Simulation-Guided Design of a MegaJoule Dense Plasma Focus

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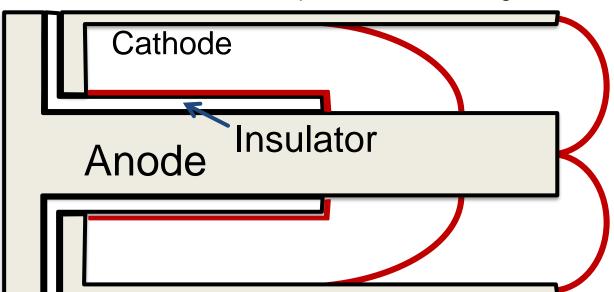
Outline

- Introduction to dense plasma focus (DPF)
- Neutron generation physics
- PIC modeling & benchmarks to measurements
- Simulation movies and restrikes
- Two ways current is diverted from pinch location:
 - Restrikes in plasma
 - Arcing behind gun
- Trends in the simulations → reduced order model
- Using the model to improve experiments
- Next questions to answer

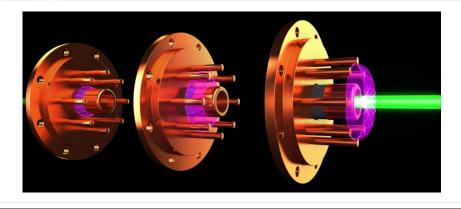


Dense Plasma Focus: A coaxial plasma railgun

The "Mather" DPF: an open ended coaxial gun



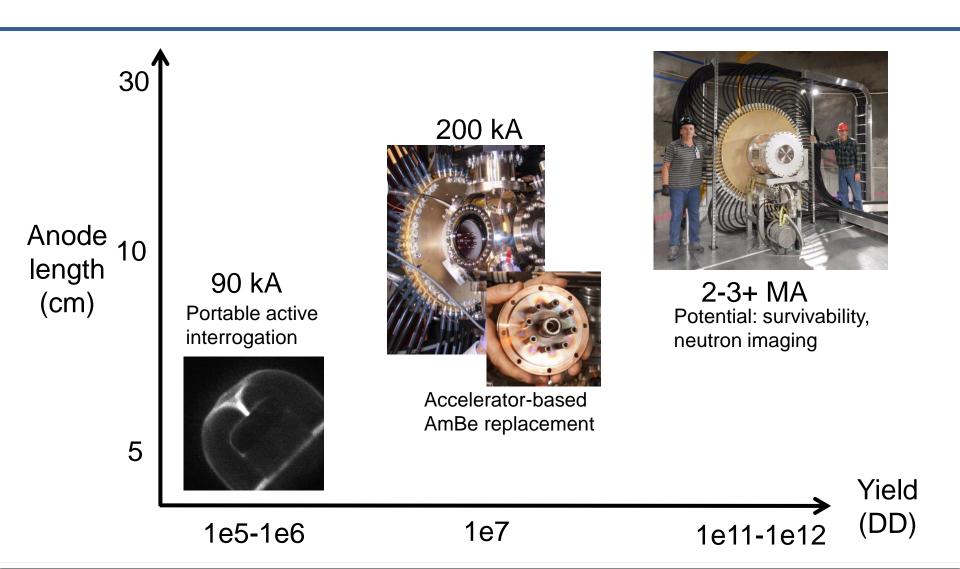




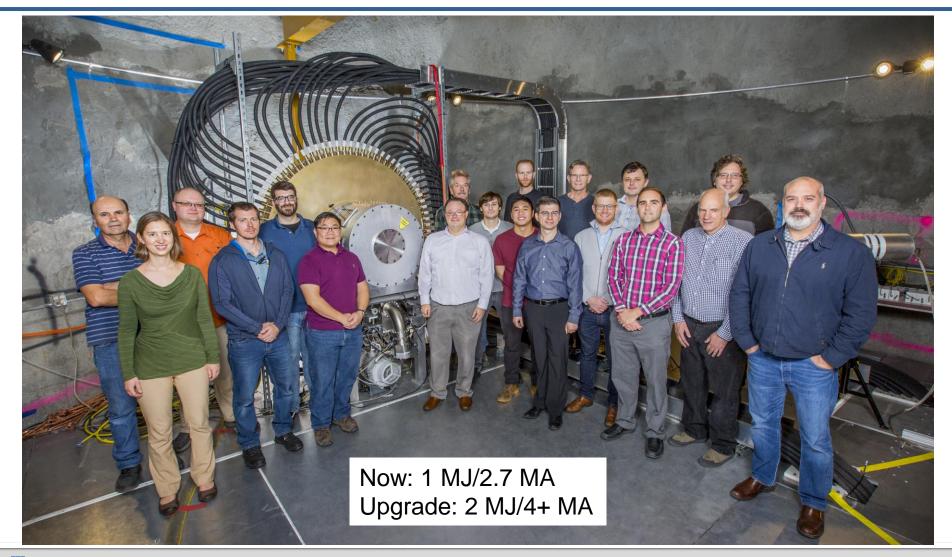
DPFs make

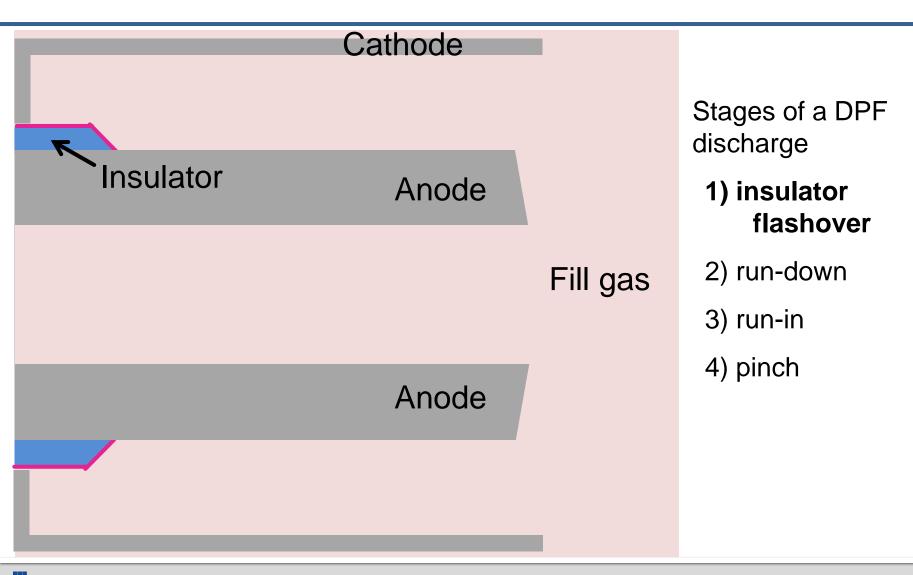
- energetic (keV to MeV) beams
- x-rays
- neutrons (for D or DT gas)

