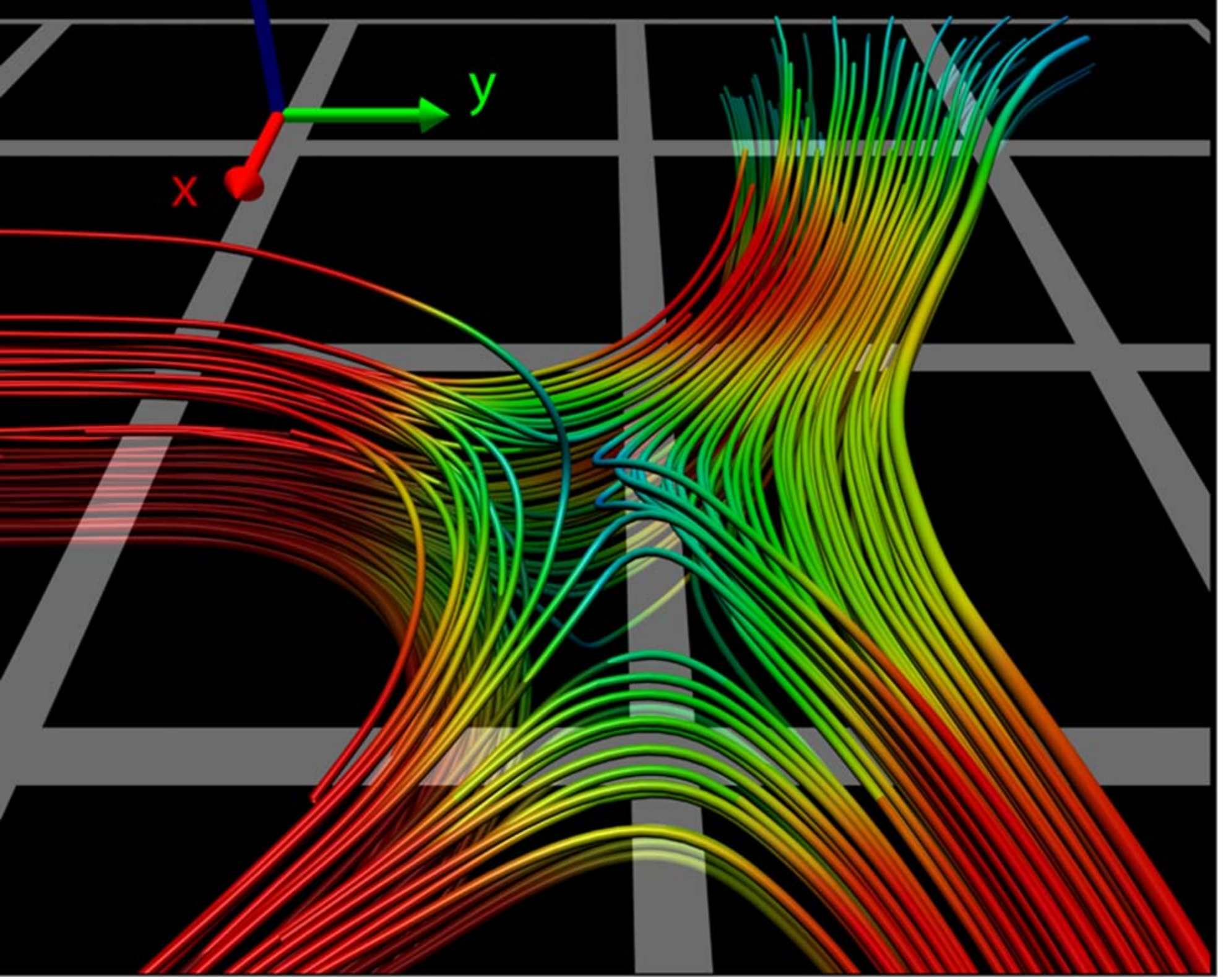
$\delta z = 94.8 \text{ cm}$

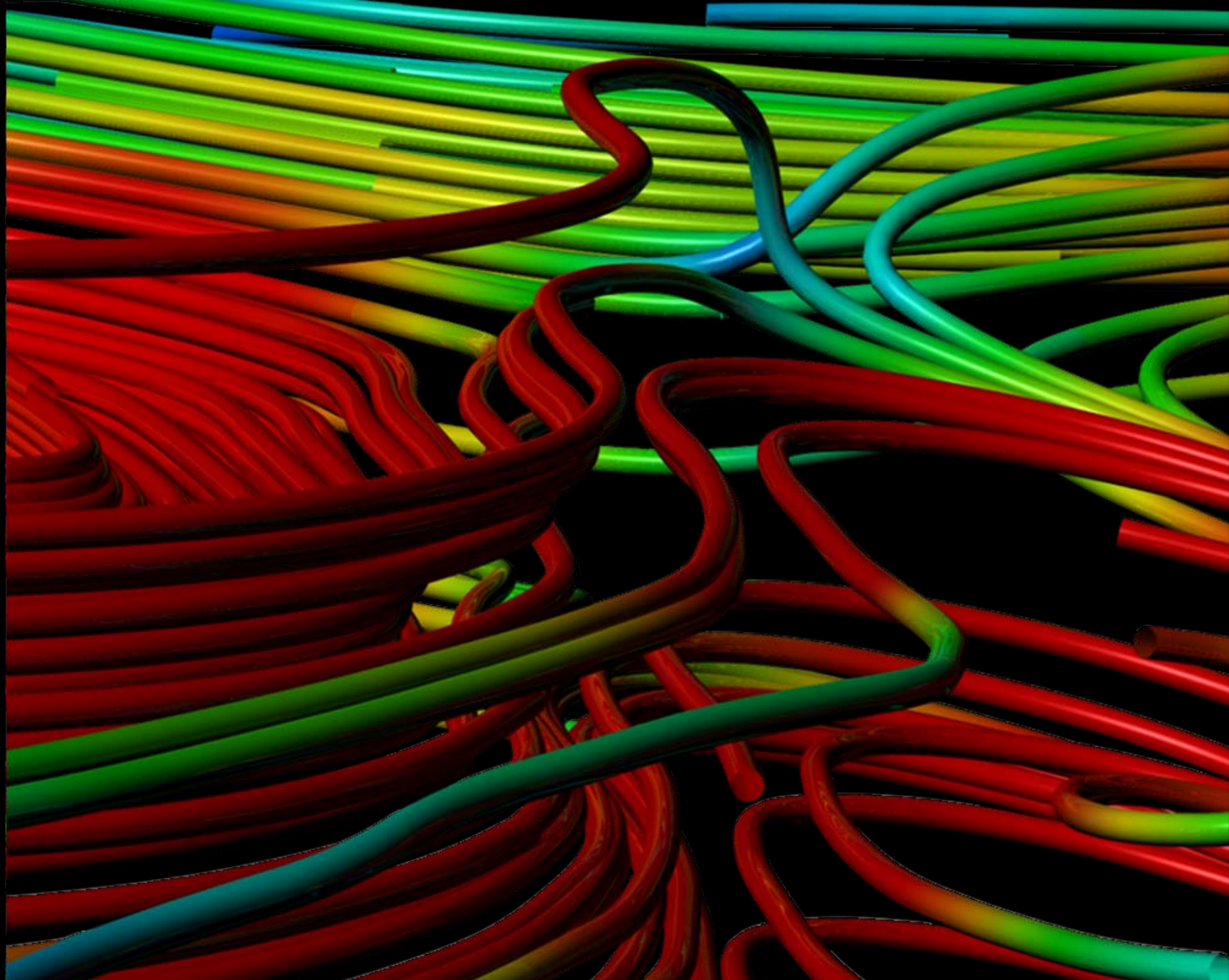


5.25
Os

$\Omega_z =$
95 cm



