

1. Introduction

- To study **minor ion heating** process in the solar wind and solar corona;
- To address **possible heating mechanisms** between wave- particle interaction;
- To investigate **nonresonant genuine, pseudo and stochastic heating of ions by low-frequency Alfvén waves** via analytic solution & particle simulation.

Pseudo heating

- Kinetic process;
- Deformation of distribution function (maintain Maxwellian shape);
- Effective broadening under Alfvén wave spectra;
- Reversible;
- Heating efficiency: depends on wave amplitude
- No threshold requirement

Nonresonant heating

- Kinetic process;
- Initial pitch-angle scattering of newly created ions or violation of the first adiabatic invariant;
- Dissipation of wave fields;
- Irreversible;
- Heating efficiency: depends on wave amplitude
- No threshold requirement

Stochastic heating

- Kinetic process;
- Chaotic particle motions when satisfying the threshold condition; phase hole forms.
- Dissipation of wave fields;
- Irreversible;
- Heating efficiency: independent of wave amplitude
- Threshold: $k_x v_z B_w / (B_0 \Omega_c) \sim 0.1$

2. Analytic Theory

Wave Magnetic Perturbation

$$\delta \mathbf{B}_w^{linr} = \sum_k B_k \cos \phi_k \mathbf{i}_y, \quad \text{linearly polarized magnetic perturbations}$$

$$\delta \mathbf{B}_w^{cir} = \sum_k B_k (\cos \phi_k \mathbf{i}_x - \sin \phi_k \mathbf{i}_y), \quad \text{circularly polarized magnetic perturbations}$$

Nonresonant Genuine Heating Temperature (linearly polarized case)

$$T_{linx}^R = T_{liny}^R = T_{lin\perp}^R$$

$$= \frac{m_p}{2k_B \sqrt{\pi} v_{th}} \int_{-\infty}^{\infty} |v_{\perp} - U_{\perp}|^2 e^{-\left[\frac{v_{\parallel}(0)}{v_{th}}\right]^2} dv_{\parallel}(0)$$

$$\simeq T_0 \left(1 + \frac{m_i}{2m_p} \frac{\delta B_w^2}{\beta_p B_0^2} \right),$$

$$T_{lin\parallel}^R = \frac{m_p}{k_B \sqrt{\pi} v_{th}} \int_{-\infty}^{\infty} |v_{\parallel} - U_{\parallel}|^2 e^{-\left[\frac{v_{\parallel}(0)}{v_{th}}\right]^2} dv_{\parallel}(0)$$

$$\simeq T_0 \left(1 + \frac{3}{8} \frac{m_i}{m_p} \frac{\delta B_w^4}{\beta_p B_0^4} \right).$$

Nonresonant Genuine + Pseudo Heating Temperature

(linearly polarized case)

$$T_{linx}^{A+R} \simeq T_0 \left(1 + \frac{m_i}{2m_p} \frac{\delta B_w^2}{\beta_p B_0^2} \right),$$

$$T_{liny}^{A+R} \simeq T_0 \left(1 + \frac{3}{2\beta_p} \frac{m_i}{m_p} \frac{\delta B_w^2}{B_0^2} \right).$$

