

# Molecular Dynamics Simulations of the Hydrodynamic Transport Coefficients



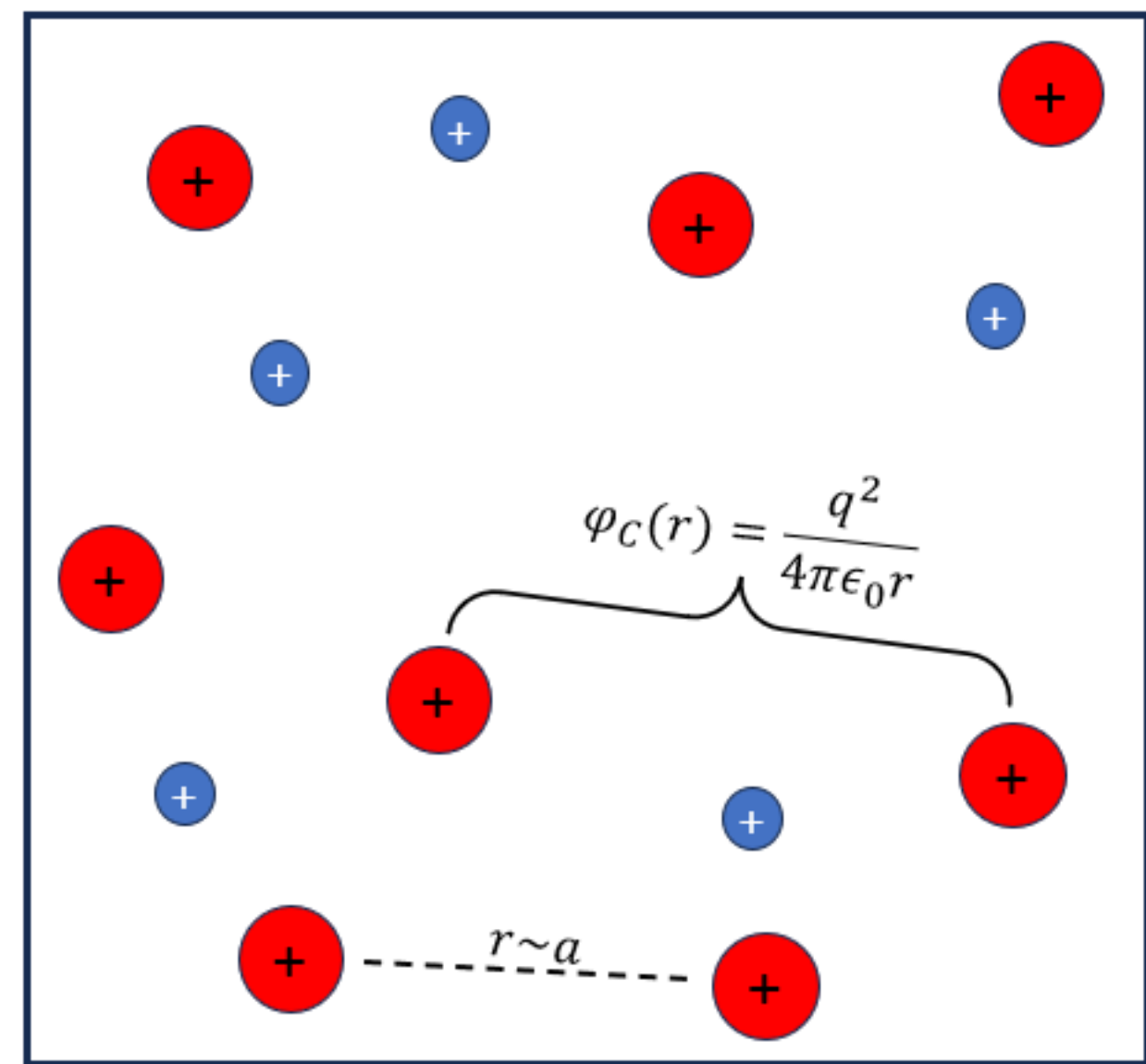
Briggs Damman\*, Jarett Levan†, and Scott Baalrud\*