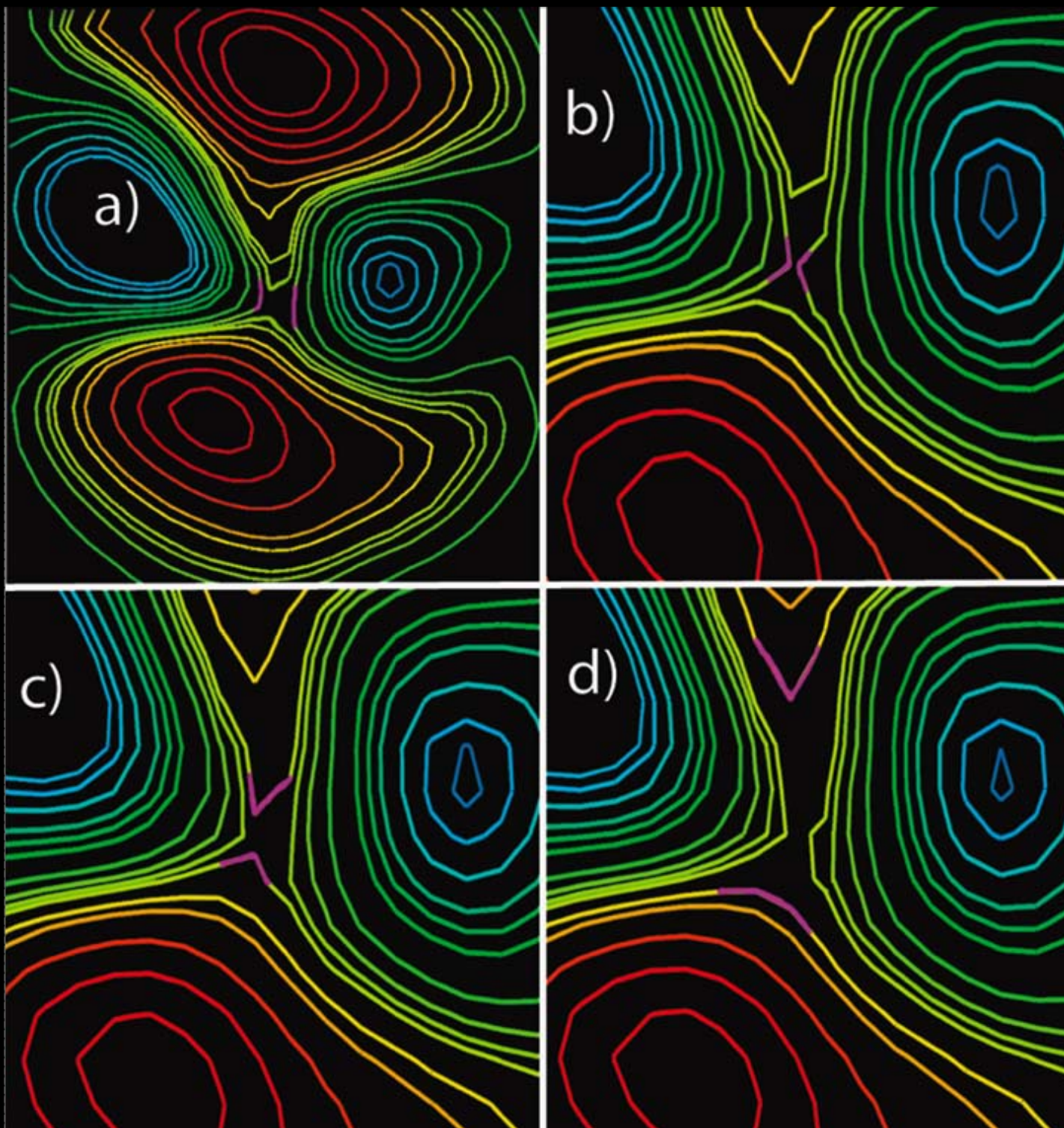




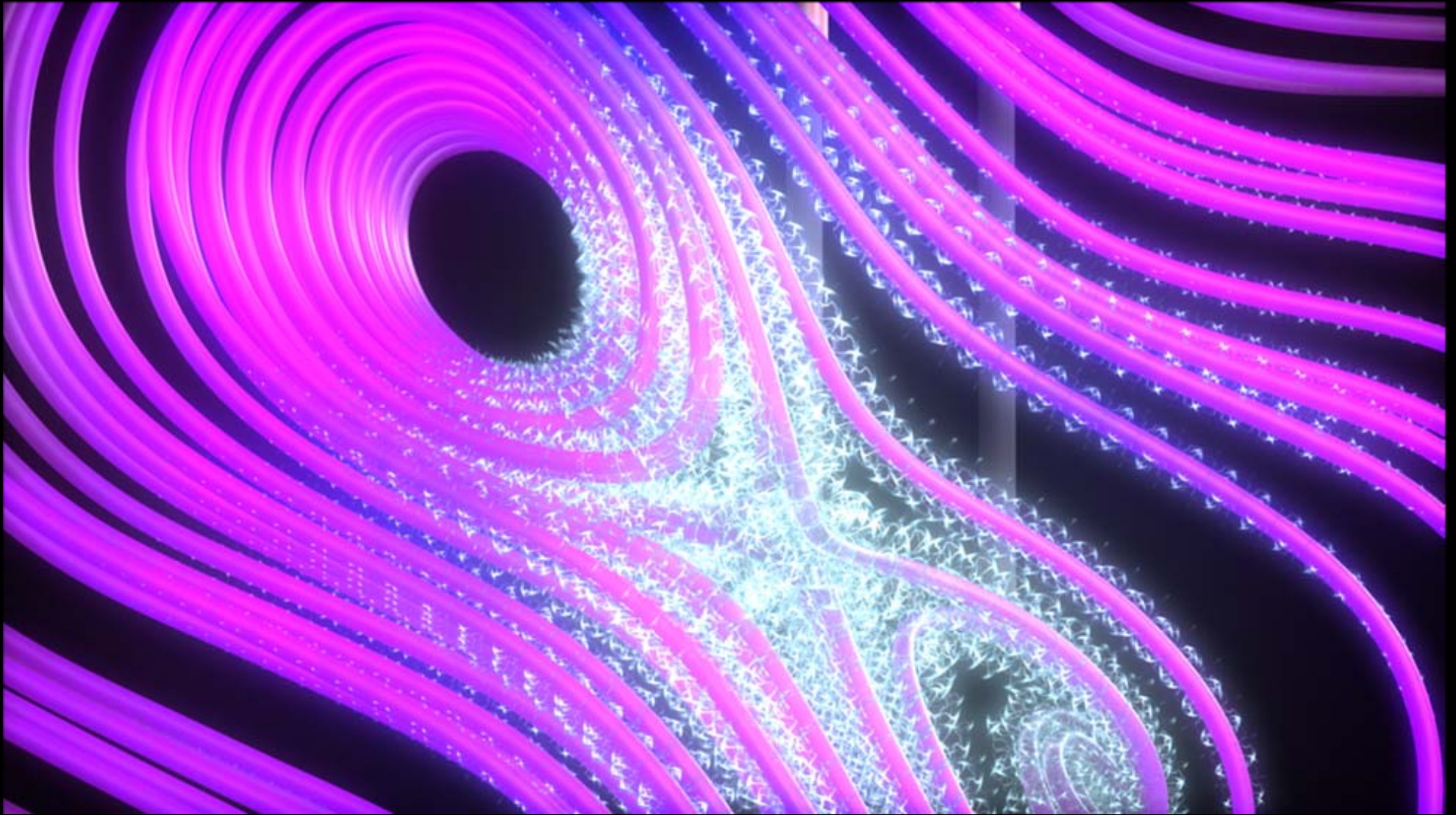
$\underline{\Omega}z =$
158 cm

◆ =
9.8
○s

$\underline{\Omega}t =$
10 ns



