

Non-invasive measurements of plasma parameters via optical emission spectroscopy

Amy Wendt

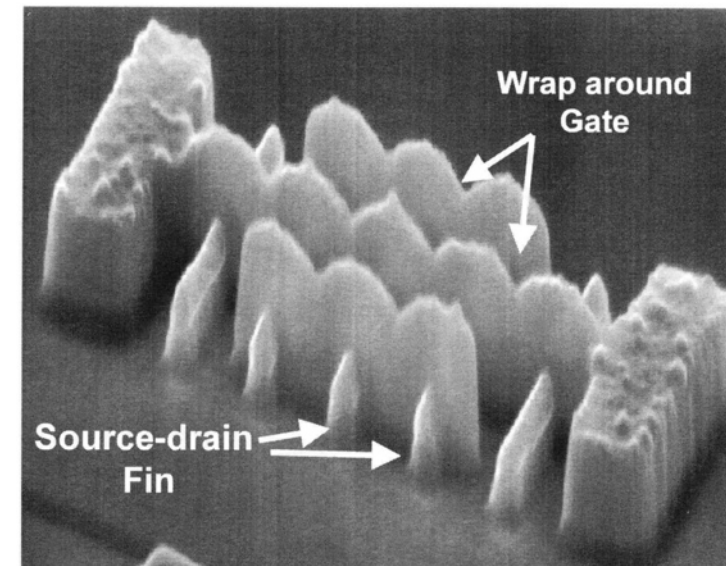
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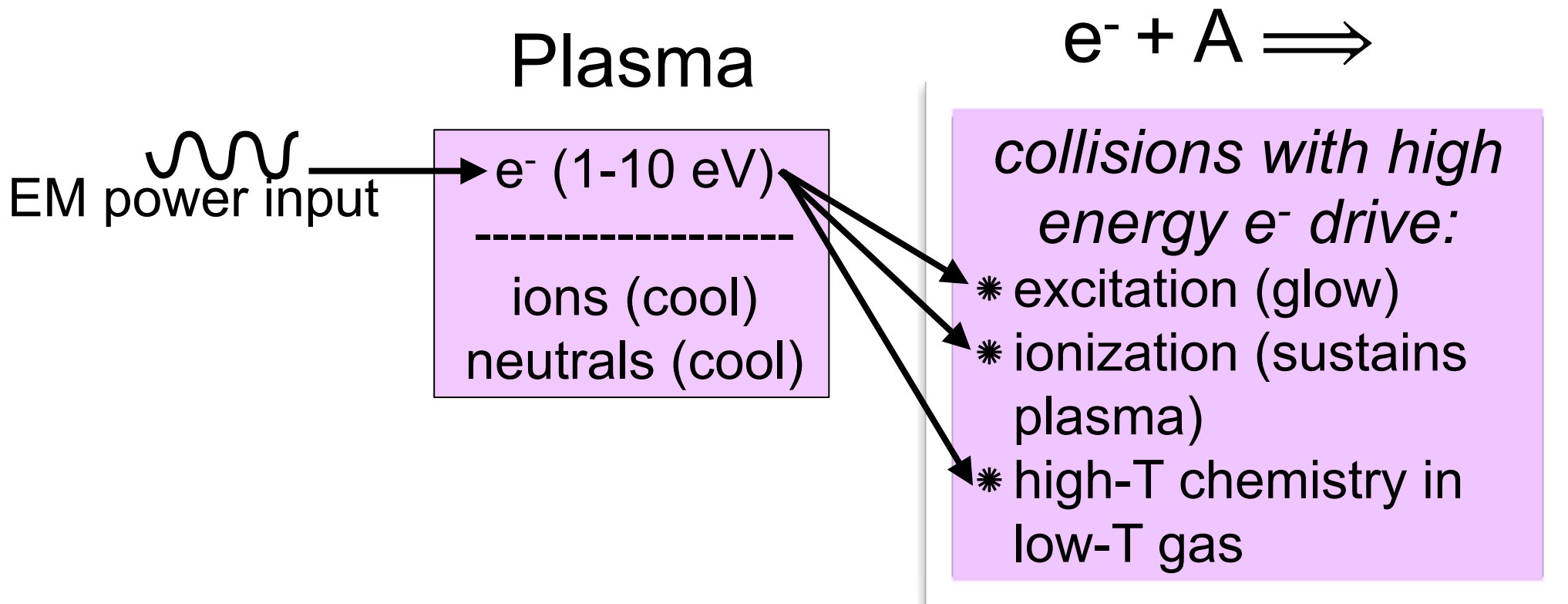
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 - Lauren Aneskavich, UW ECE Dept. (now at Rockwell Automation)
 - Dr. Svetlana Radovanov, Applied Materials Corp.
 - Dr. Harold Persing, Applied Materials Corp.
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Low-temperature plasma applications

- applications are diverse
- manipulating plasma is a means to control technological outcomes
- diagnostic tools are needed to improve plasma science understanding



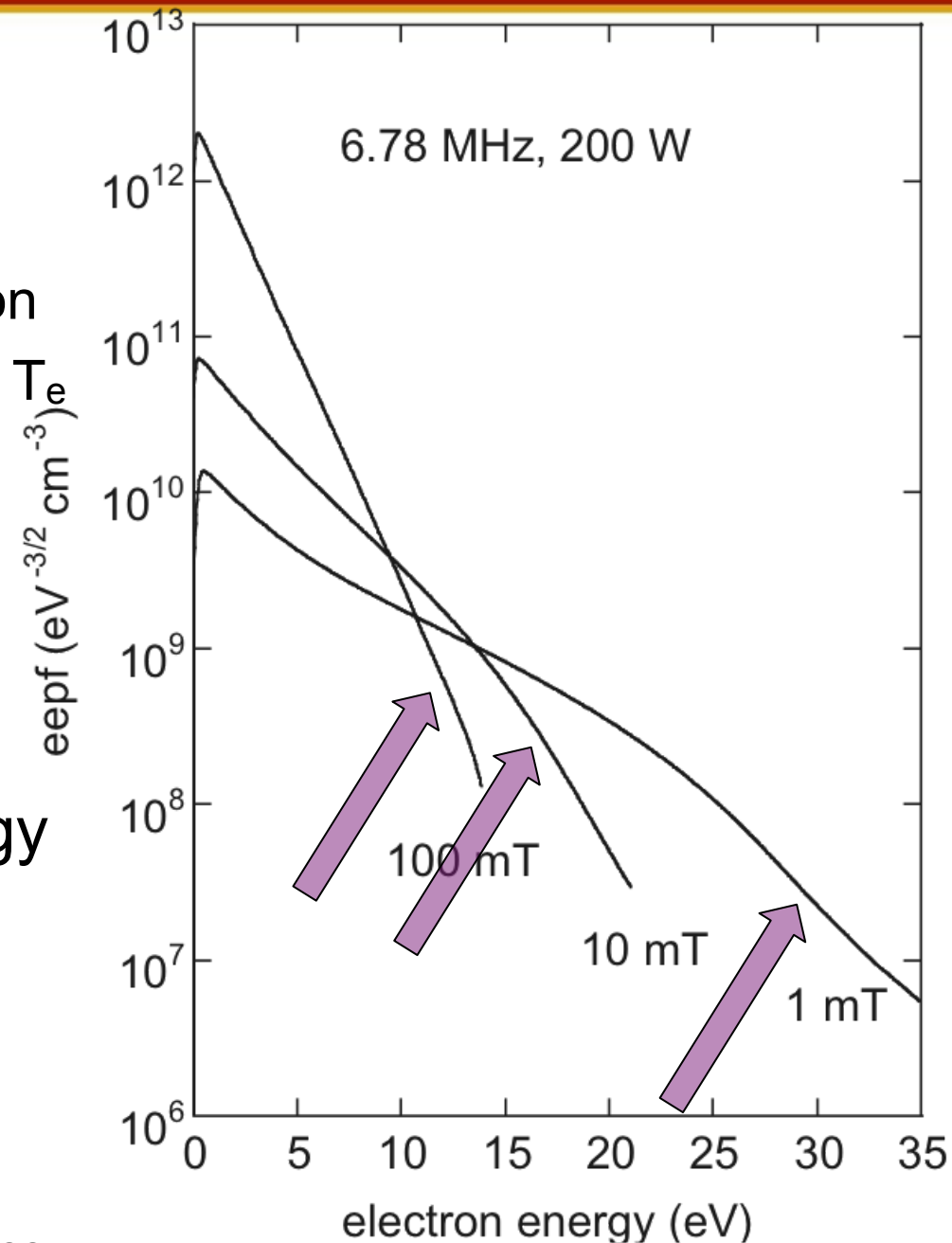
LTPs are non-equilibrium systems



- *energetic electrons are key link in process outcomes*

Electron energy distribution functions

- Maxwellian EEPF appears as straight line -
 - corresponds to thermodynamic equilibrium *within* electron population
 - characterized by single parameter - T_e
- RF inductively coupled plasmas: EEDFs with 'depleted' high energy 'tails'
 - attributed to:
 - ♦ inelastic collisions
 - ♦ escape to walls



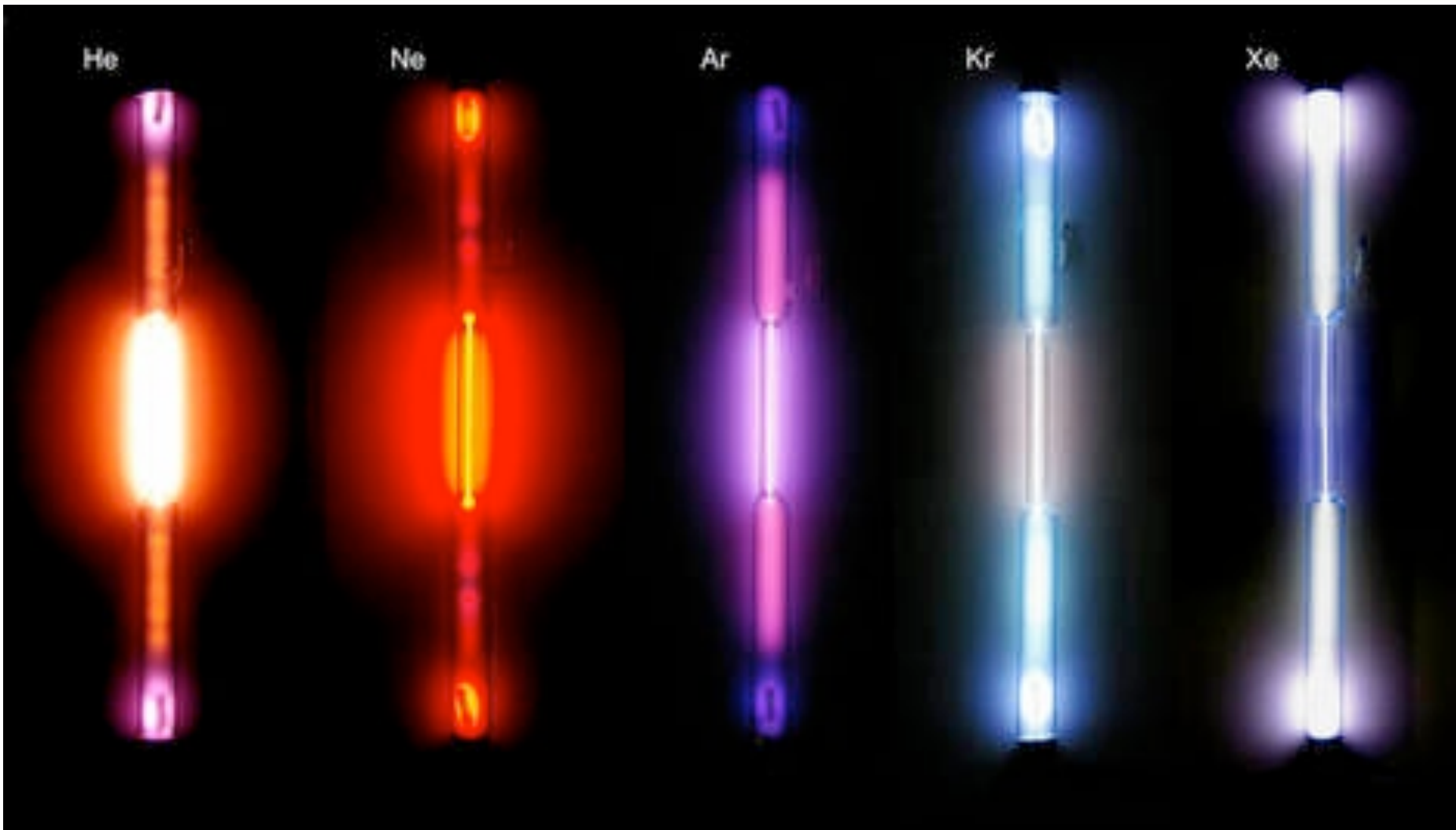
Godyak *et al.*, PSST 11, 525 (2002)