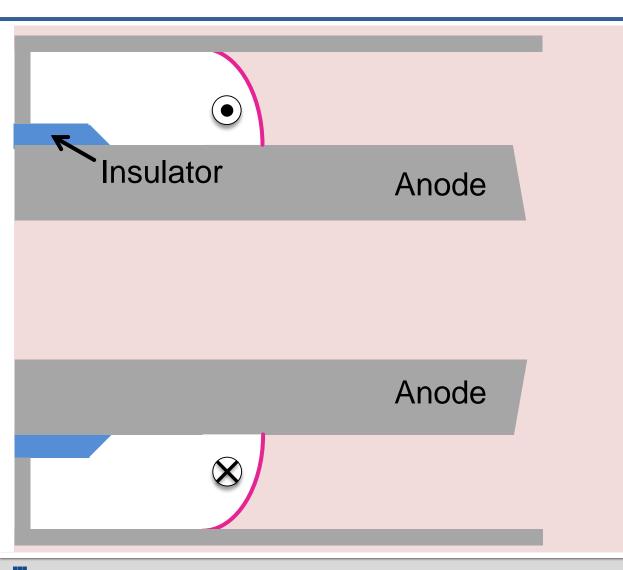
DPFs can be sized for relevant yield



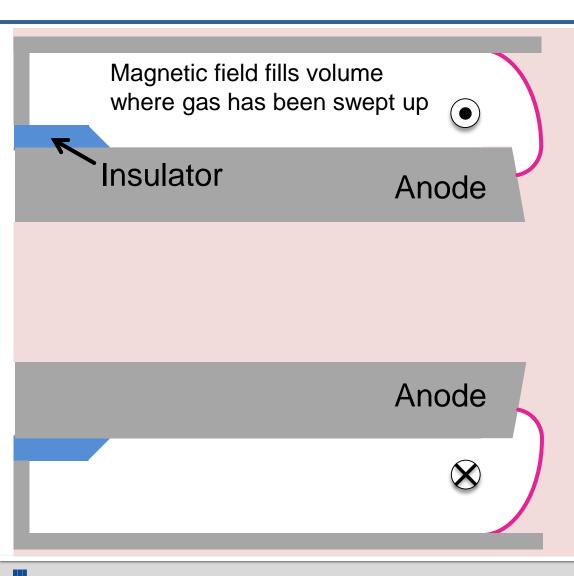
MegaJOuLe Neutron Imaging Radiography (MJOLNIR) design & build team



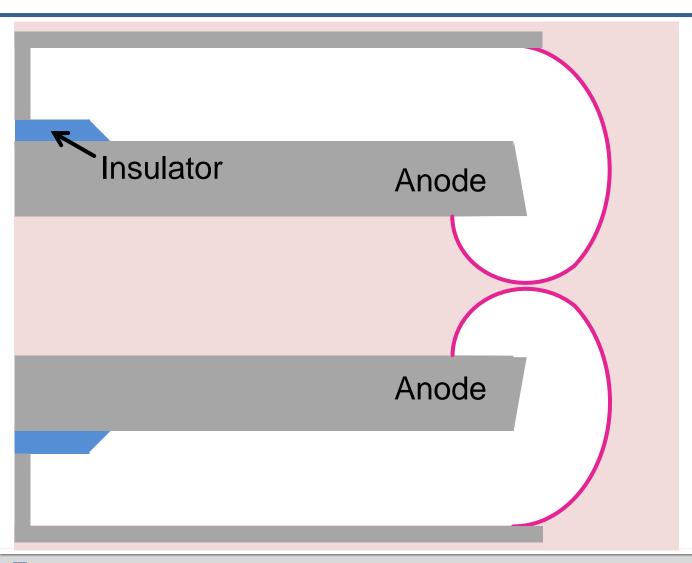




- 1) insulator flashover
- 2) run-down
- 3) run-in
- 4) pinch

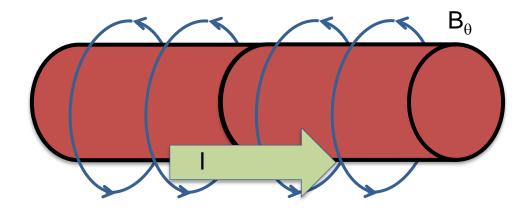


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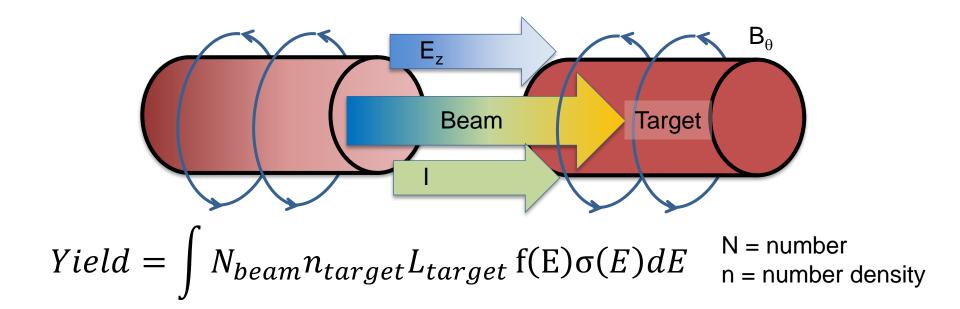