\*Nuclear Engineering and Radiological Sciences, University of Michigan; †Applied Physics, University of Michigan

## Motivation

- The Chapman-Enskog solution of the Boltzmann equation has yet to be thoroughly validated
- Results are needed in strong coupling to provide a testbed for theories
- Two-component systems are historically overlooked in favor of one-component systems

## The Two-Component Plasma Model



- Repulsive two-component system
- Using open source code LAMMPS
- 5000 particles per species in 3D box with periodic boundary conditions
- Thermostat, then evolve particles in NVE ensemble

Electron coulomb coupling parameter:

$$\Gamma_e = \frac{q^2}{4\pi\epsilon_0 a_e k_B T} \quad (1)$$

## Non-Equilibrium Solutions

Transport coefficients are evaluated using the Green-Kubo relations:

$$\text{Electrical Conductivity: } \sigma = \frac{V}{3k_B T} \int_0^\infty dt \langle \mathbf{j}(t) \cdot \mathbf{j}(0) \rangle, \quad (2a)$$

$$\text{Electrothermal Coefficient: } \varphi = -\frac{V}{3k_B T^2} \int_0^\infty dt \langle \mathbf{j}(t) \cdot \mathbf{q}(0) \rangle, \quad (2b)$$

$$\text{Thermal Conductivity: } \lambda = \frac{V}{3k_B T^2} \int_0^\infty dt \langle \mathbf{q}(t) \cdot \mathbf{q}(0) \rangle, \quad (2c)$$

$$\text{Thermoelectric Conductivity: } \phi = \frac{V}{3k_B T} \int_0^\infty dt \langle \mathbf{q}(t) \cdot \mathbf{j}(0) \rangle, \quad (2d)$$

$$\text{Shear Viscosity: } \eta = \frac{V}{6k_B T} \sum_{i=1}^3 \sum_{j \neq i}^3 \int_0^\infty dt \langle \hat{P}_{ij}(t) \hat{P}_{ij}(0) \rangle, \quad (2e)$$

with the microdynamics described using the Irving-Kirkwood procedure:

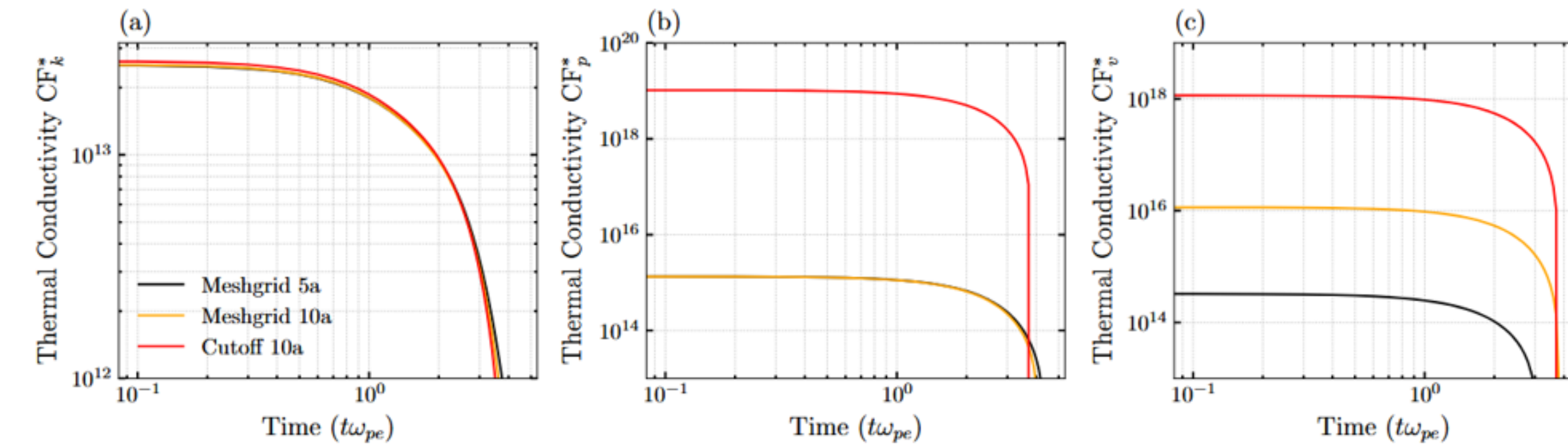
$$\mathbf{j} = \frac{1}{V} \sum_{i=1}^N q_i \mathbf{v}_i, \quad (3a)$$