Presentation Outline: Optical diagnostics for plasma characterization



- Optical emission spectroscopy of argon-containing plasmas
 - emission model for determination of
 - ♦ EEDF, T_{eff}
 - ♦ metastable and resonance level densities
 - ♦ electron density - n_e
 - ♦ Example: non-Maxwellian EEDFs in inductively coupled plasma
- Real-time plasma monitor
 - metastable and resonance level densities - n_m and n_r
 - effective electron temperature - T_{eff}
 - electron density - n_e
- VUV detection
 - resonance level densities - n_r - as a proxy for Ar VUV flux

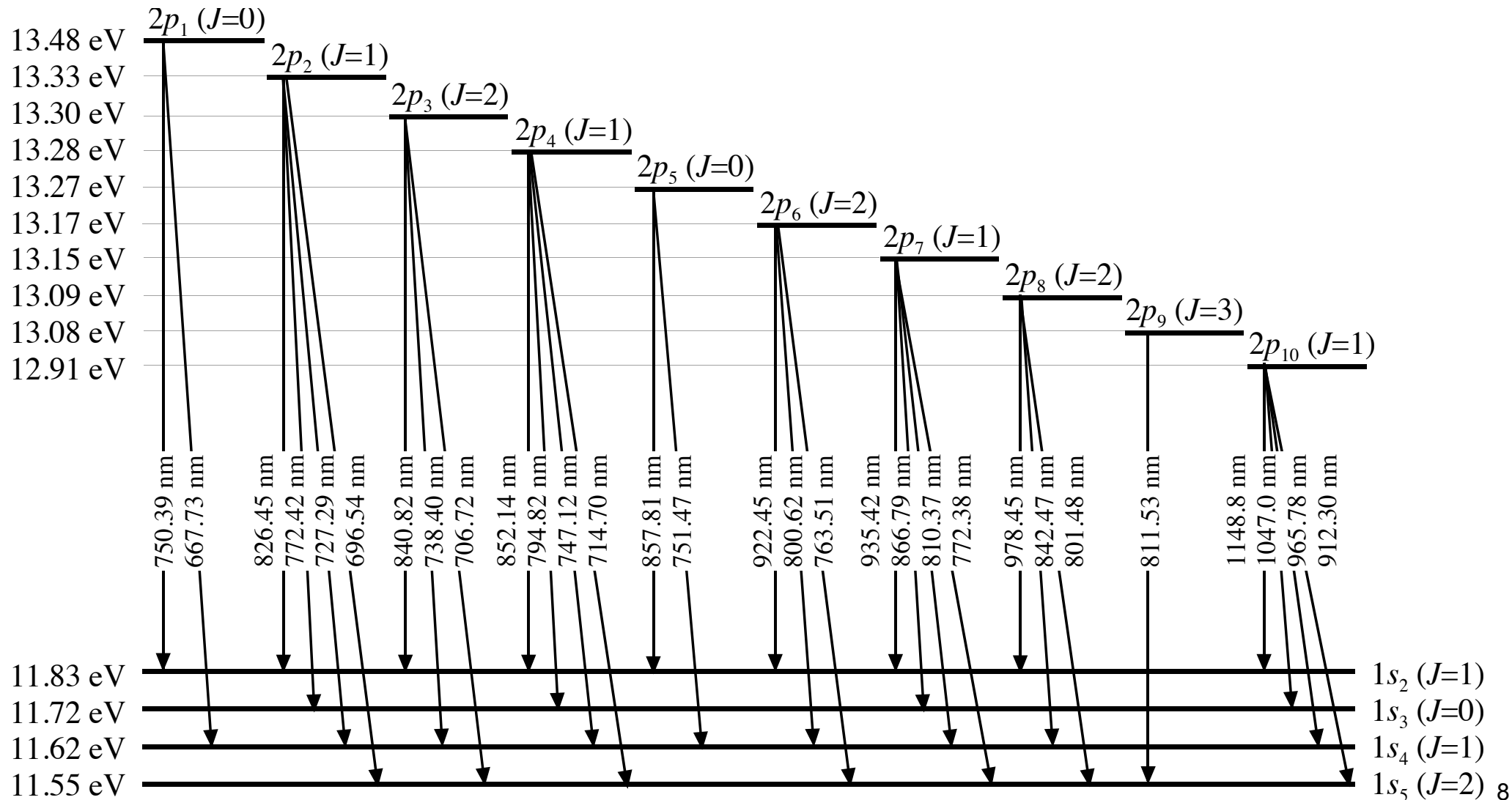
Each plasma emits distinctive spectrum



- http://wapedia.mobi/en/Gas_filled_tube - different rare gases
- Also, plasmas with the *same* gas changes colors with conditions

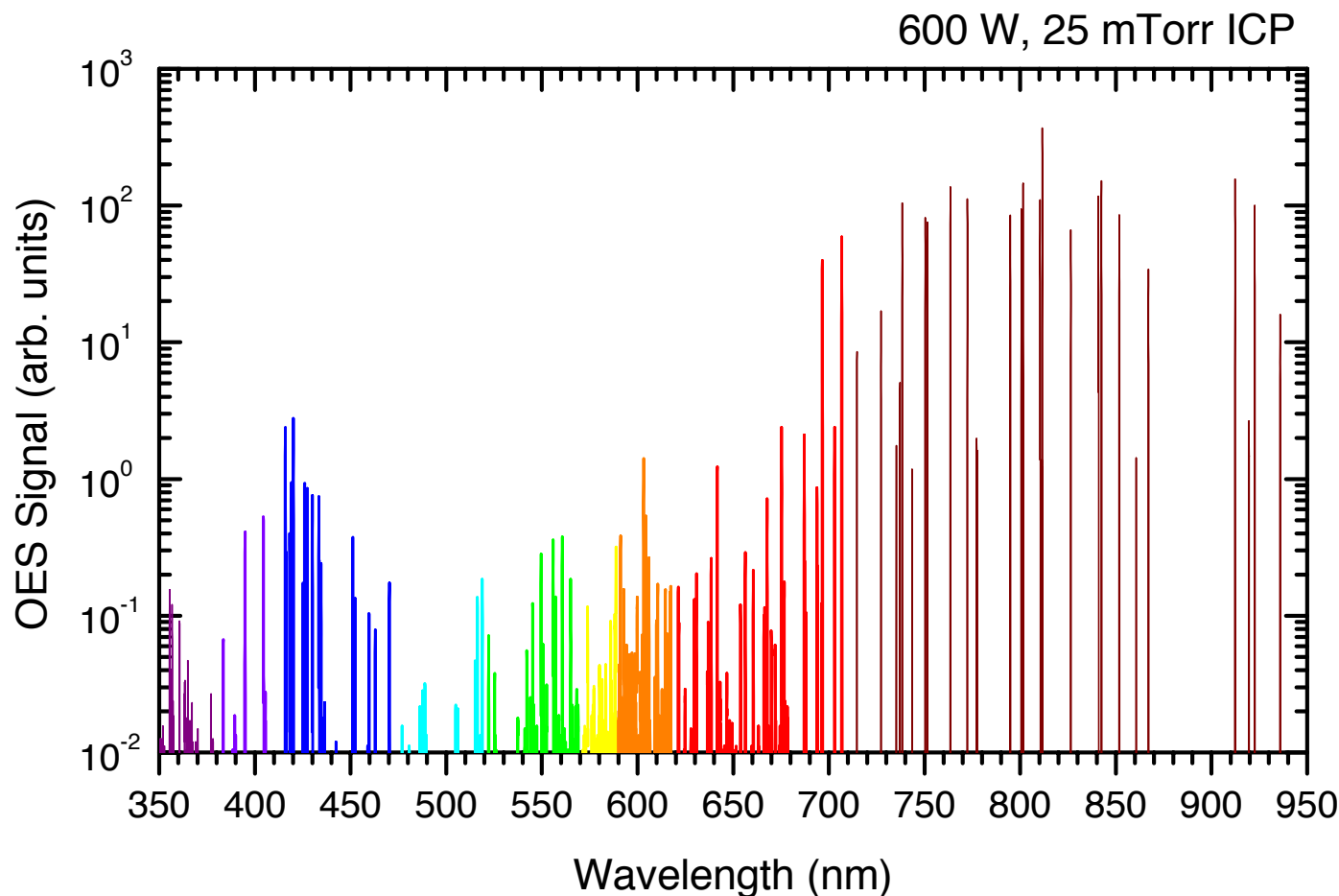
Partial Ar energy level diagram

- Allowed transitions resulting in photon emission



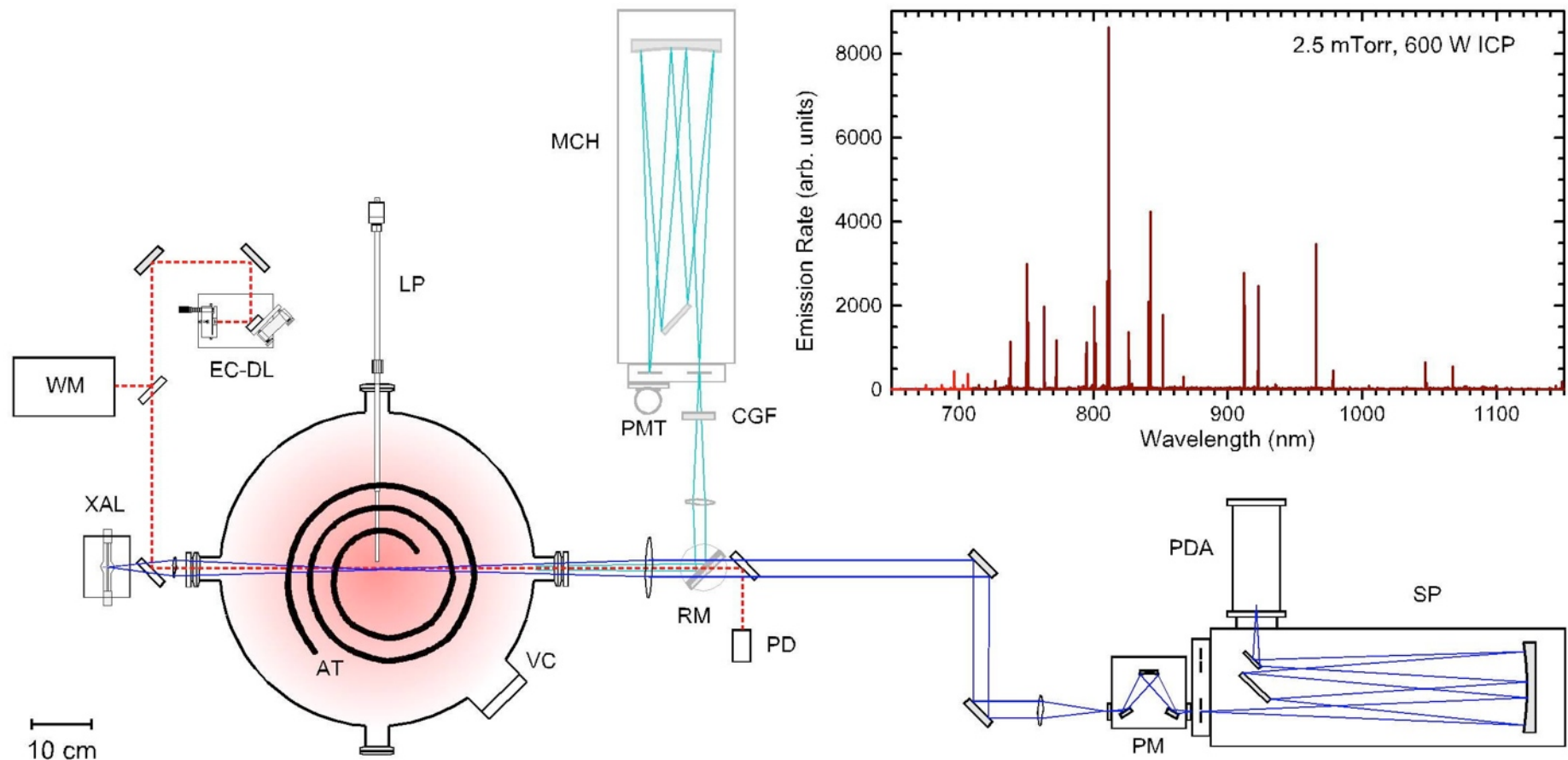
Argon emission spectrum

- Ultimate Goal: use emission spectra to extract encoded information about plasma conditions



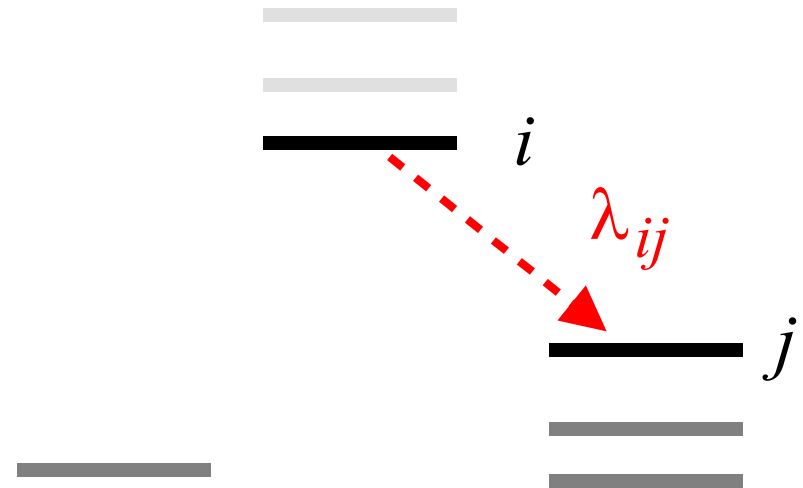
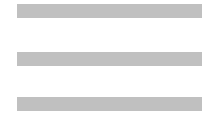
Inductively coupled plasma (ICP)