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Physics of neutron generation during z-pinch phase (according to 2D PIC simulations)

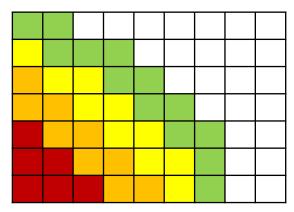


Physics of neutron generation during z-pinch phase (according to 2D PIC simulations)



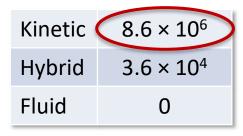
Frame:1 Delay:6.147us Exposure: 3ns

Kinetic (particle) code captures anomalous resistivity and beam formation in plasmas

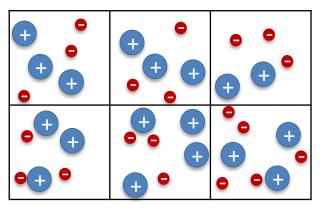


Fluid picture: each "pixel" is a fluid element with a density, temperature, and velocity

Kinetic model needed to get correct neutron yields in dense plasma focus (DPF)



Agrees with experiment



Kinetic picture: each "pixel" is a collection of particles; density, internal energy, and velocity are derived from collection

- Each "pixel" in the pinch region is really 1,000-10,000 particles
- 100-500 million particles per simulation
- We resolve electron cyclotron motion (~femtosecond time-steps)
- ~1 million time steps

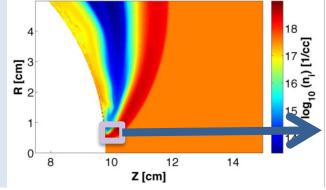
Details of LLNL PIC modeling

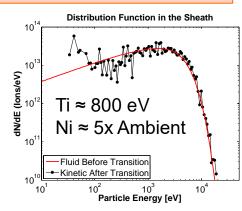
Simulations are performed in two coupled stages:

- 1. Single Fluid with MHD which establish the initial conditions and circuit dynamics
- 2. Kinetic PIC with Full Maxwell's equations starting 10-20 ns prior to the pinch phase

Fluid to Kinetic Transition

- Maps Local plasma distribution function to a drifting Maxwellian of kinetic particles
- Preserve during transition
 - Currents and Fields
 - 2. Plasma Conditions





Note: A typical 3 MA simulation takes 60-120 days of wall time A few million cpuHrs

Resolution

 Δx , $\Delta z = 100-400 \mu m$ $\Delta t = 0.25-250 \text{ fs}$

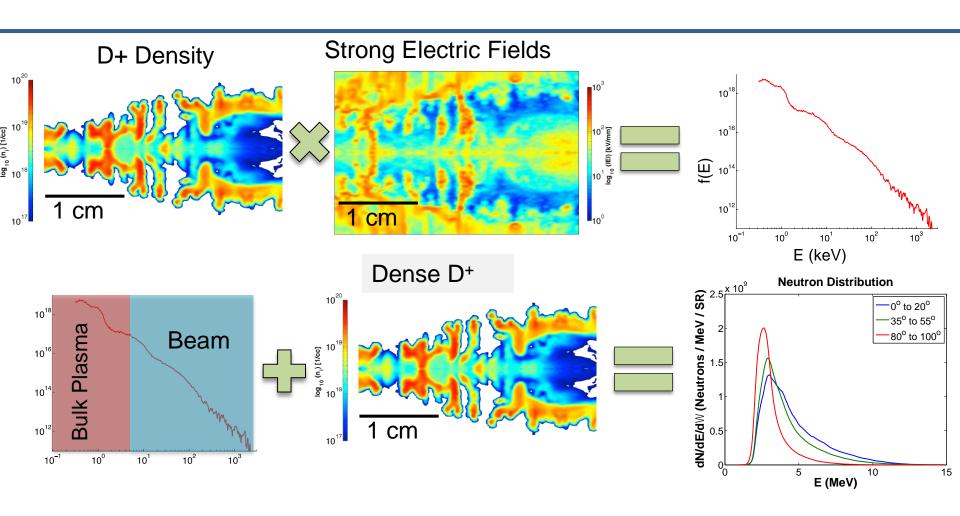
Physics Models

- Collisions
- D-D/T Fusion packages

Circuit Driver Full Maxwell Equations Implicit Particle advance

- Direct Implicit Scheme
- Full Matrix Inversion for Field Advance
- Currently resolving ω_c across most of the simulation

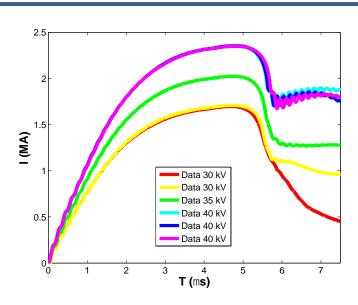
How a Dense Plasma Focus produces Neutrons



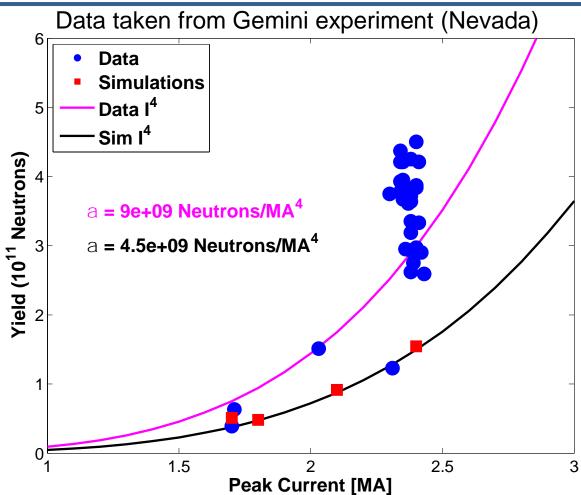
Beam Target dominates the total yield over thermonuclear