thermal speed $v_{th}^{tot} = \sqrt{v_{th}^R^2 + \xi^2}$ non-thermal speed

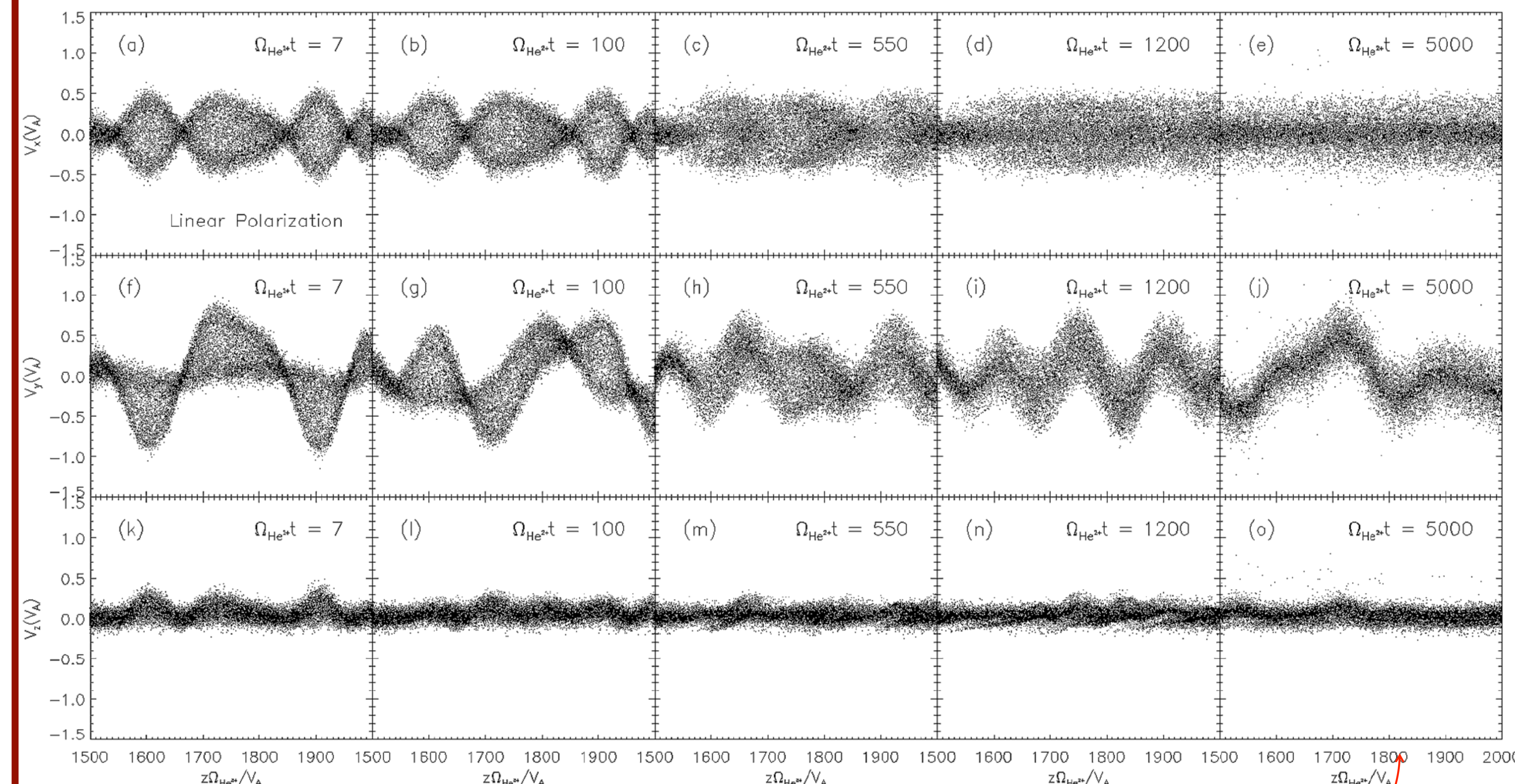
$$= \sqrt{\frac{2k_B T_{cir\perp}^{A+R}}{m_i}},$$