- Operates in Ar and Ar mixtures at 13.56 MHz
- Electron energy distribution functions (EEDF)
 - Langmuir probes
 - optical diagnostics



Determining plasma properties from spectrum: emission model

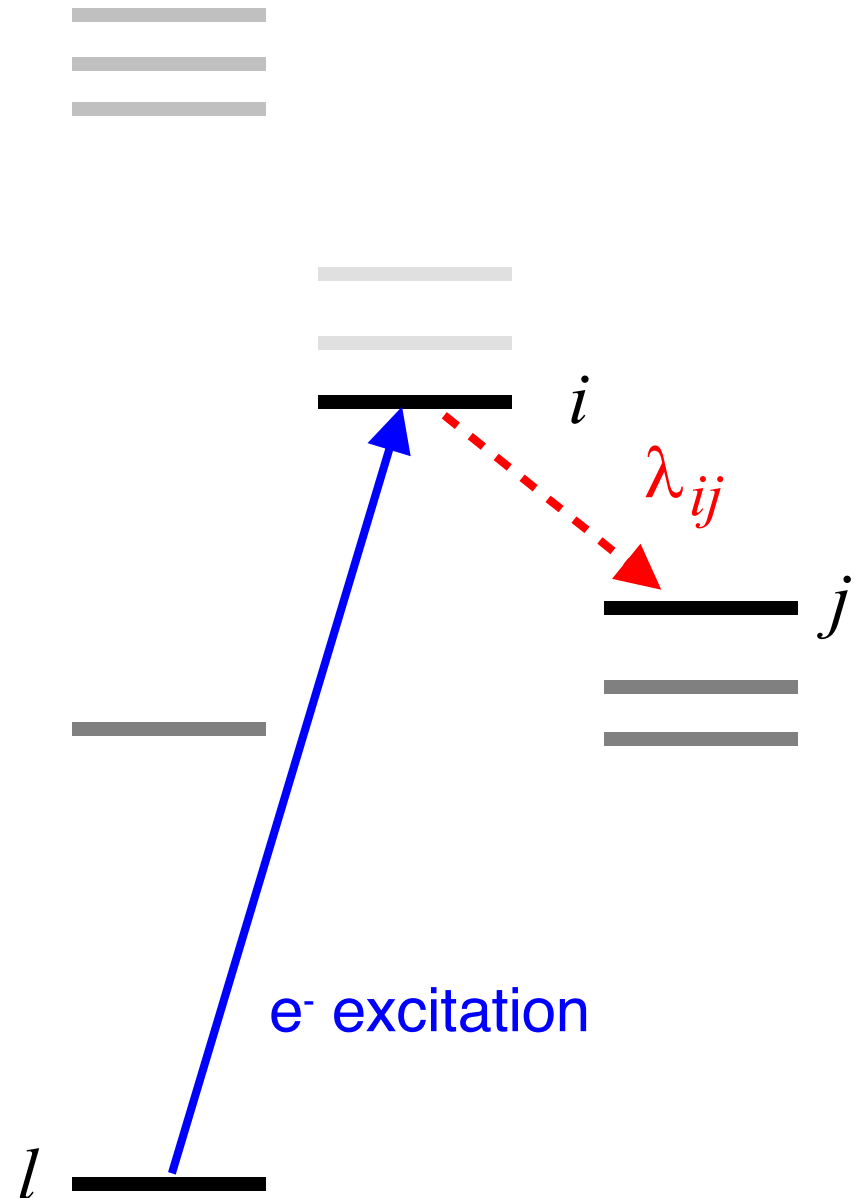
- Goal:
 - predict photon emission rate Φ_{ij} for $i \rightarrow j$
- Approach:
 - account for mechanisms that populate emitting state - level i



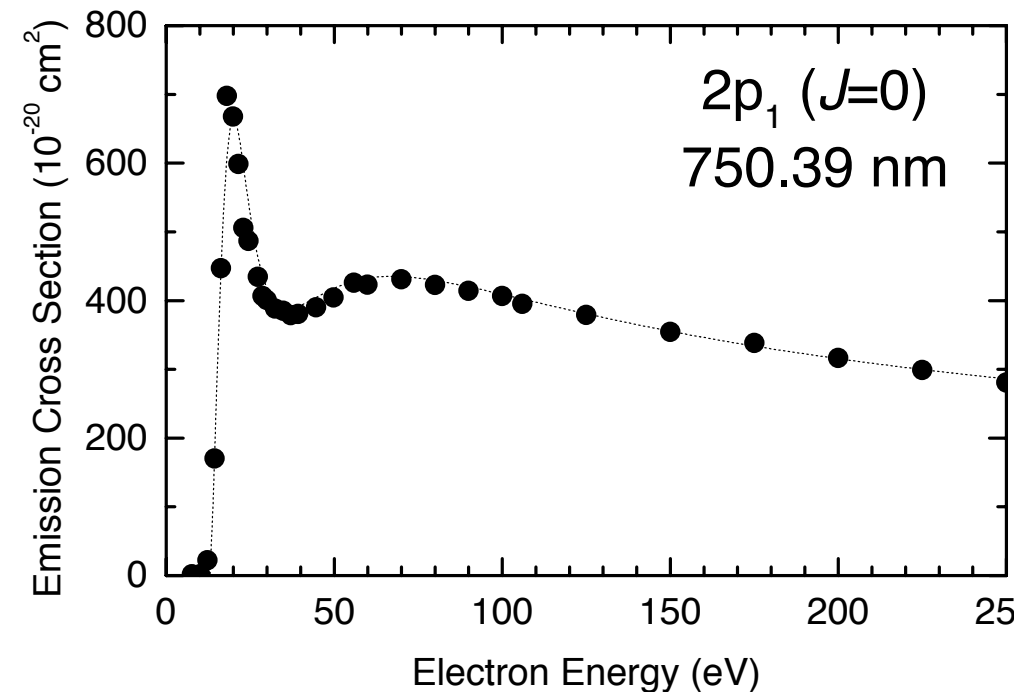
l —

Emission model

- Goal:
 - predict photon emission rate Φ_{ij} for $i \rightarrow j$
- Include:
 - electron excitation from level l
 - ♦ excitation probability is a function of electron energy
 - ♦ measured “cross sections” used to characterize excitation probabilities

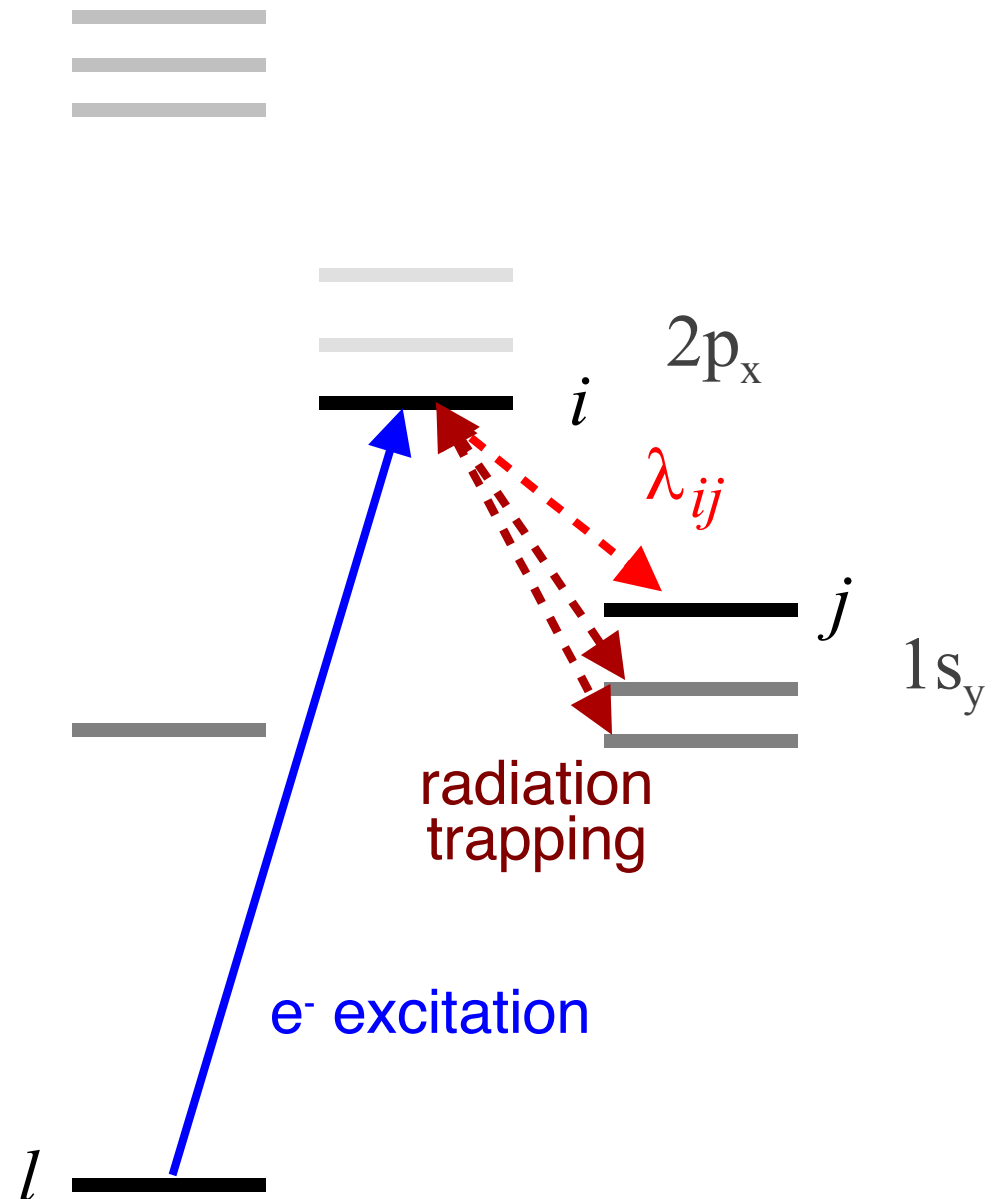


- cross section vs. electron energy
 - measure of electron impact excitation probability
 - key link between emission intensity and EEDF
 - measured by Prof. Chun Lin's group:
 - ♦ excitation from ground state
 - *Phys. Rev. A* **57**, 267 (1998)
 - *Phys. Rev. A* **68**, 032719 (2003)
 - *Ar. Data Nucl. Data Tables* **93**, 831 (2007)
 - ♦ excitation from metastable levels
 - *Phys. Rev. Lett.* **81**, 309 (1998)
 - *Phys. Rev. A* **59**, 2749 (1999)
 - *Phys. Rev. A* **75**, 052707 (2007)



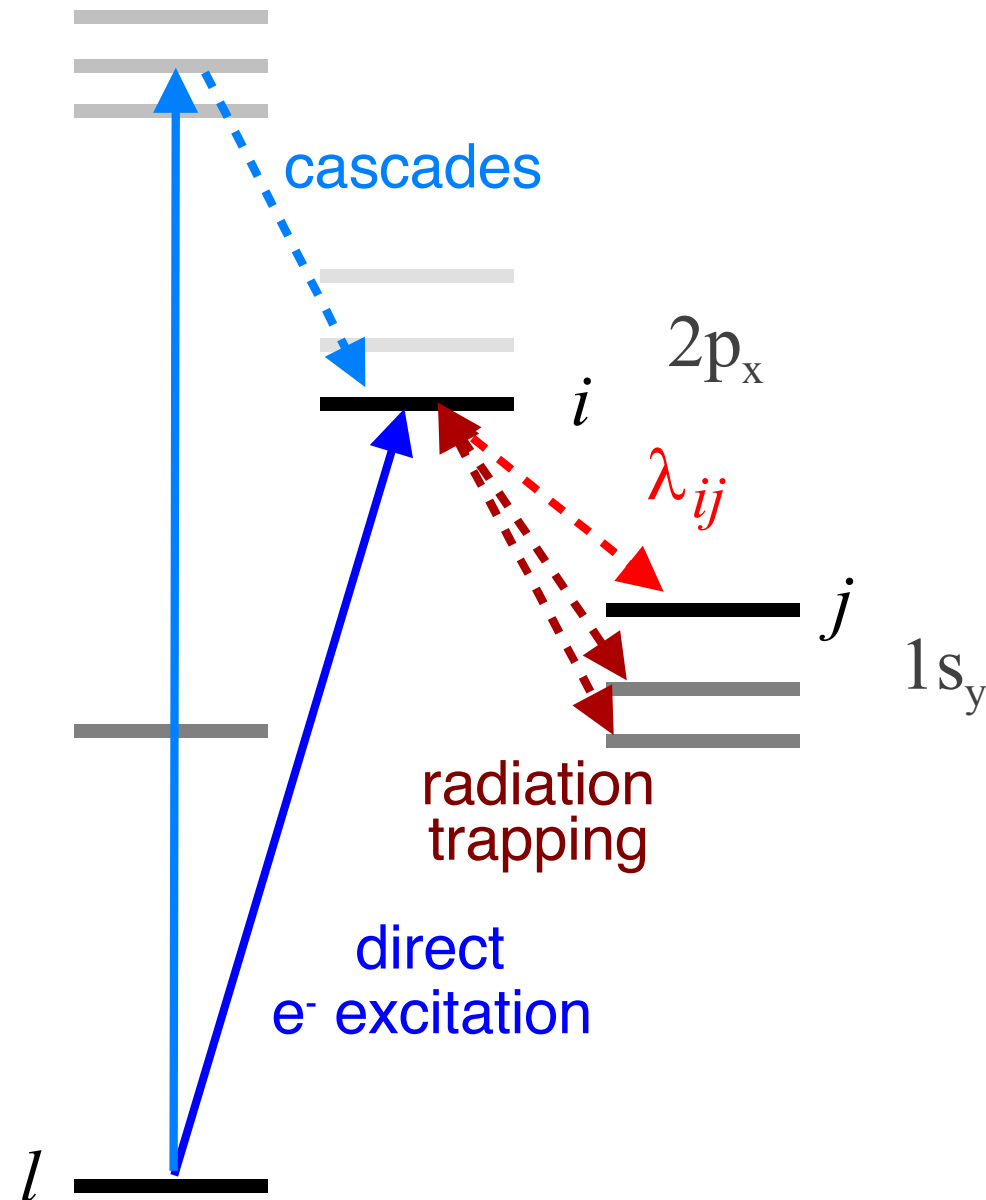
Emission model - reabsorption

- Goal:
 - predict photon emission rate Φ_{ij} for $i \rightarrow j$
- Include mechanisms for populating level i :
 - electron excitation from level l
 - radiation trapping: excitation of metastables by photon absorption