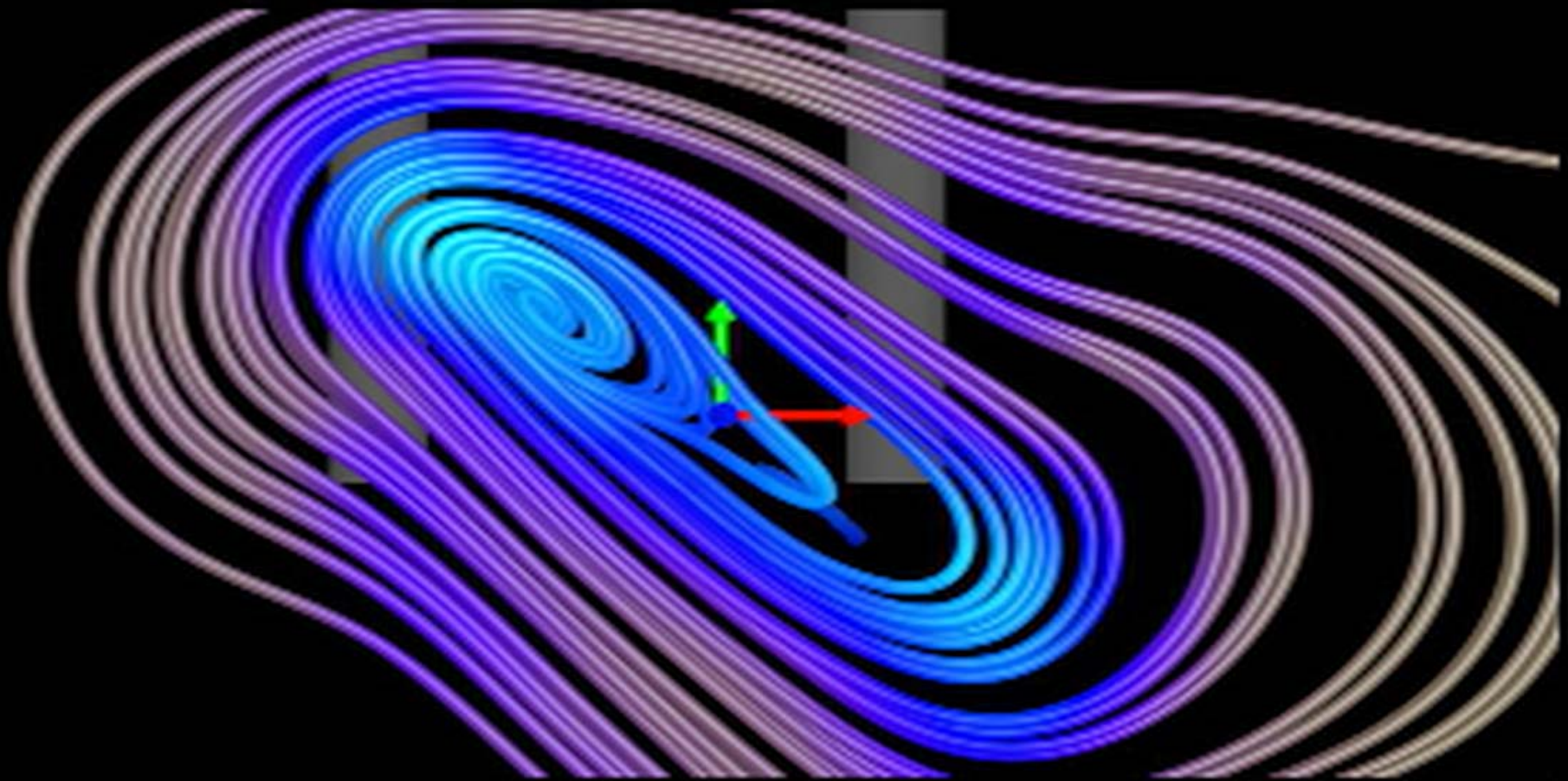
LAPD



$t = 5.12 \text{ } \text{\textcircled{O}}\text{s}$

Blines, E field

Colliding laser produced plasmas