Model agrees within a factor of 2 with experimental yields



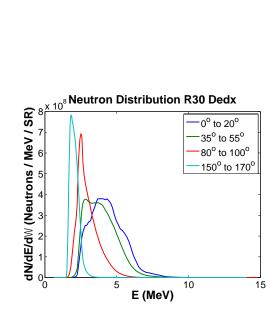
We have typically observed overall fairly good agreement but typical under predict the yield by roughly a factor of 2x for higher current MJ DPF shots

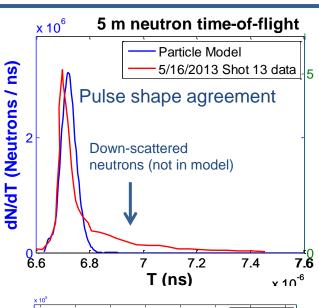


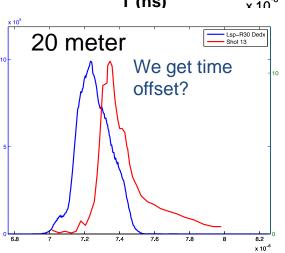
Similar scaling to Experiment but at half the coefficient

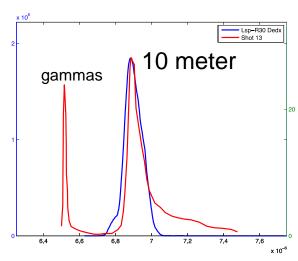


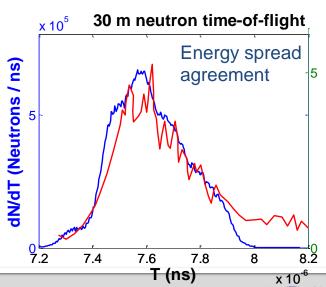
Near/far nToF agreement indicates that simulated pulse shape/neutron energies are reasonable



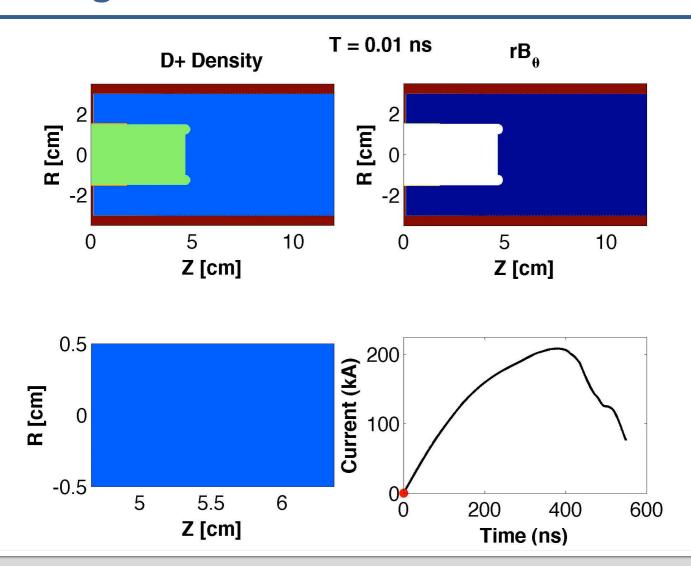




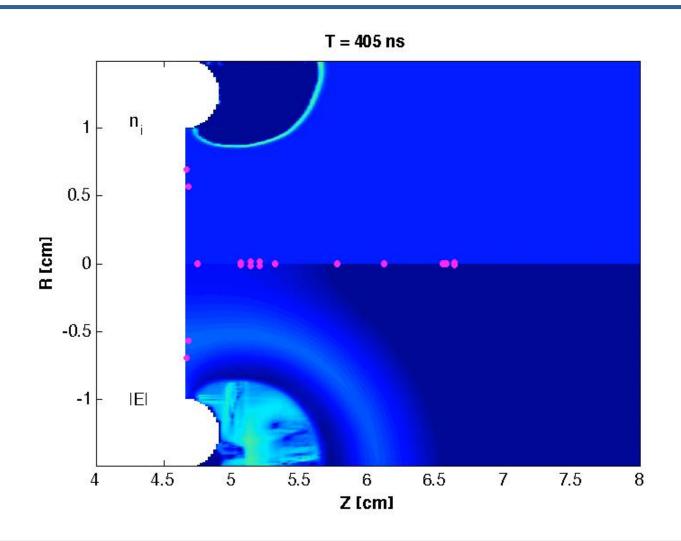




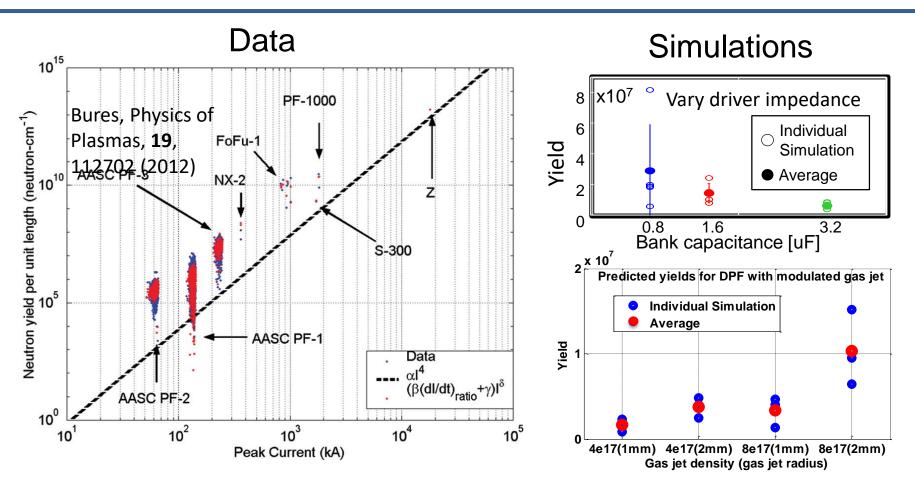
Example of a 200 kA fluid-to-kinetic simulation including restrikes