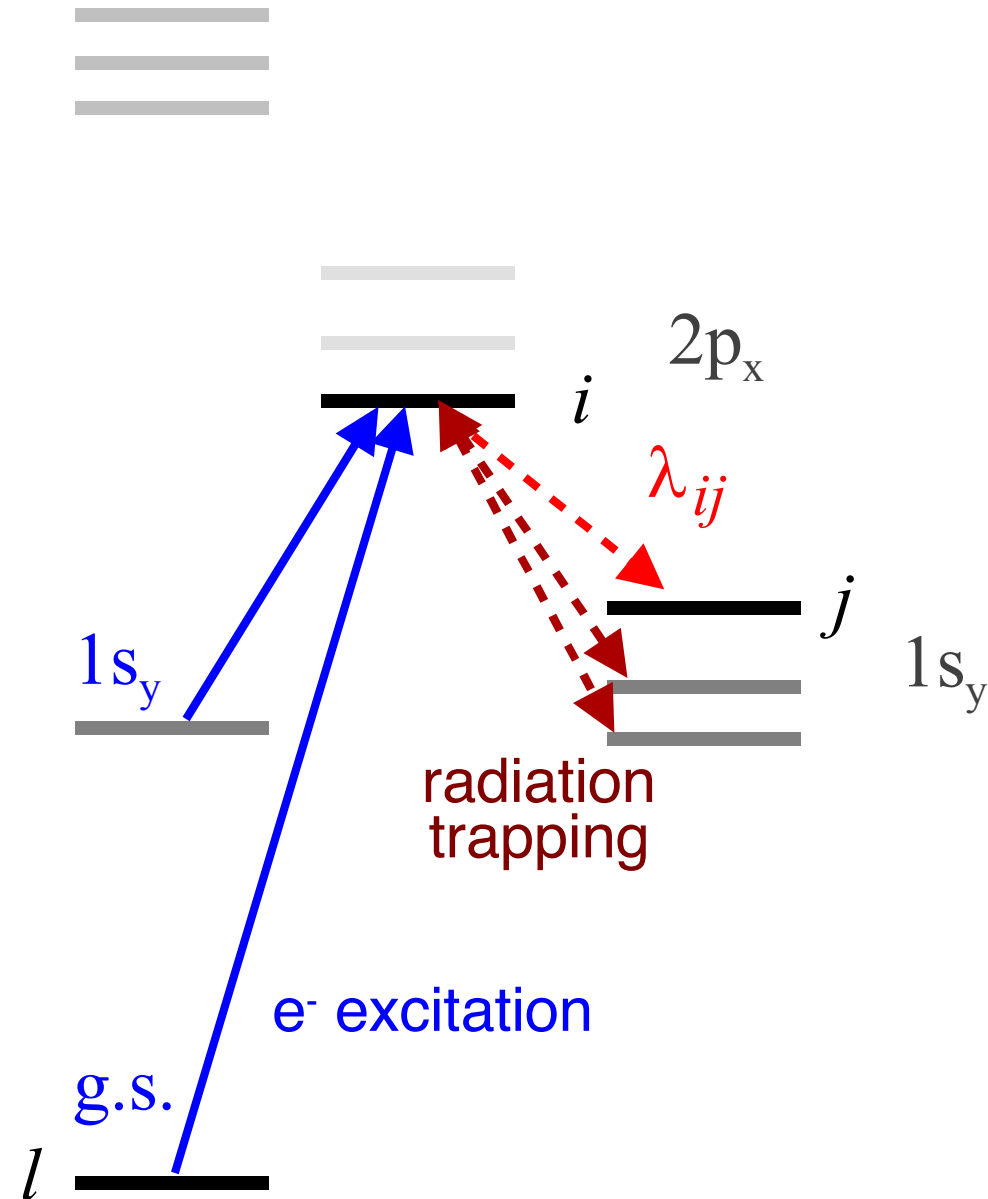


Emission model - cascades

- Goal:
 - predict photon emission rate Φ_{ij} for $i \rightarrow j$
- Include mechanisms for populating level i :
 - indirect electron excitation from level l
 - cascades from higher lying states
 - measured cross sections include cascades

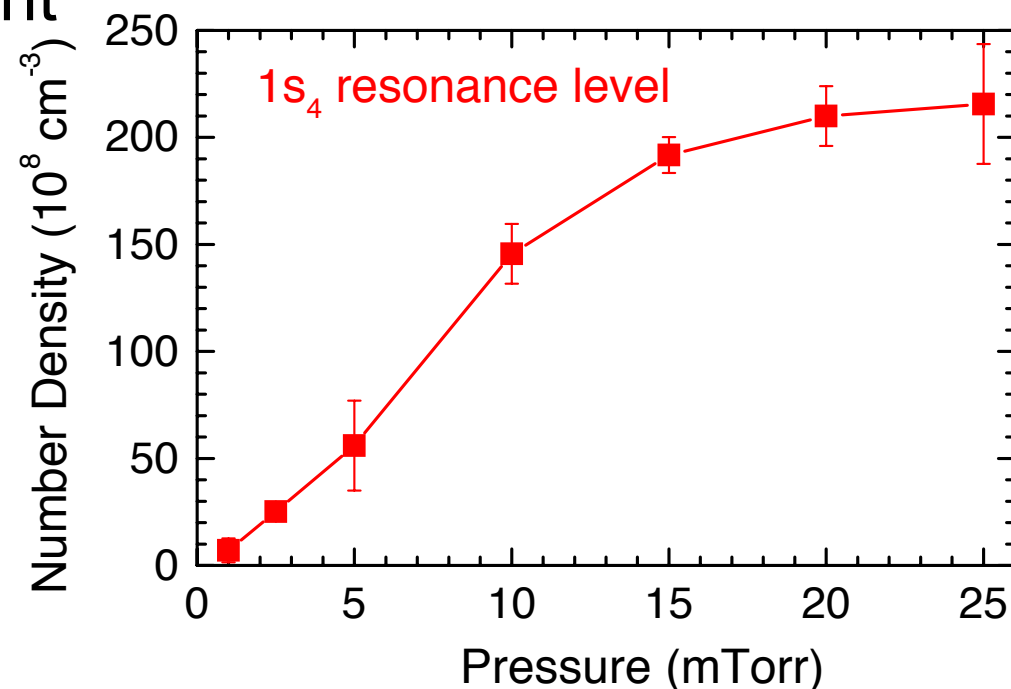


- Goal:
 - predict photon emission rate Φ_{ij} for $i \rightarrow j$
- Include mechanisms for populating level i :
 - electron excitation:
 - ♦ from ground state
 - ♦ from metastable/resonance levels
 - photon reabsorption/trapping
 - ♦ from metastable/resonance levels



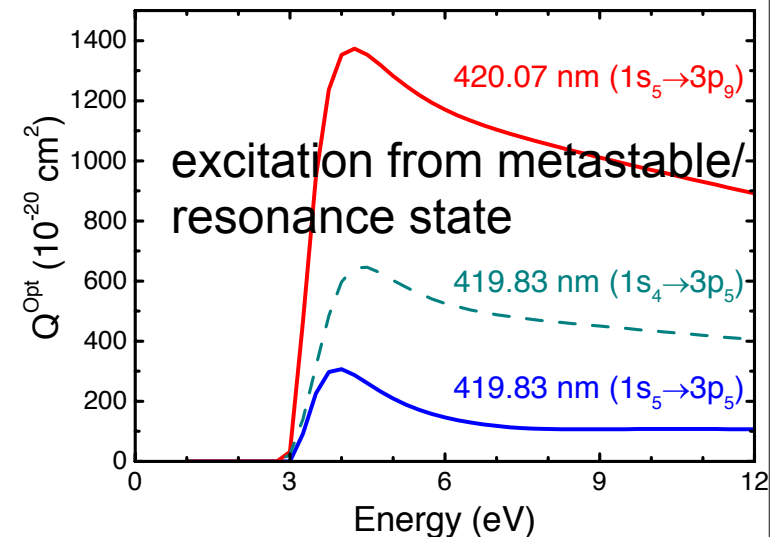
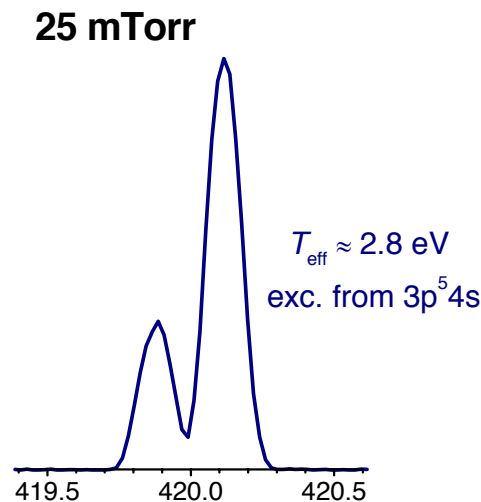
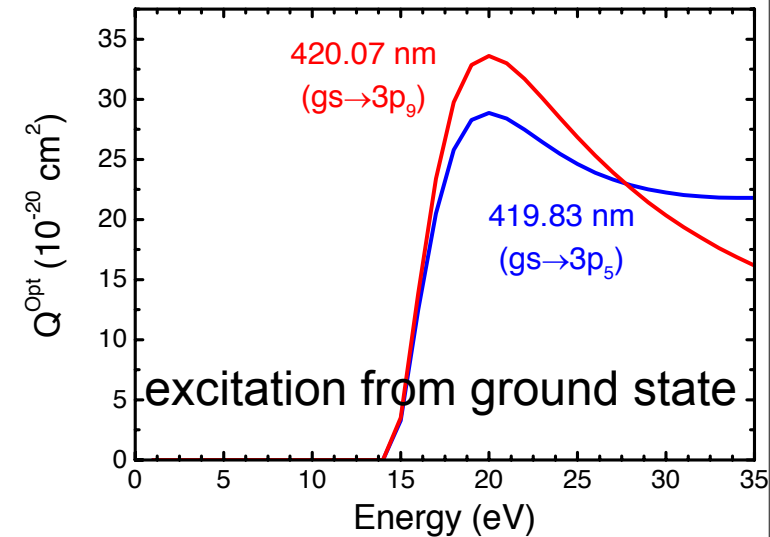
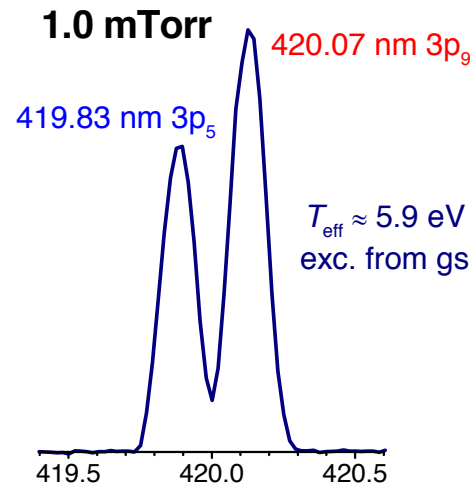
Excitation from resonance levels $1s_2$ & $1s_4$