$$\mathbf{P} = \frac{1}{V} \sum_{i=1}^N \left( m_i \mathbf{v}_i \mathbf{v}_i + \frac{1}{2} \sum_{j \neq i}^N \mathbf{r}_{ij} \frac{\partial \phi_{ij}}{\partial \mathbf{r}_i} \right), \quad (3b)$$

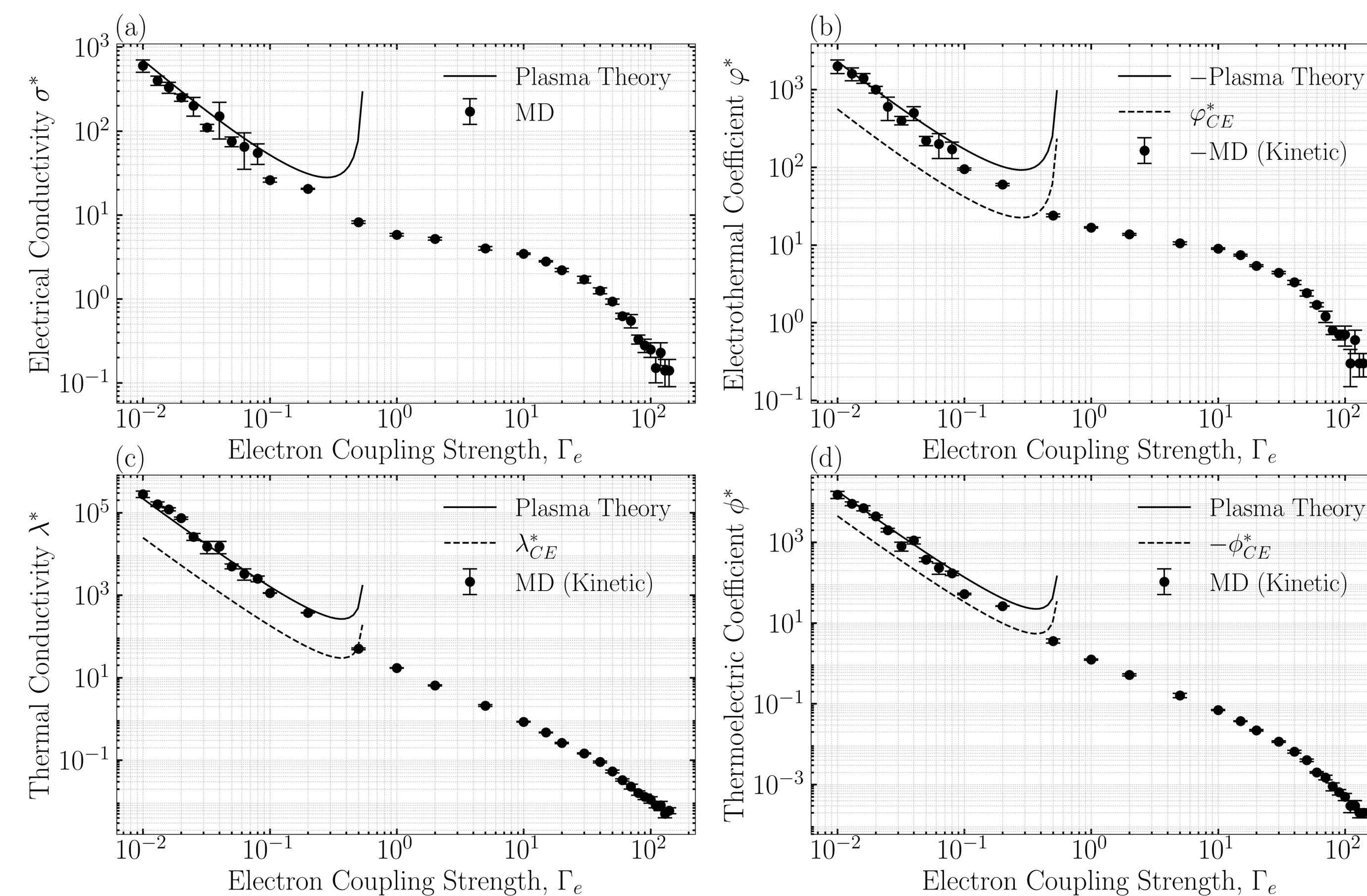
$$\mathbf{q} = \frac{1}{V} \sum_{i=1}^N \left[ \underbrace{\frac{1}{2} m_i |\mathbf{v}_i|^2}_{\text{kinetic}} + \underbrace{\frac{1}{2} \sum_{j \neq i}^N \phi_{ij}}_{\text{potential}} + \underbrace{\frac{1}{2} \sum_{j \neq i}^N (\mathbf{r}_i \cdot \mathbf{v}_i) \frac{\partial \phi_{ij}}{\partial \mathbf{r}_i}}_{\text{virial}} \right]. \quad (3c)$$