Particles accelerated across the gap to >1 MeV

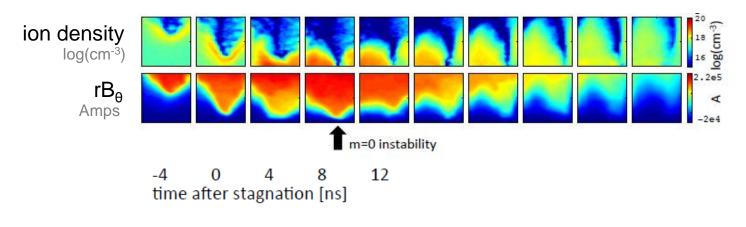


DPFs show variable measured yields and kinetic simulations reproduce variability/stochasticity

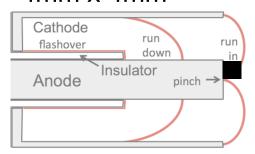


<u>Challenge:</u> Find ways to increase average yield and simultaneously improve shot-to-shot consistency

High-yield pinches always exhibit strong m=0 instability in simulations

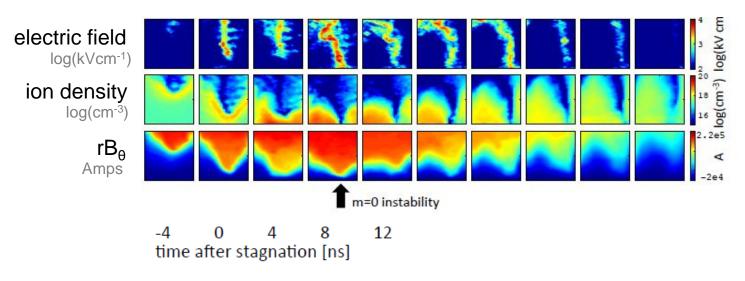


zoom region 4mm x 4mm

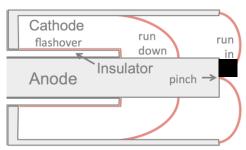


High-yield pinch

Electric field sets up in the gap between two separated blobs of plasma



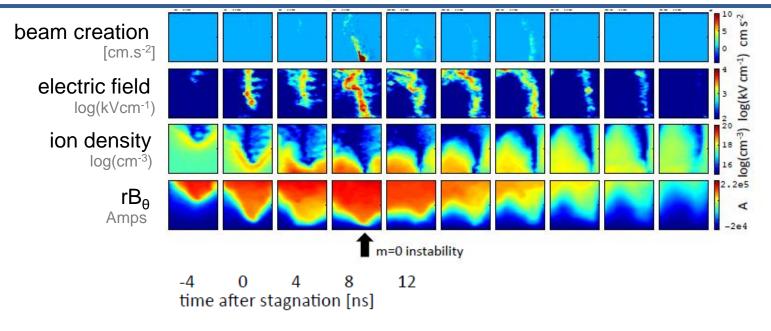
zoom region 4mm x 4mm



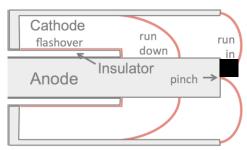


High-yield pinch

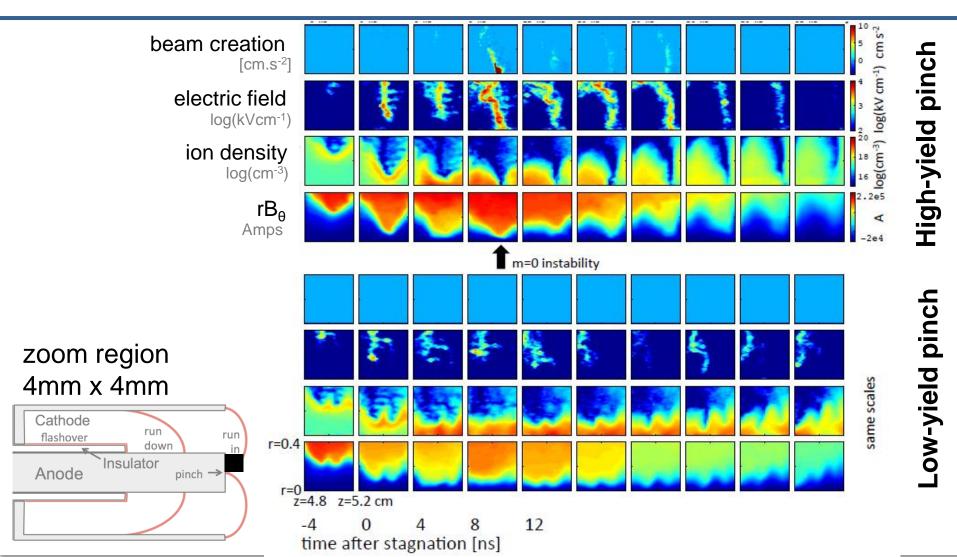
Beam forms on axis when electric field sets up there



zoom region 4mm x 4mm

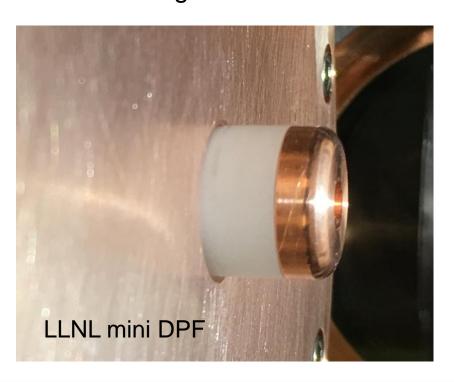


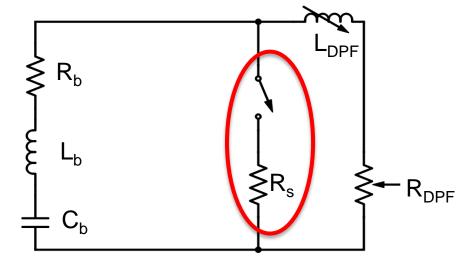
In a low yield simulation, a full m=0 never sets up, possibly due to restrike currents