Alfvén wave magnetic field

Colliding Laser Plasmas

◆ = 3.8 Os

QSL

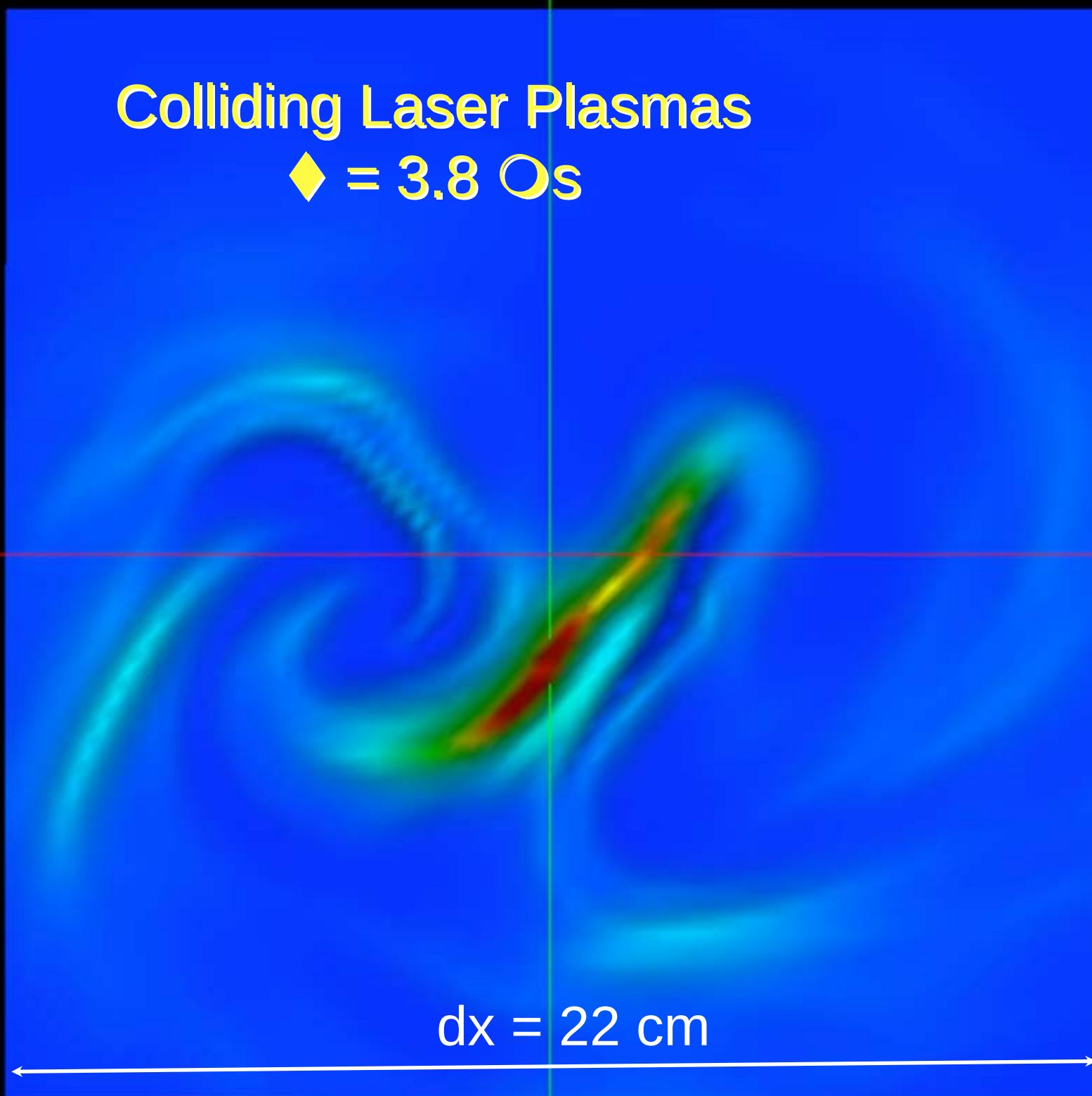
100.0

x

1.0

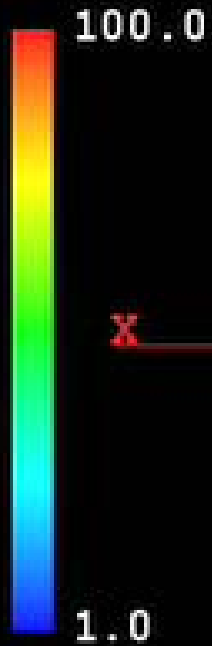