## Separation of Heat Flux Components

- Potential is infinite in repulsive system  $\Rightarrow$  heat flux is infinite
- Component significance:
  - Kinetic components dominate at weak coupling
  - OCP shows potential and virial components dominate when  $\Gamma \geq 10$



## Results



## Chapman-Enskog Solution

The Chapman-Enskog solution is the textbook kinetic theory closure to the MHD equations and relate the transport coefficients to system parameters:

$$\sigma_{CE} = 1.93 \frac{n_e q_e^2 \tau_e}{2m_e}, \quad (4a)$$

$$\varphi_{CE} = 0.78 \frac{k_B n_e q_e \tau_e}{m_e}, \quad (4b)$$

$$\lambda_{CE} = 1.02 \frac{n_e k_B^2 T \tau_e}{m_e}, \quad (4c)$$

$$\eta_{CE} = 0.96 n_i k_B T \sqrt{\frac{m_i}{2m_e}} \tau_e, \quad (4d)$$

$$\tau_e = \frac{3}{2\sqrt{2}\pi} \frac{(4\pi\epsilon_0)^2 \sqrt{m_e} (k_B T)^{3/2}}{n_e q_e^4 \ln \Lambda} \quad (5)$$

where

is the electron coulomb collision time.

## Comparison of Linear Constitutive Relations

The fluxes in both theories are equivalent; however, the organization of the fluxes are different leading to different transport coefficients. After relating the fluxes, the following equations may be obtained:

$$\sigma = \sigma_{CE} = 1.93 \frac{n_e q_e^2 \tau_e}{2m_e} \quad (6a)$$

$$\varphi = -\varphi_{CE} - \frac{5k_B}{2q_e} \sigma_{CE} = -3.19 \frac{k_B n_e q_e \tau_e}{m_e} \quad (6b)$$