Similar to restrikes, current arcing behind the gun diverts current from pinch

- Mystery: why wasn't LLNL mini DPF producing expected yields?
- Hypothesis: current is flowing behind the gun

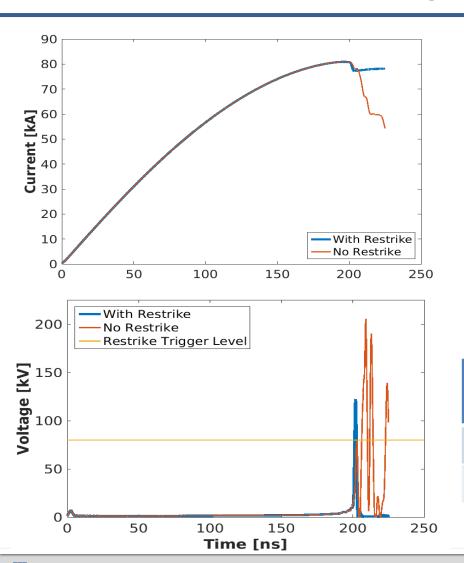


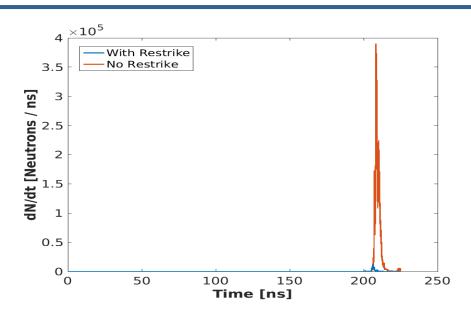


| C_b | 400 nF |
|----------------|---|
| L _b | 50.3 nH |
| R_b | $80~\text{m}\Omega$ |
| V_s | 80 kV |
| R_s | $1 \text{ m}\Omega - 1 \text{ k}\Omega$ |
| L_DPF | 4.5 nH |

We can add a parasitic current path in the simulation behind the gun to simulate arcing behind the gun

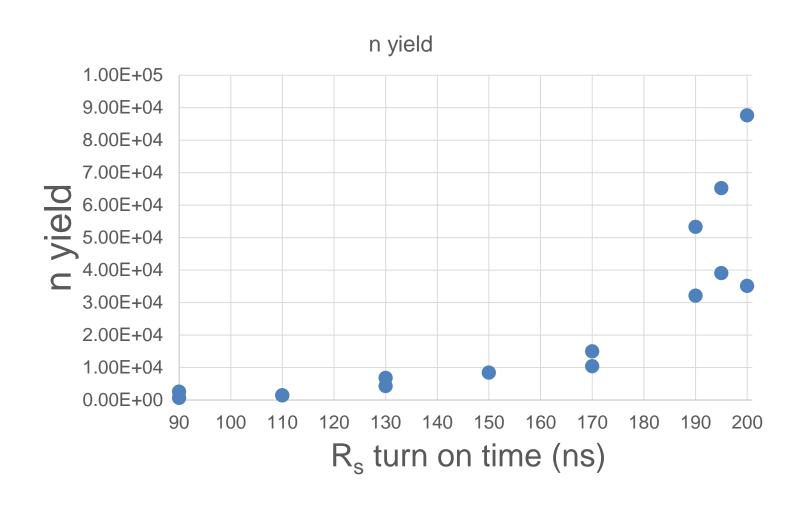
Arcing behind the gun can introduce significant reduction in neutron yield



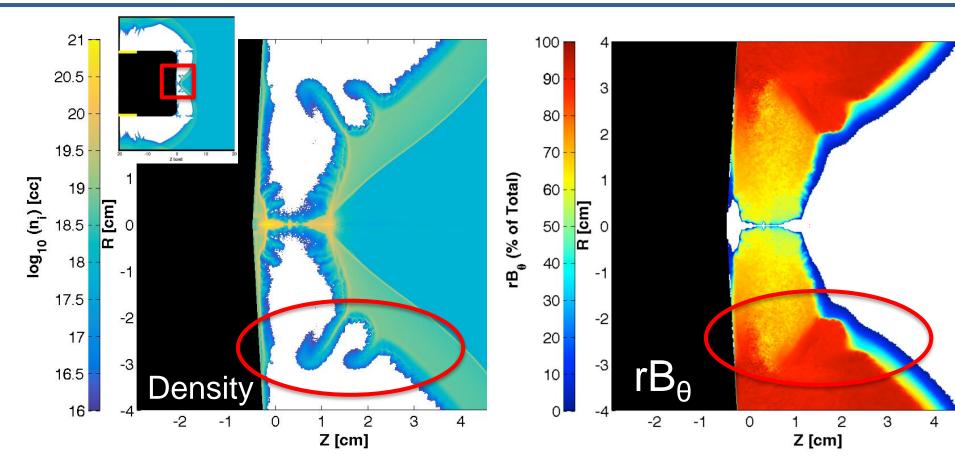


| Case | Trigger Voltage [kV] | Yield (10 ⁶ Neutrons) |
|-------------|-------------------------|-------------------------------------|
| restrike | 80 | .02 |
| no restrike | 8000 | 0.9 |