- $1s_3$ and $1s_5$ are *metastables*:
 - no radiative losses
 - high concentrations possible
- $1s_2$ and $1s_4$ populations significant when radiation is trapped:
 - $1s_2$ and $1s_4$ resonance levels:
 - ♦ have radiative decay channel to ground state
 - ♦ generally lower concentrations
 - ♦ exception: reabsorption when ground state density is high
 - *Electron impact excitation from resonance levels significant*
 - ♦ cross sections not measured
 - ♦ estimates used
 - ♦ *Adv. At. Mol. Opt. Phys.* **54**, 319 (2006)



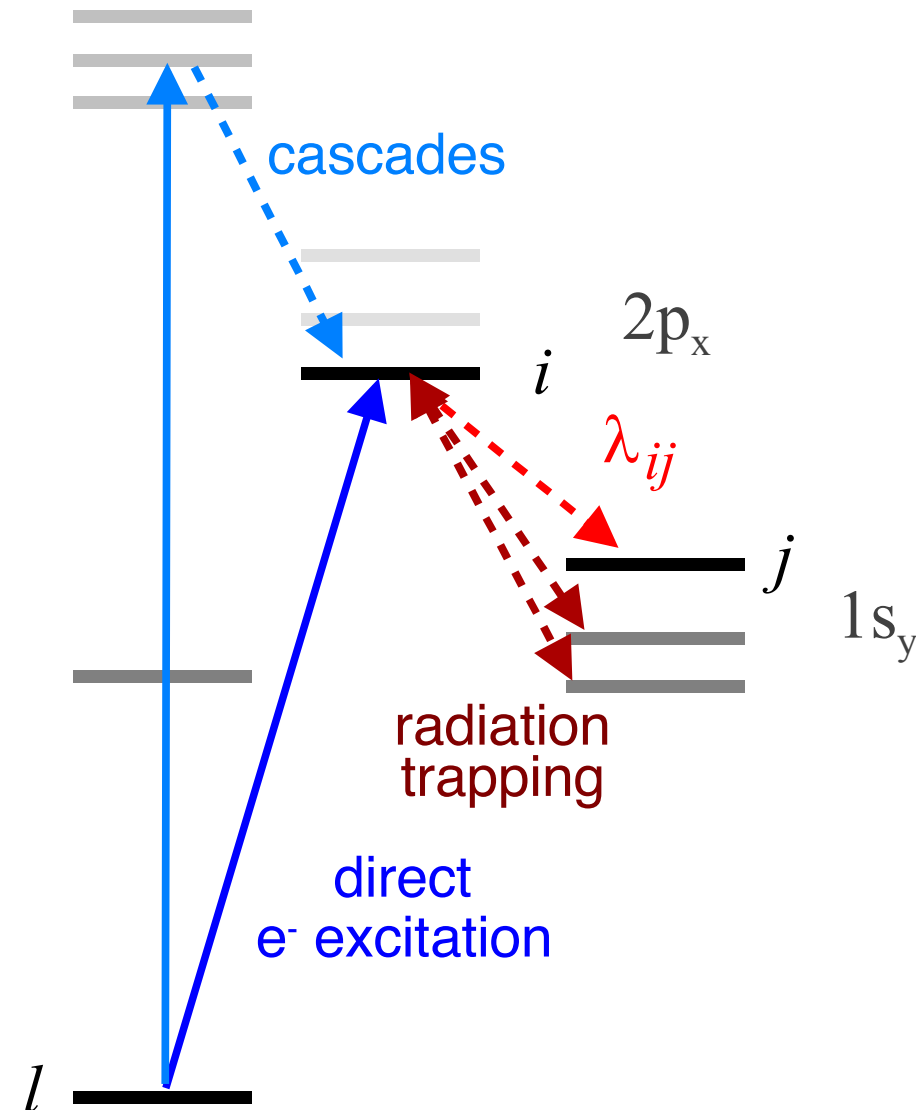
Cross sections: variations in threshold and energy dependence used in model

- 419.8 and 420.1 nm - close in wavelength
- difference between ground state and metastable cross sections
- large difference in peak intensities at 1 mTorr vs. 25 mTorr
 - higher ground state contribution at higher T_e



Emission model: summary

- Model goal:
 - predict photon emission rate Φ_{ij} for $i \rightarrow j$
- Effects included:
 - electron excitation from level l
 - ♦ cascades
 - ♦ metastable contribution
 - radiation trapping
- Needed inputs
 - atomic data (A_{ij} , cross sections,...)
 - electron energy distribution function
 - number densities (ground state AND $1s_y$ for argon)



Emission model: intensity of $i \rightarrow j$ transition

$$\Phi_{ij} = K n_e R_{ij} \left[\sum_l n_l k_{ij}^l \right]$$

- sum over initial level l (ground state, $1s_2$, $1s_3$, $1s_4$, $1s_5$)
- k_{ij}^l is the electron impact excitation rate
 - depends on cross sections, EEDF

$$k_{ij}^l = \sqrt{\frac{2}{m_e}} \int_0^\infty Q_{ij}^l(E) f_e(E) \sqrt{E} dE$$

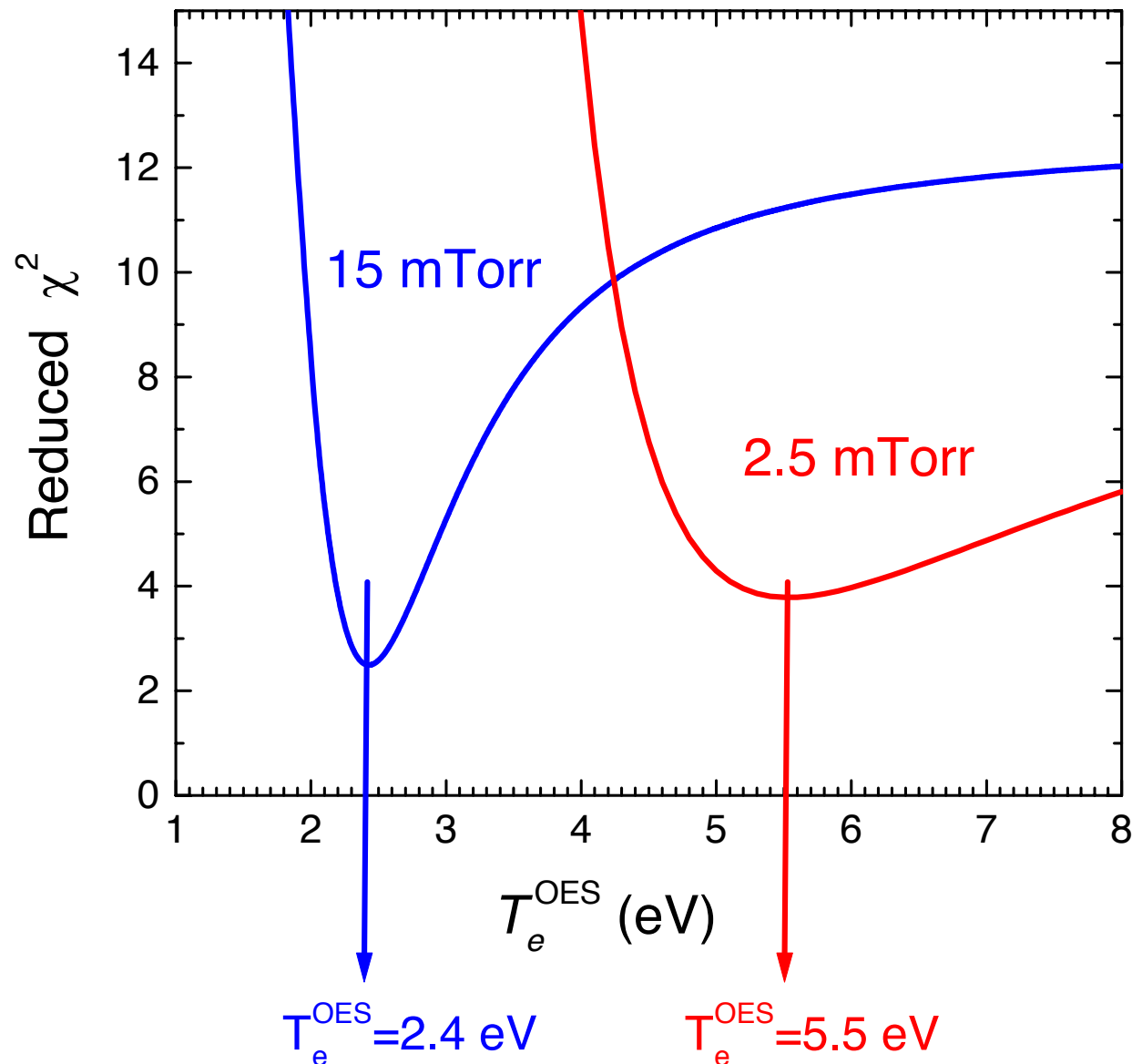
- reabsorption correction factor R_{ij} is function of $1s_y$ concentrations
- use ratios to standard line to eliminate K , n_e

$$\frac{\Phi_{ij}}{\Phi_{i'j'}} = \frac{R_{ij} \sum_l n_l k_{ij}^l}{R_{i'j'} \sum_l n_l k_{i'j'}^l}$$

- calculate line ratios for multiple trial EEDFs
- seek EEDF giving best match to observed line ratios

Trials for Maxwellian EEDF

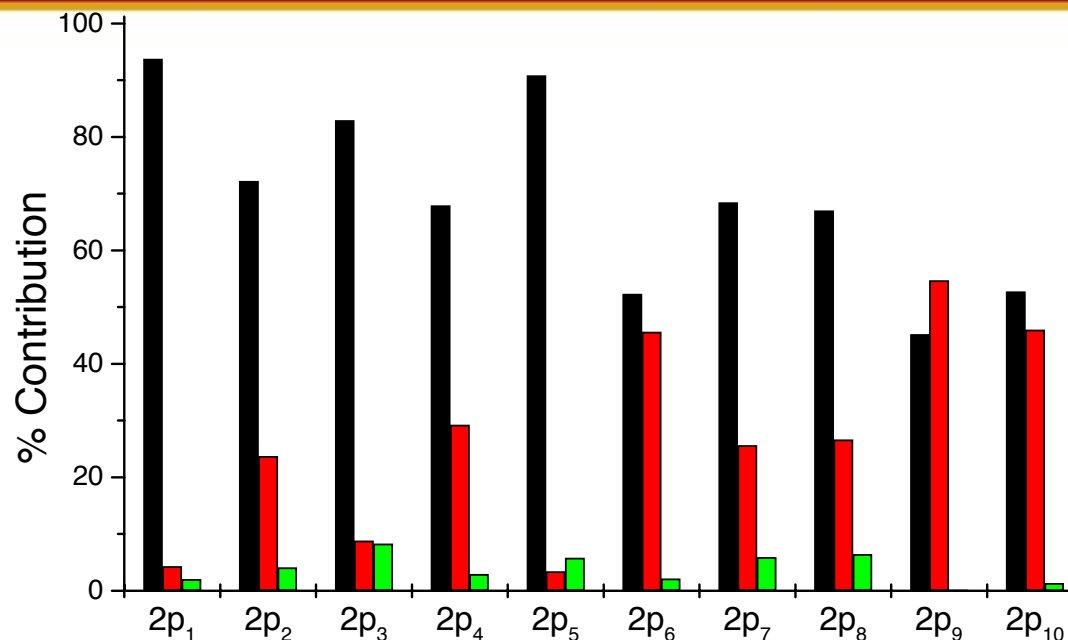
- minimum χ^2 gives best fit electron temperature