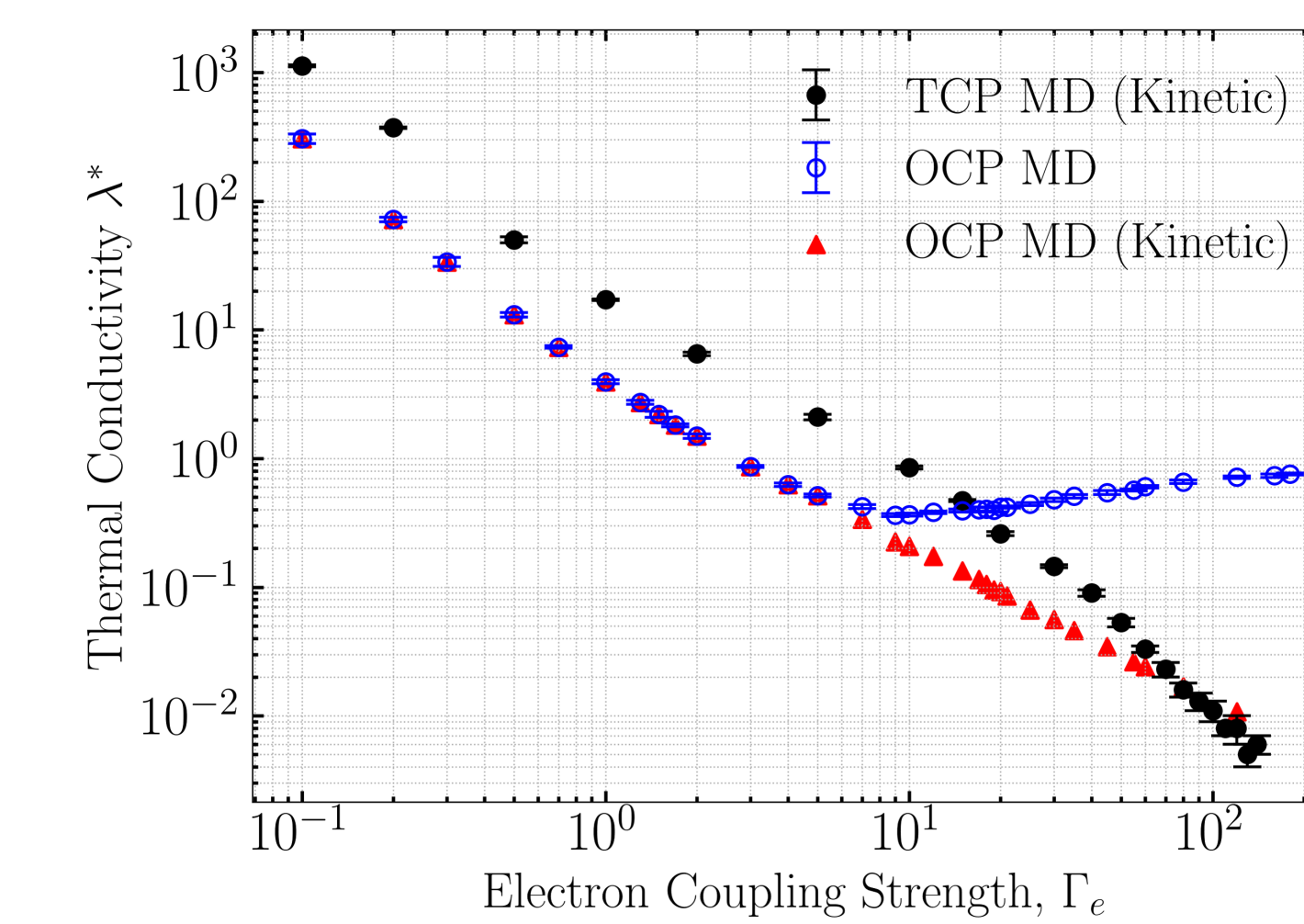
$$\phi = T \varphi_{CE} + \frac{2q_e}{5k_B T} \sigma_{CE} = 3.19 \frac{k_B T n_e q_e \tau_e}{m_e} \quad (6c)$$

$$\lambda = \lambda_{CE} + T \frac{\varphi_{CE}^2}{\sigma_{CE}} + \frac{5k_B T}{q_e} \varphi_{CE} + \frac{25k_B^2 T}{4q_e^2} \sigma_{CE} \quad (6d)$$