The earlier the arcing, the worse the yield

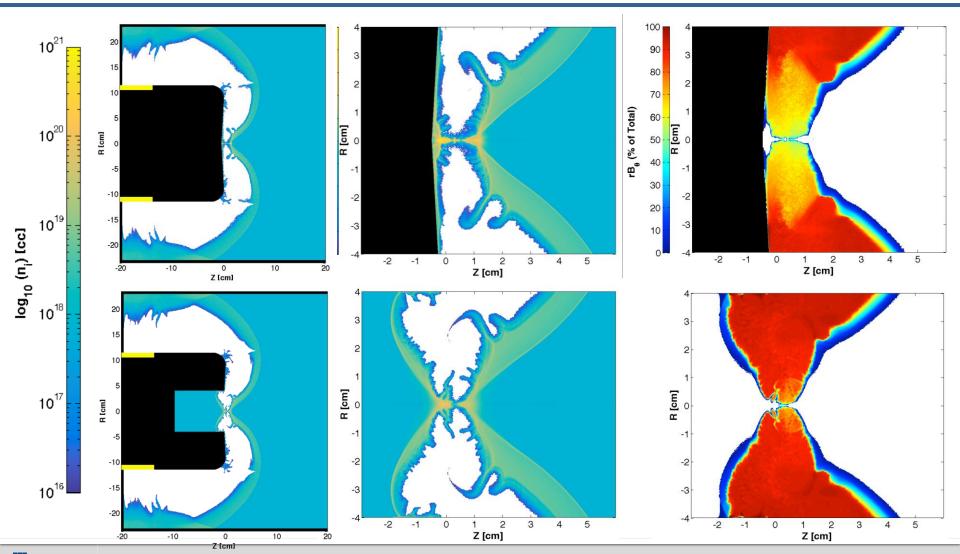


Simulations indicate that anode-to-sheath restriking is common

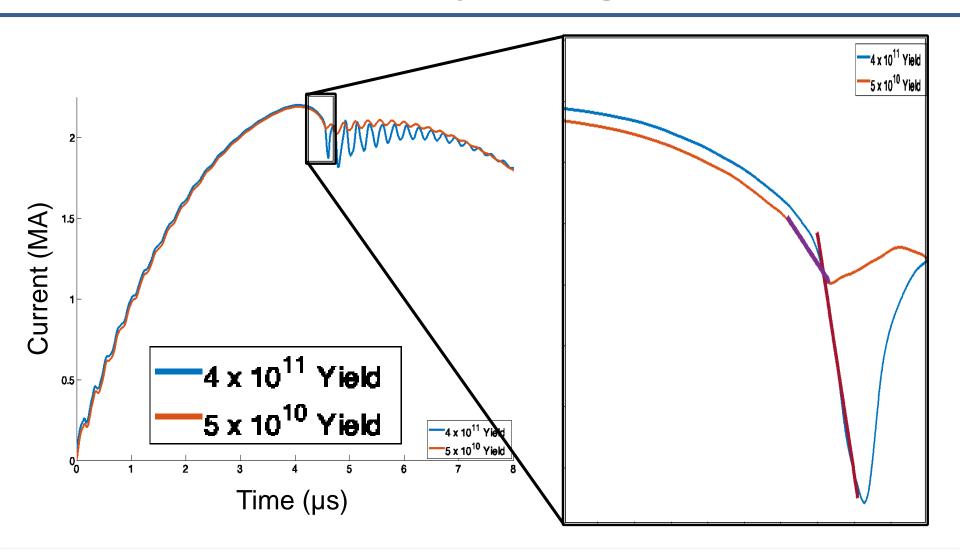


Trailing mass left behind by hydrodynamic instabilities create lower-inductance current pathways that divert current from the pinch region

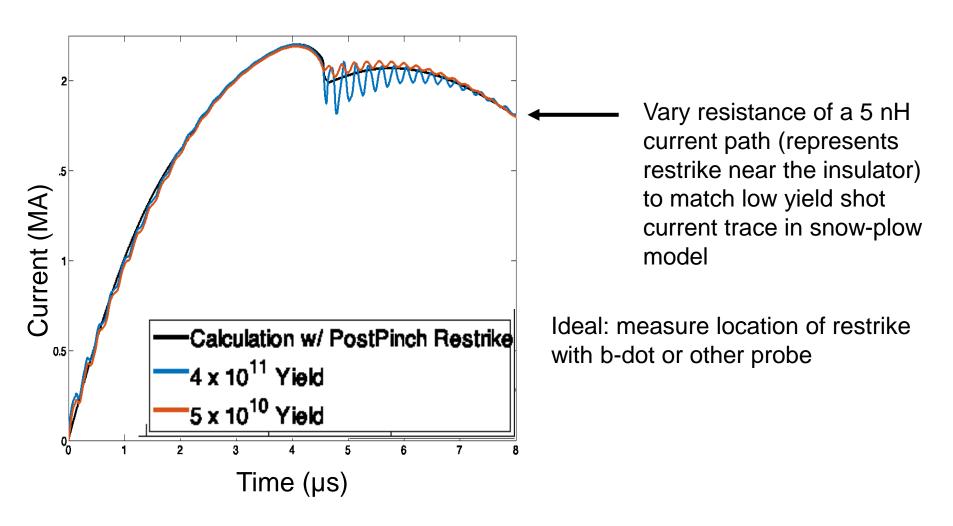
Hollow anode mitigates anode-to-sheath restrikes



Lower yield shots show evidence of parasitic current diversion from pinch region



Lower yield shots show evidence of parasitic current diversion from pinch region

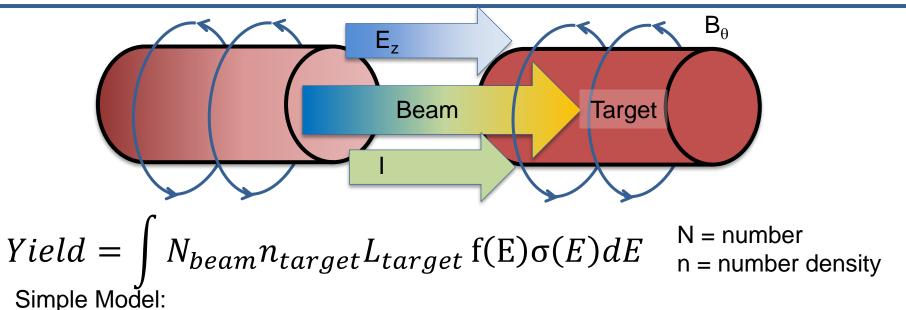


Insights from PIC modeling to be applied to lower order model

- Beam temperature for a given DPF current doesn't change much
- For a given stored energy class of DPFs, conversion efficiency of gun energy to beam energy is somewhat constant
 - → Main influence on yield is through aerial density and temperature of target
- Large radius anode/long implosion time leads to hydrodynamic instabilities that appear early in run-in
 - We can mitigate these hydro instabilities with a tapered anode, where the taper stops at a particular radius. Mass is swept up starting from the radius where the taper stops (the "implosion radius").
- Plasma target needs to be hot to minimize stopping power (increases aerial density average cross-section)
 - Too much mass in the implosion (high gas fill or large implosion radius) can cause target to be cold
- A hollow anode can help mitigate anode-to-sheath restrikes