References

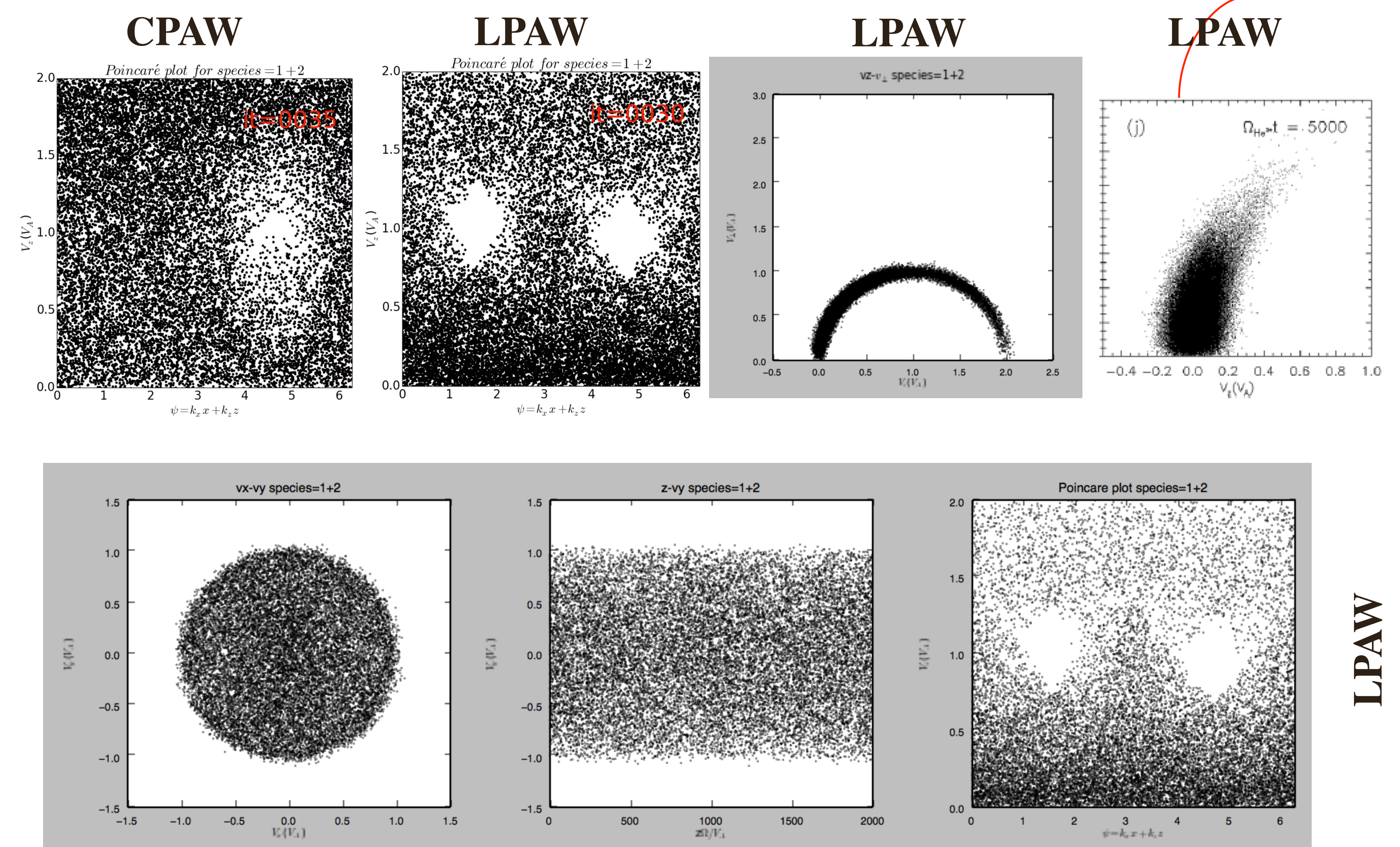
- C. F. Dong & N. Singh, Phys. Plasma **20**, 012121 (2013);
- C. F. Dong, Phys. Plasma **21**, 022302 (2014);
- C. S. Wu & P. H. Yoon, Phys. Rev. Lett. **99**, 075001 (2007);
- C. F. Dong & C. S. Paty, Phys. Plasma **18**, 030702 (2011);
- C. B. Wang et al, Phys. Rev. Lett. **96**, 125001 (2006);
- Z. H. Guo, C. Crabtree, and L. Chen, Phys. Plasma **15**, 032311 (2008);
- B. Chandran, B. Li, B. Rogers, E. Quataert, & K. Germaschewski, Astrophys. J., 720, 503 (2010)

3. Test Particle Simulation Results & Discussions

Scatter plots of the test particles for LPAW (nonresonant genuine and pseudo heating)



Scatter plots and the Poincaré plots of the test particles (stochastic heating)



4. Conclusions

- The pseudoheating is associated with wave fluctuations that contribute to the nonthermal broadening of VDFs due to the wave field forces;
- Although **the stochastic heating** is more **significant**, it occurs only when certain threshold condition is satisfied, and **takes much longer time** than nonresonant heating. Thus, **nonresonant heating** may also be **important** in some circumstances;
- The pseudoheating can be **tested** for solar wind **minor ions**, such as alpha particles at 1 AU based on current available data from *Wind* and *Helios* spacecrafts;
- For the near solar region, *in situ* observations from **high spatial and temporal resolution** space instruments are required, such as **the upcoming NASA Solar Probe Plus mission**.