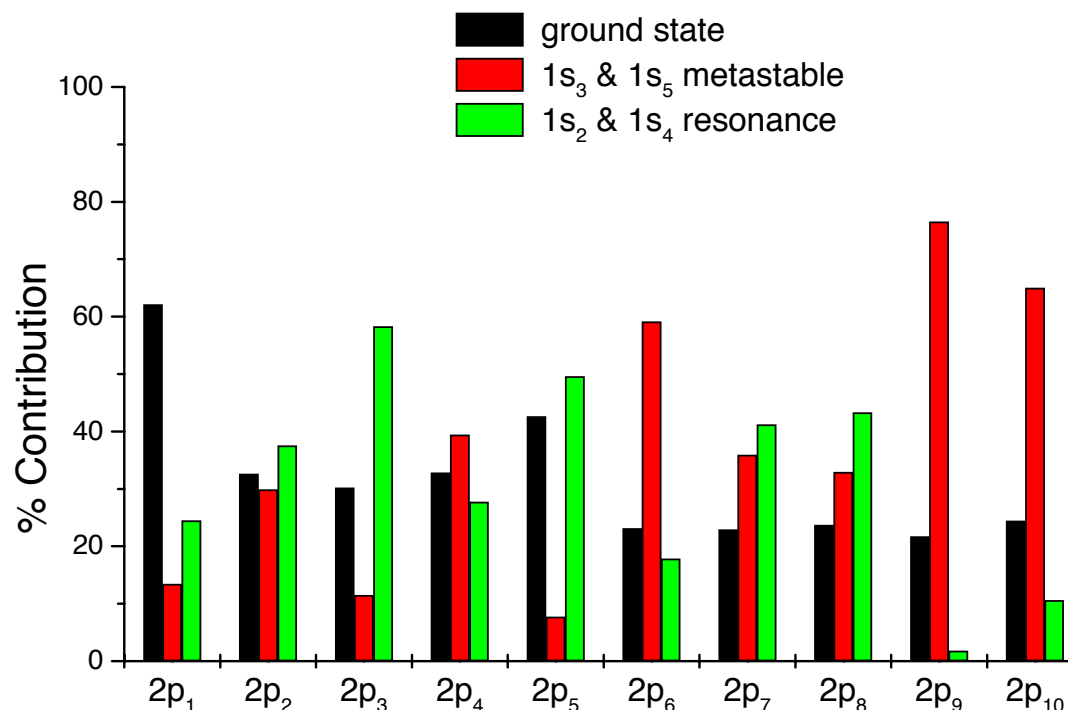
Emission model predictions

- contributions from ground state vs. $1s_y$ levels

2.5 mTorr
 $T_{\text{eff}} = 5.4 \text{ eV}$
 $n_m/n_0 = 1.4 \times 10^{-3}$



15 mTorr
 $T_{\text{eff}} = 2.7 \text{ eV}$
 $n_m/n_0 = 3.2 \times 10^{-4}$

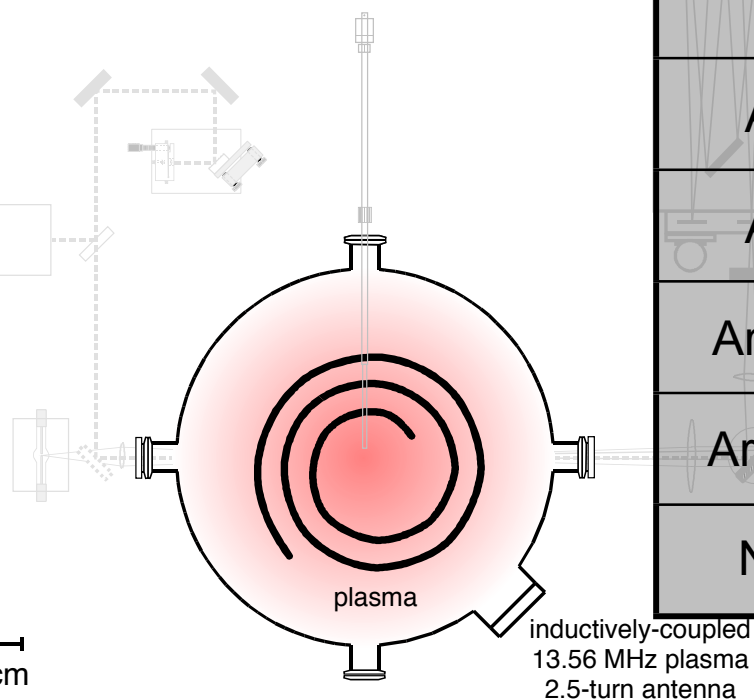


Application: Non-maxwellian EEDFs in ICPs

Inductively coupled plasma source

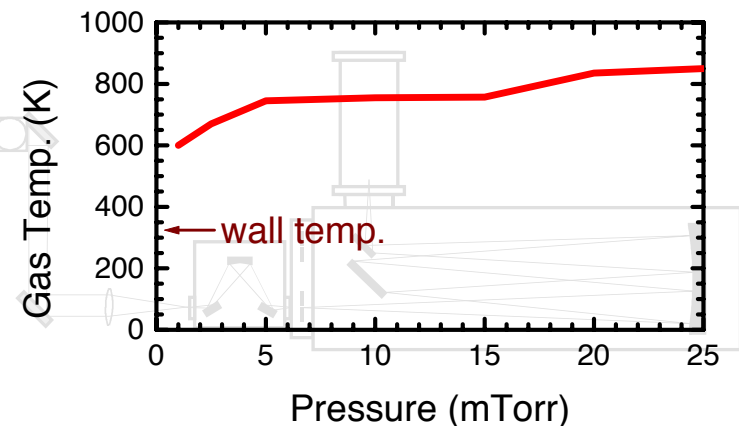
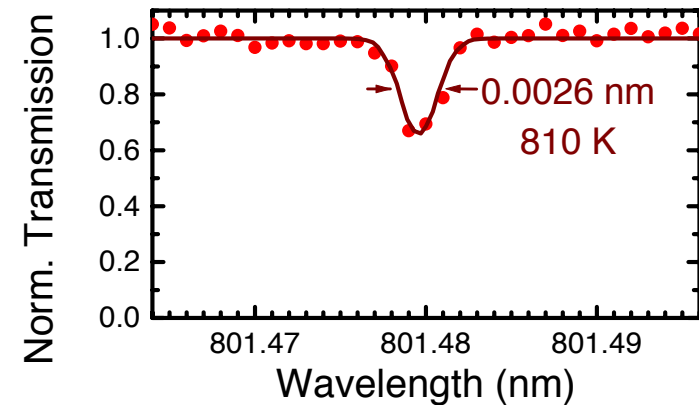
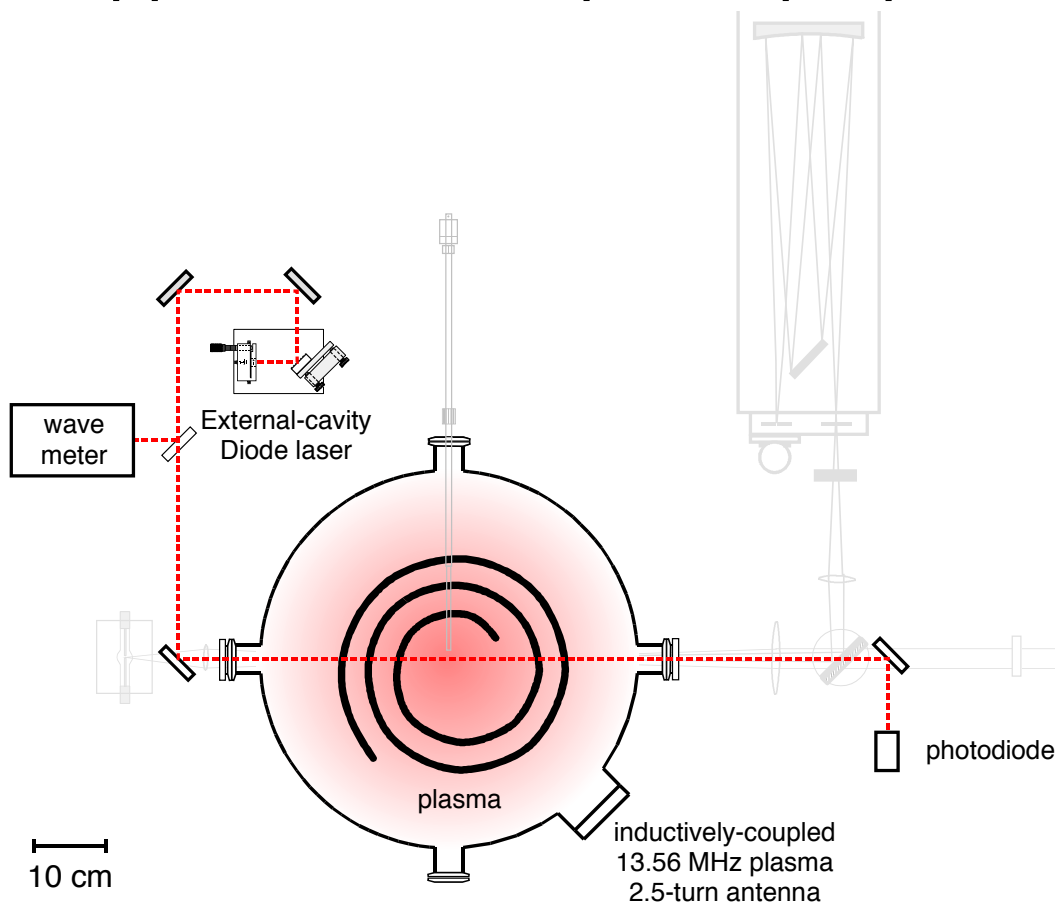
- 2.5 turn flat coil antenna
- 13.56 MHz
- 50 cm diameter chamber
- measure emission model inputs
- compare EEDFs with independent measurements

gas	variable	pressure (mTorr)	Power (W)	n_e (cm ⁻³)	T_e (eV)
Ar	pressure	1-50	600	10^{10} - 10^{12}	2-6
Ar	power	2.5, 15	20-1000	10^9 - 10^{12}	5, 3
Ar/N ₂	0-86% N ₂	2.5, 15	600	10^{10} - 10^{11}	5, 3
Ar/Ne	1-40% Ar	10	600	10^{10} - 10^{11}	3-6
Ne	pressure	5-25	600	1 - 3×10^{10}	5-7



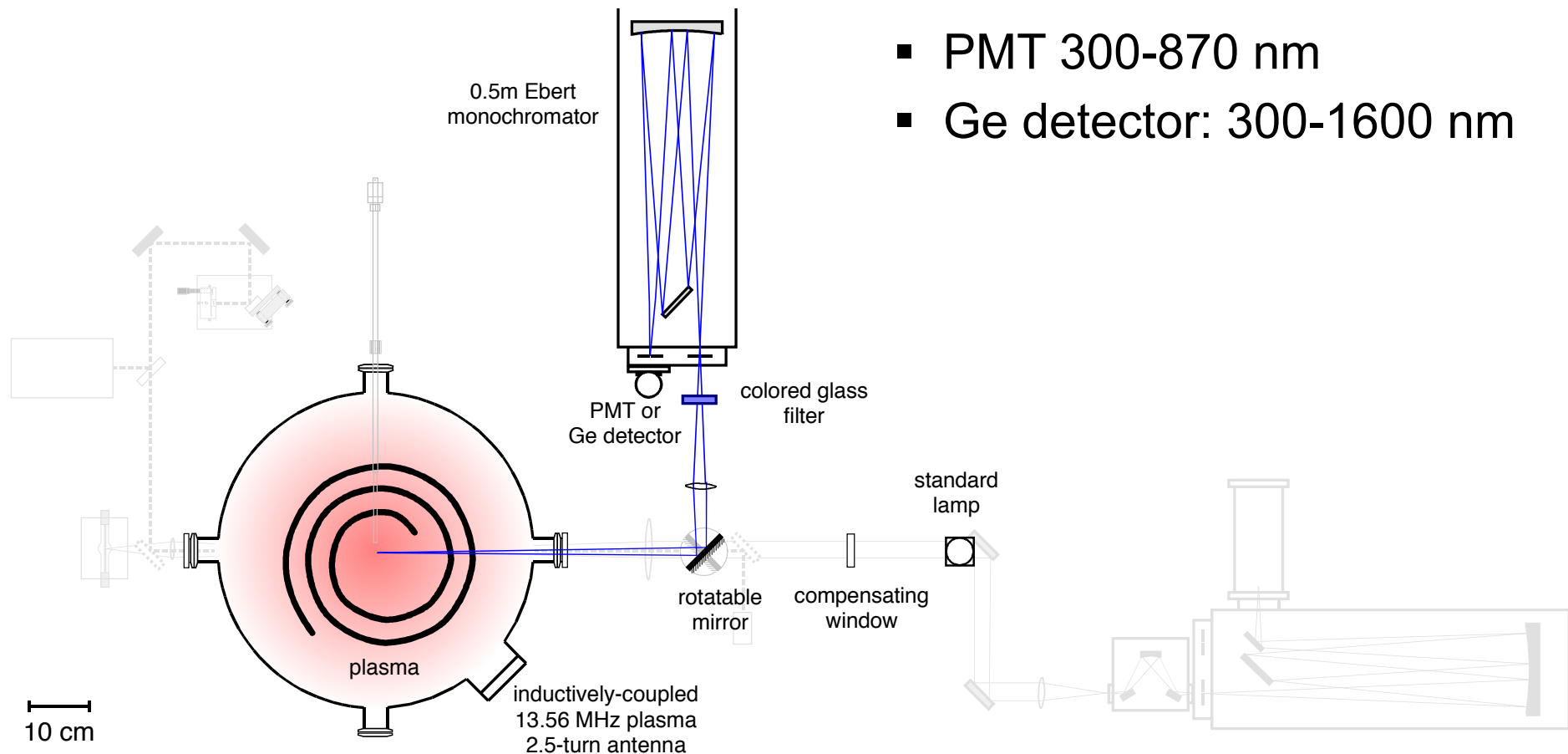
Model input: ground state density

- $n_0 = p/kT_{\text{gas}}$
- gas temperature determined with laser absorption spectroscopy
- Doppler width of $3p^54s-3p^54p$ transition line