$$= 11.58 \frac{n_e k_B^2 T \tau_e}{m_e}$$

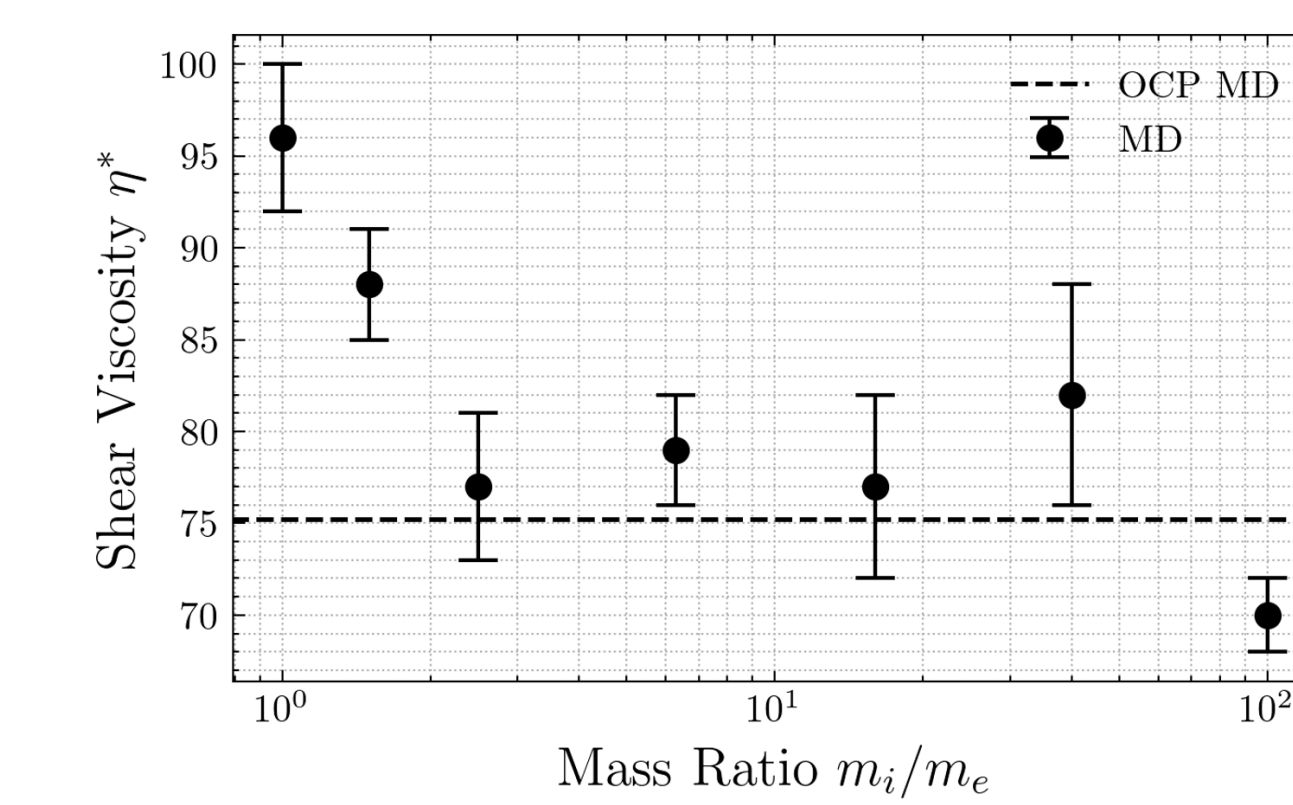
$$\eta = \eta_{CE} = 0.96 n_i k_B T \sqrt{\frac{m_i}{2m_e}} \tau_e. \quad (6e)$$