$B_{0z} = 10\text{G}$

$dx = 22\text{ cm}$

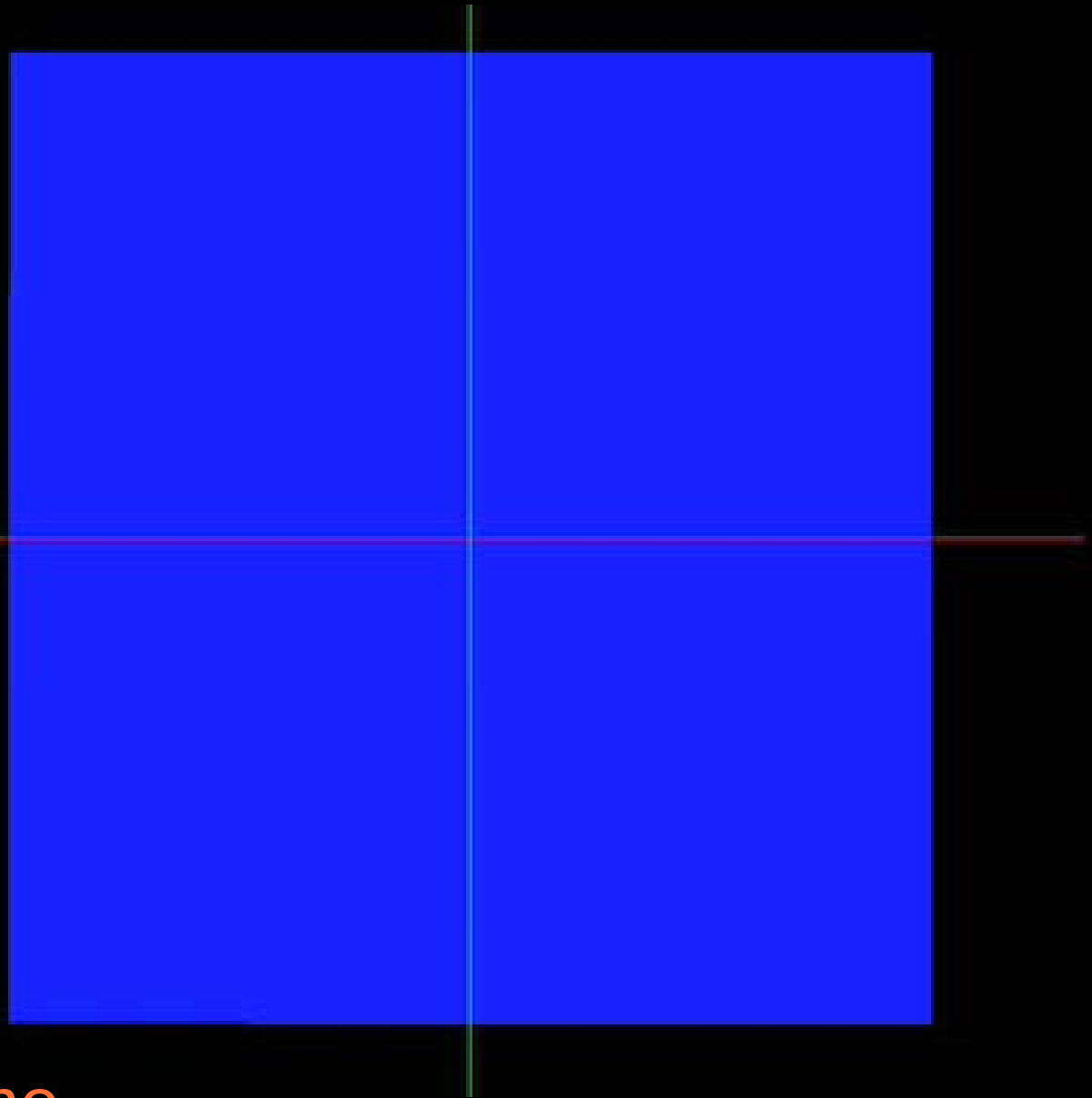


QSL for Alfvén waves

QSL

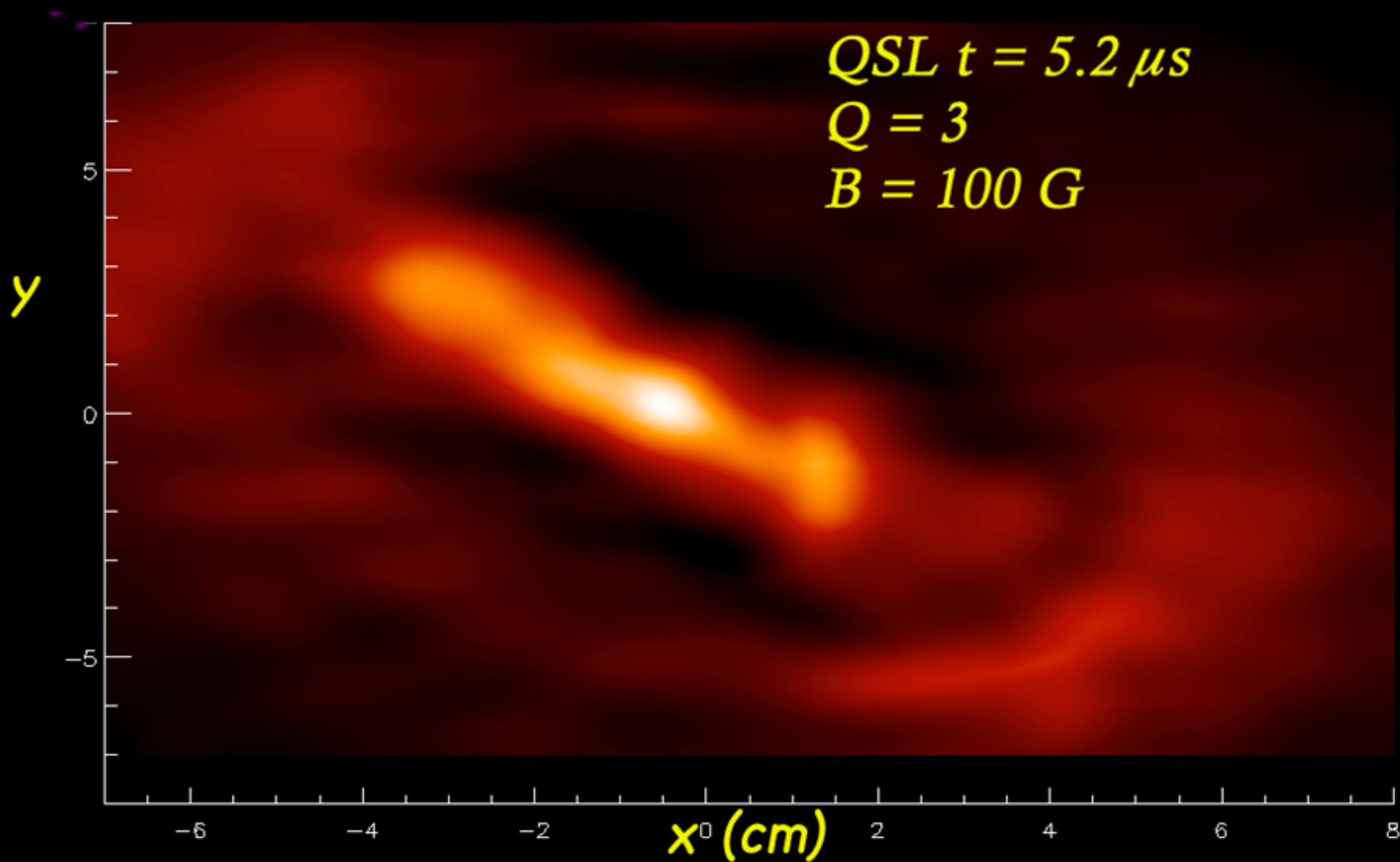


x



dt = 20 ns/frame
movie = 1 usec

$B_{0z} = 10\text{G}$



$$B_{0z} = 100\ G$$

In nature most reconnection is probably three dimensional and the concept of a separatrix is of limited use

A QSL has been observed when magnetic flux ropes merge **and** there is reconnection

The flux rope current system becomes, sheetlike, complex and return currents appear.

The QSL is an indicator that reconnection is occurring, we await a relationship between Q and the reconnection rate.