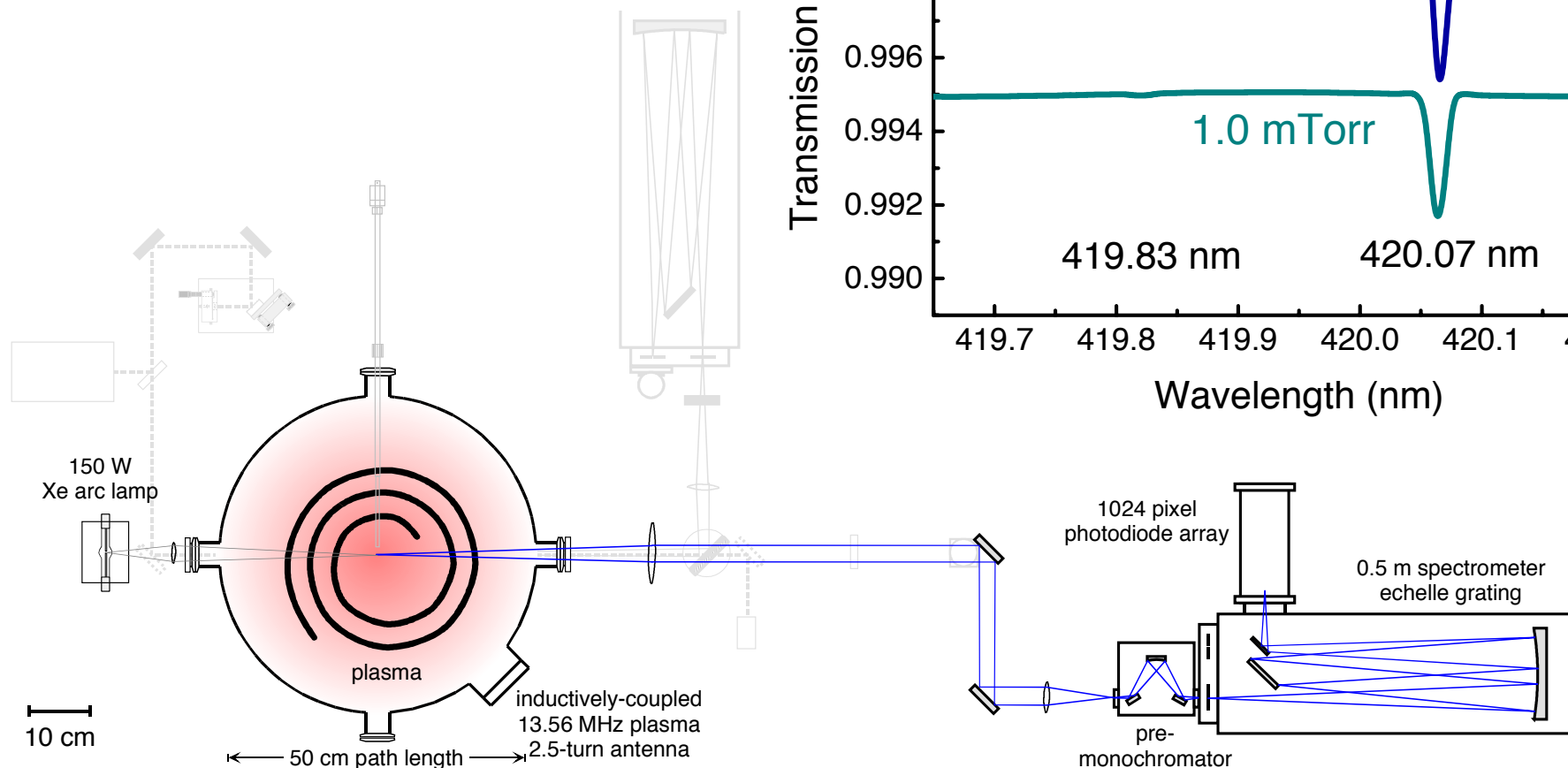
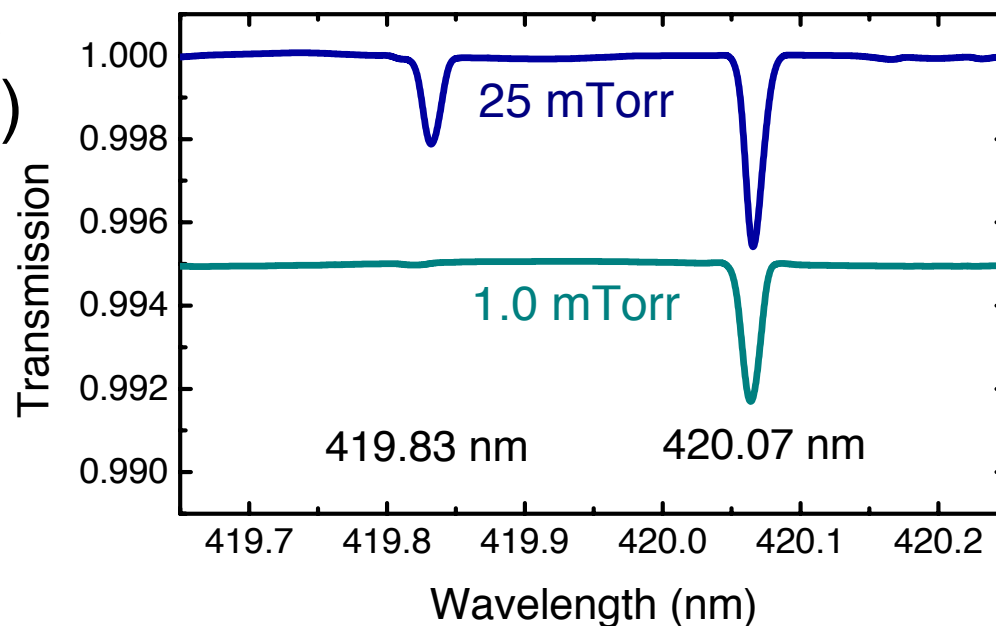
Model input: optical emission spectrum (OES)

- 300-1600 nm wavelength range
- 0.13 nm resolution
- Detectors:
 - PMT 300-870 nm
 - Ge detector: 300-1600 nm



Model input: $1s_y$ concentrations

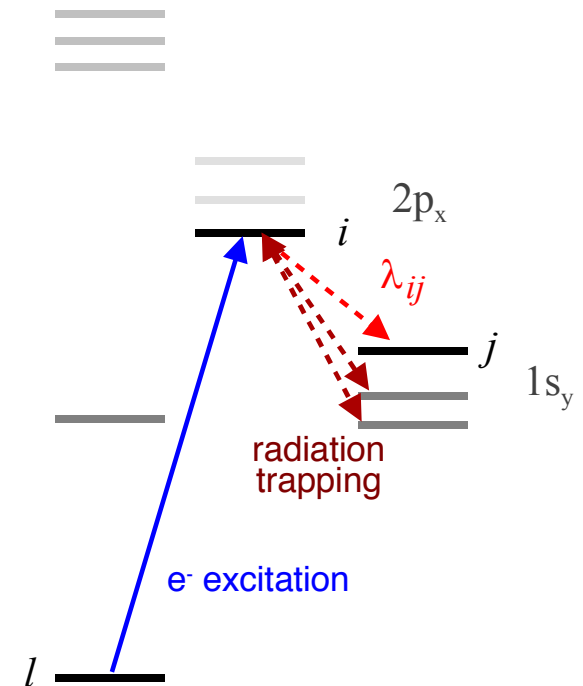
- Method #1: optical absorption spectroscopy (OAS)
 - white light absorption
 - column density measurement
- Method #2: OES (next slide)



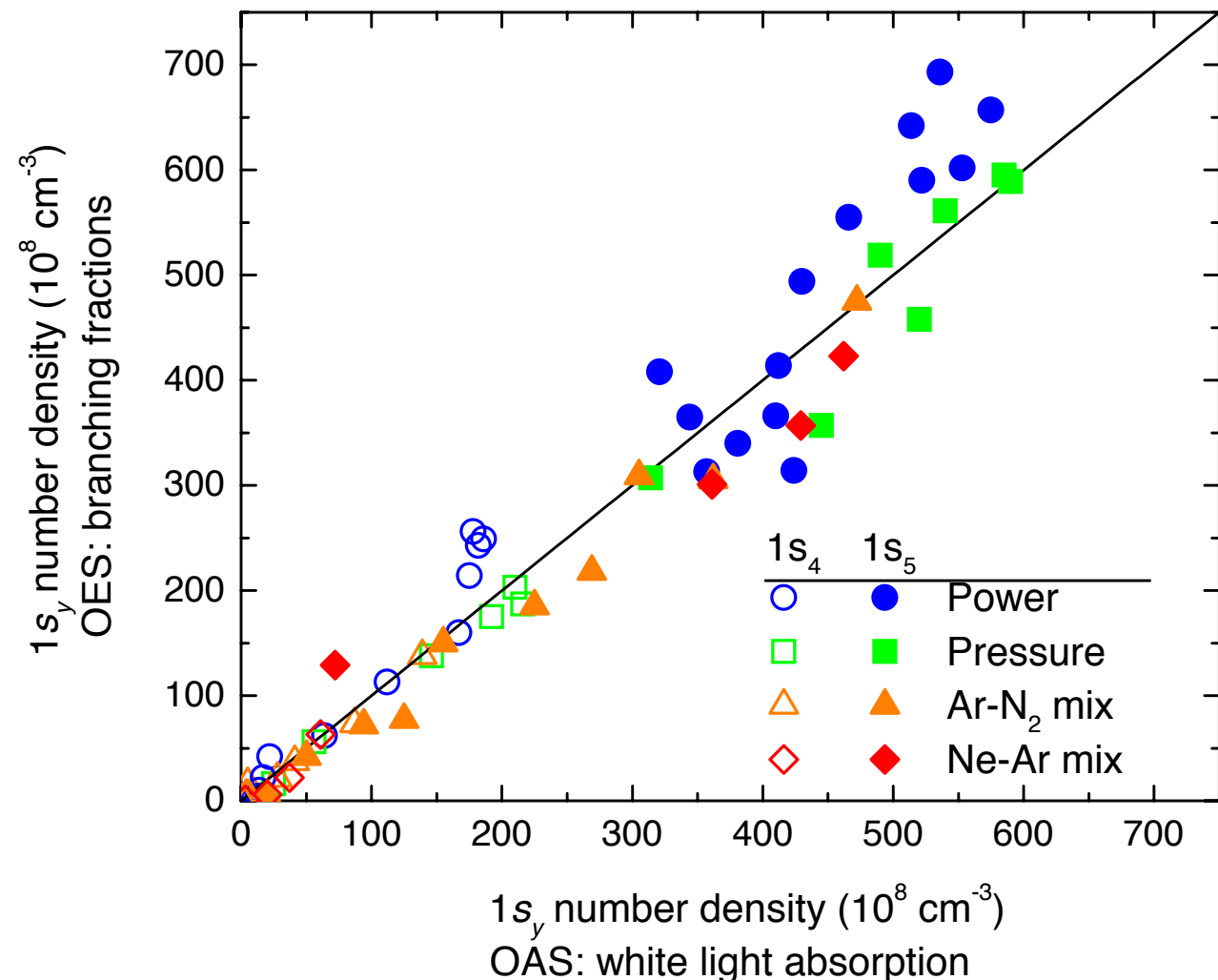
- Method #2: OES
- OES preferable due to simplified instrumentation
- OES branching fraction technique:
 - takes advantage of reabsorption by $1s_y$ metastable and resonance levels
 - reabsorption changes branching fractions
 - use ratio of lines from same *upper* state

$$\frac{\Phi_{ij}}{\Phi_{ij'}} = \frac{A_{ij}^0}{A_{ij'}^0} \frac{g[\kappa_{ij}(n_j)\rho]}{g[\kappa_{ij'}(n_{j'})\rho]}$$

- extract set of n_j densities from fit of model to observed ratios of different line pairs



- OES method agrees with OAS over wide parameter range
 - simple, non-invasive method to determine $1s_y$ densities