## Diffusive Contributions Dominate in the TCP

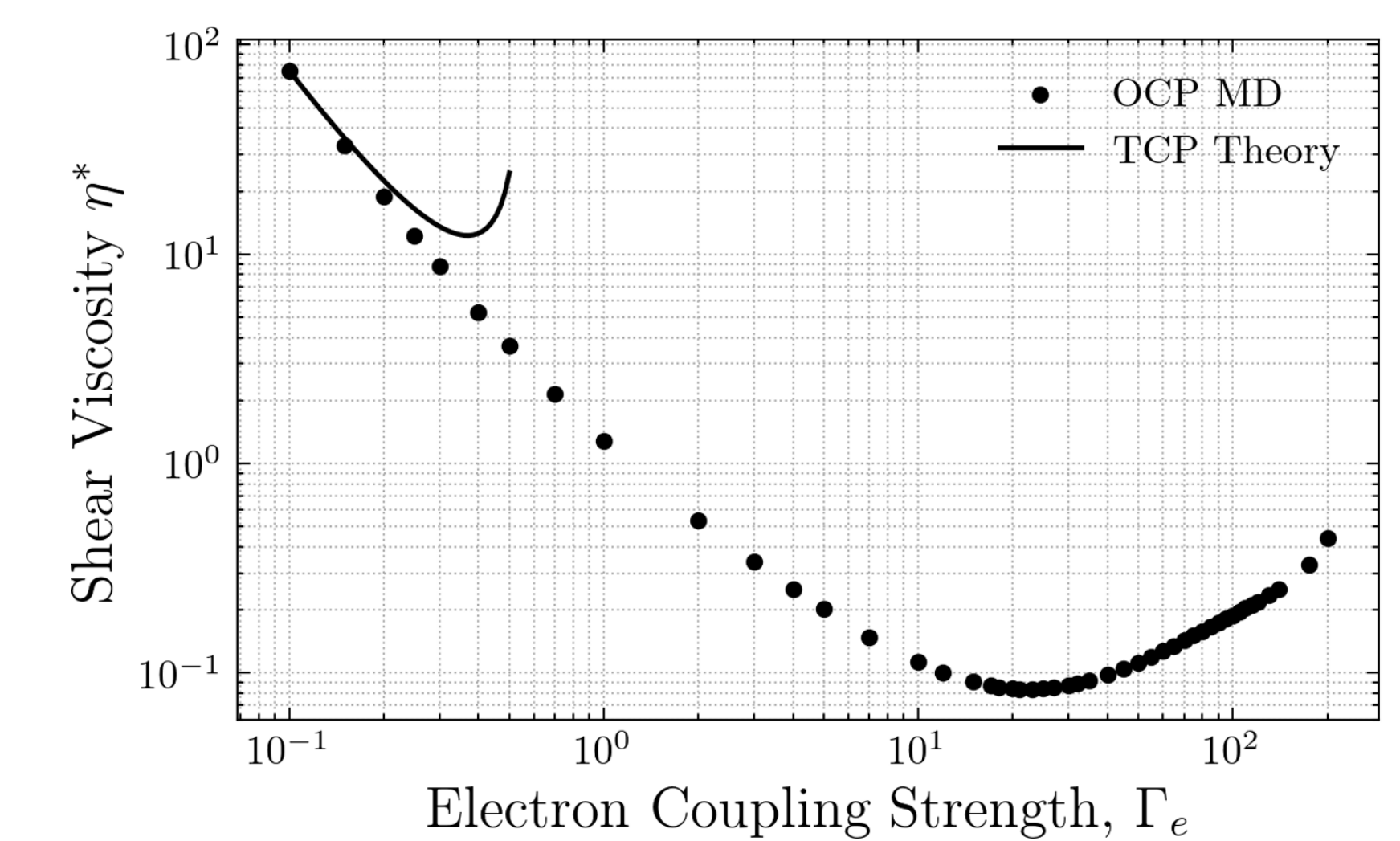


- One-component and two-component systems differ due to diffusive contributions
- Virial and potential components dominate at  $\Gamma \geq 10$

## Shear Viscosity Model



- MD Simulation of TCP Shear is intractable
- As mass ratio increases, TCP approaches OCP value
- Justifies using OCP as a model and we see good agreement with this comparison



## References

[1] J. Daligault, and S. D. Baalrud. "Determination of the shear viscosity of the one-component plasma." Physical Review E 90: 033105 (2014).