From analytic shock physics: maximum achievable convergence ratio appears to be about 10 in a shock-driven cylindrical implosion

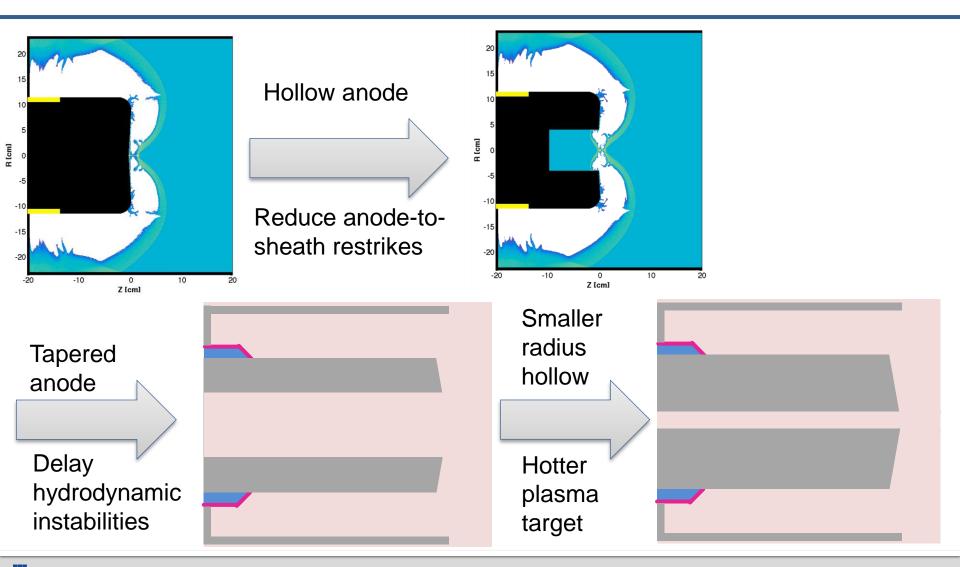
With a few assumptions, we can make a reduced order model to explore wide parameter space



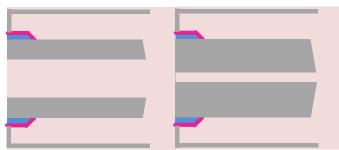
- Assume the hydrodynamic disassembly time of the "target" >> duration of ion beam and acceleration time of ion beam
- Assume beam can't miss the target, i.e. partially magnetized
 - Getting to sufficient areal density will probably mostly guarantee this
 - Larmor Radius for 1 MeV D+ near the pinch is about 0.5-1 mm
- Useful pinch length is ≈ 2 cm long
- Beam spectrum is decaying exponential

$$f(E) = e^{-E/E_b}$$

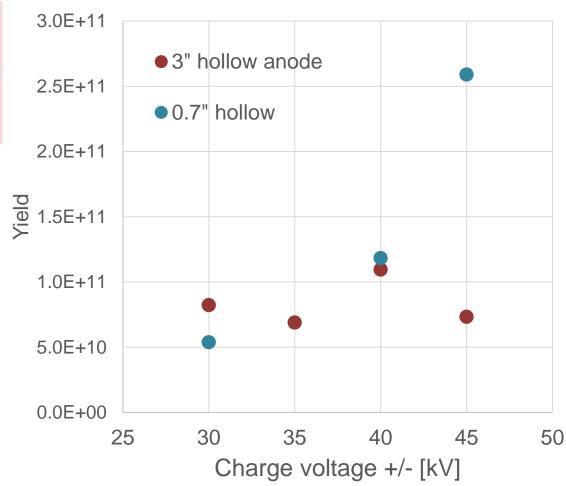
Anode design evolution influenced by both kinetic and reduced order models



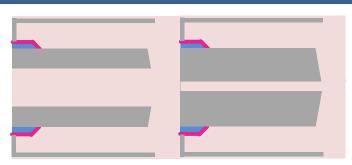
Example of how modeling insights helped improve anode design



- First anode fielded on MJOLNIR was unsuccessful at high current
- Modeling gave us insight that plasma target was not getting hot enough
- We reduced the hollow radius and recovered performance at high currents

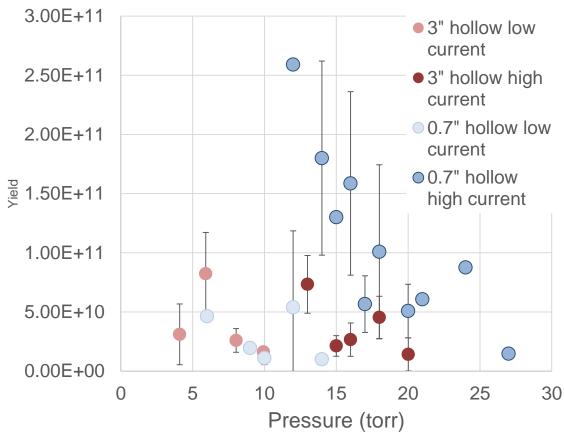


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Yields from first 2 anodes on MJOLNIR



Key future questions

- In MJOLNIR, where do restrikes occur?
 - At base of DPF/insulator
 - In A-K gap
 - From anode-to-sheath
- How do we avoid restrikes?
 - Operating pressure
 - Could be problematic since high pressure is needed for high yields
 - Increase A-K gap
 - Increases head inductance, lowers peak current what is effect on pinch current?
 - Anode shape
- What limits performance at higher pressures?