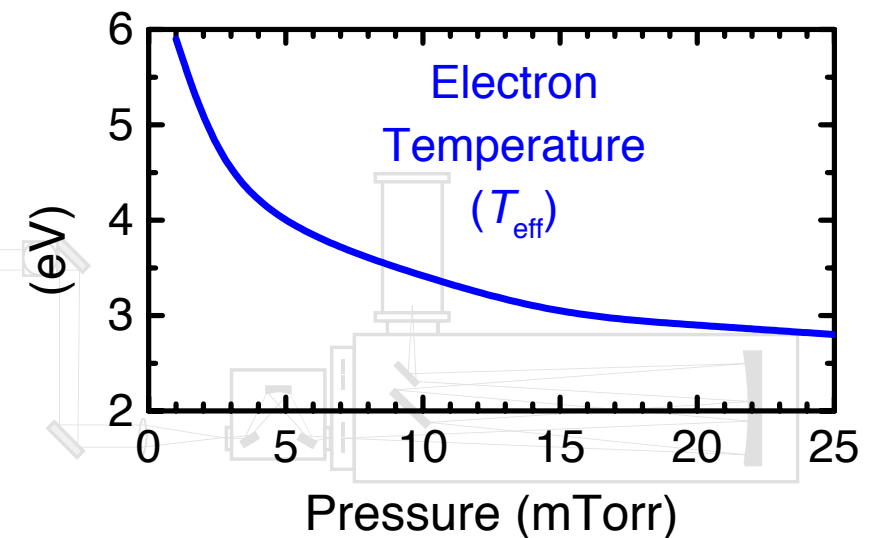
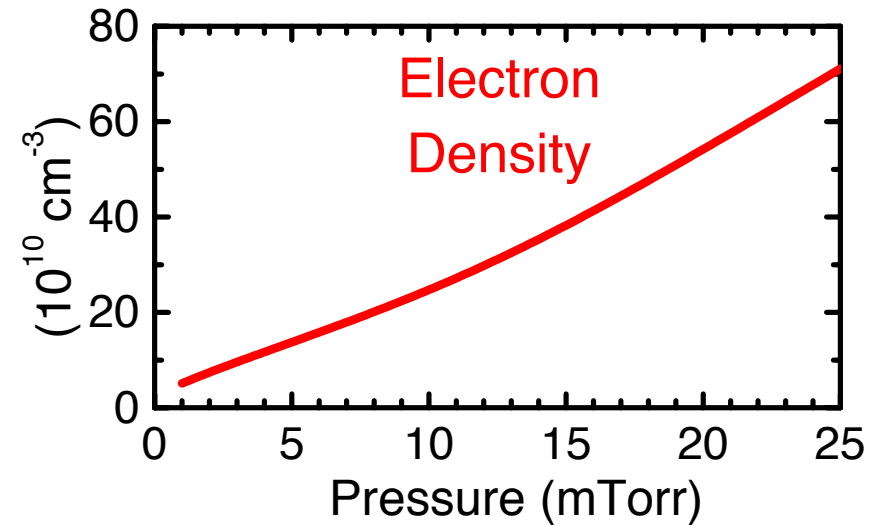
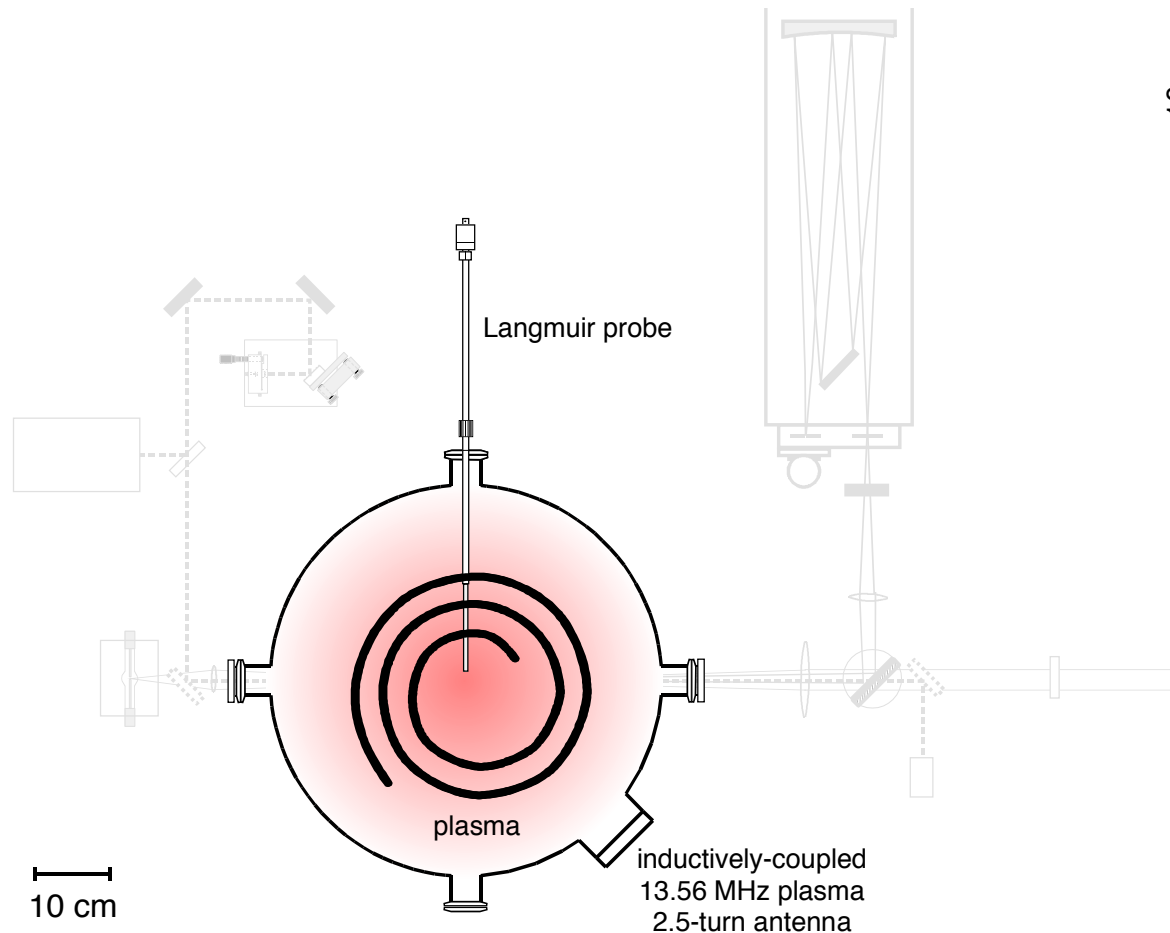


J. B. Boffard *et al.*, *PSST* **18**:035017 (2009).

Compare emission model with EEDF from Langmuir probe

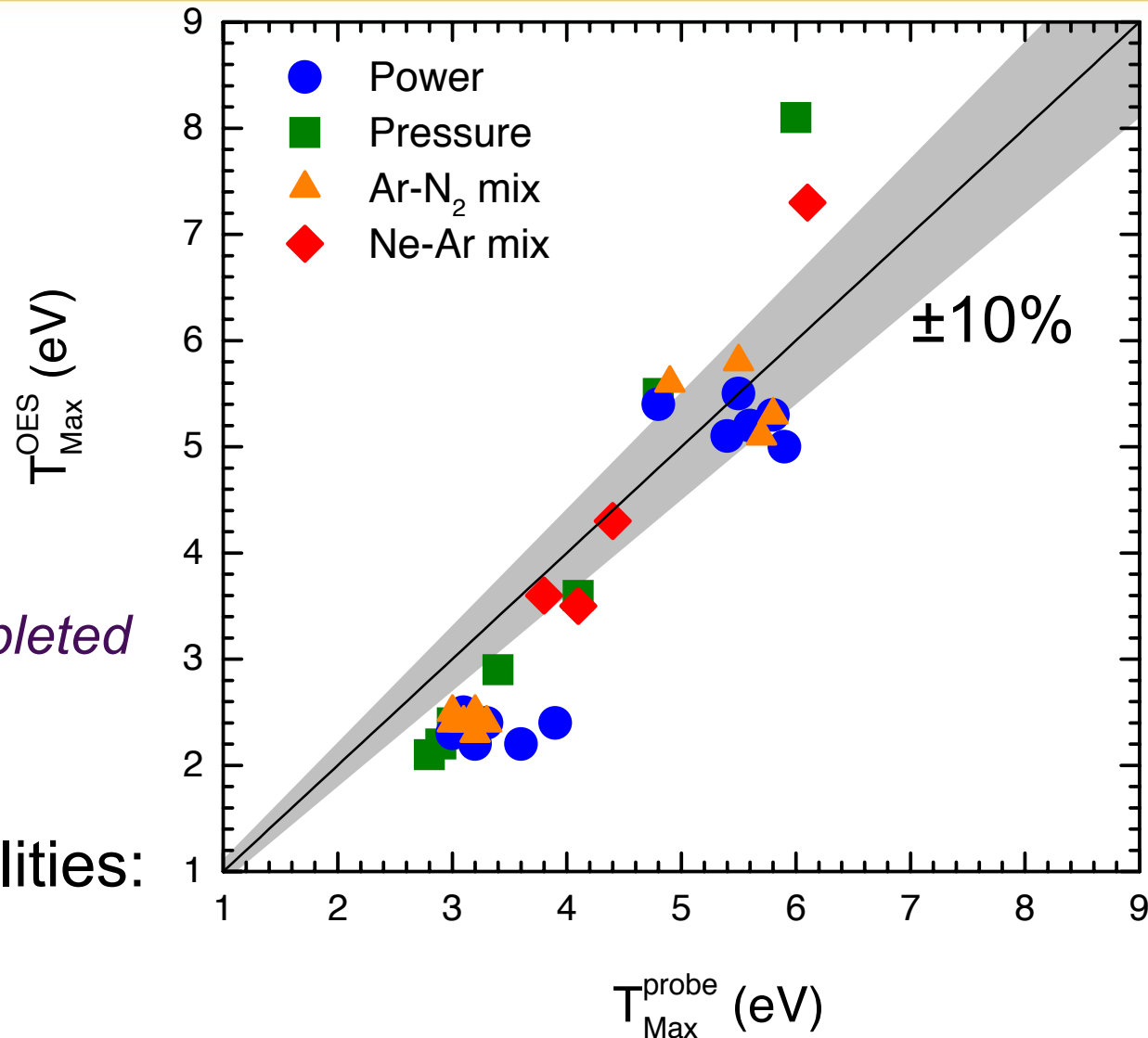
$$f_e(V) = \frac{2m_e}{e^2 A} \sqrt{\frac{2eV}{m_e}} \frac{d^2 I}{dV^2}$$

- compare to emission model output

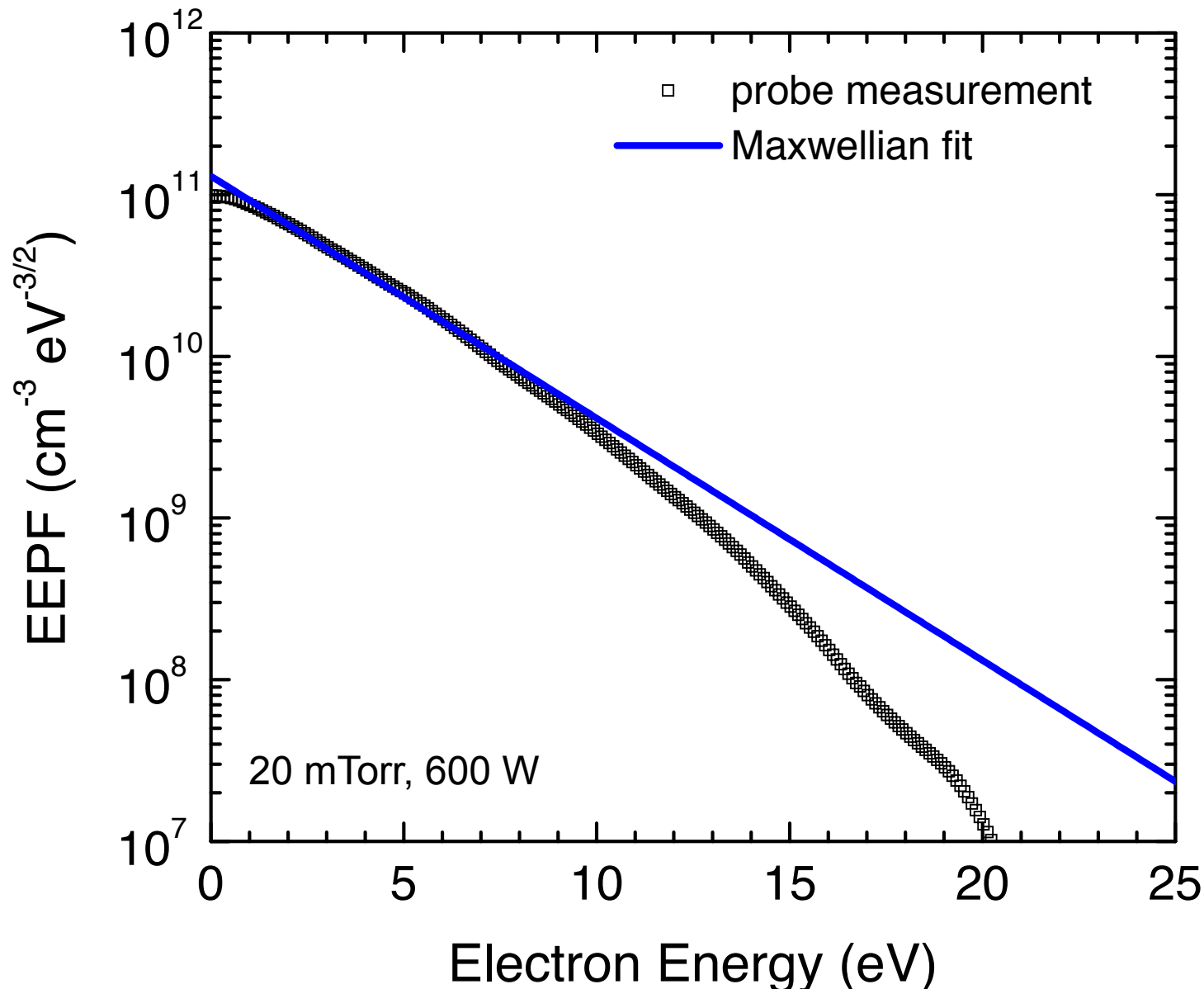


Comparison assuming *Maxwellian* EEDF

- T_e computed from
 - emission model
 - Langmuir probe fit
 - *agreement not great*
- reason for poor fit
 - EEDF is *not* Maxwellian
 - *actually, EEDF “tail” is depleted compared to Maxwellian*
- other (less likely) possibilities:
 - incorrect n_I values
 - incorrect cross sections



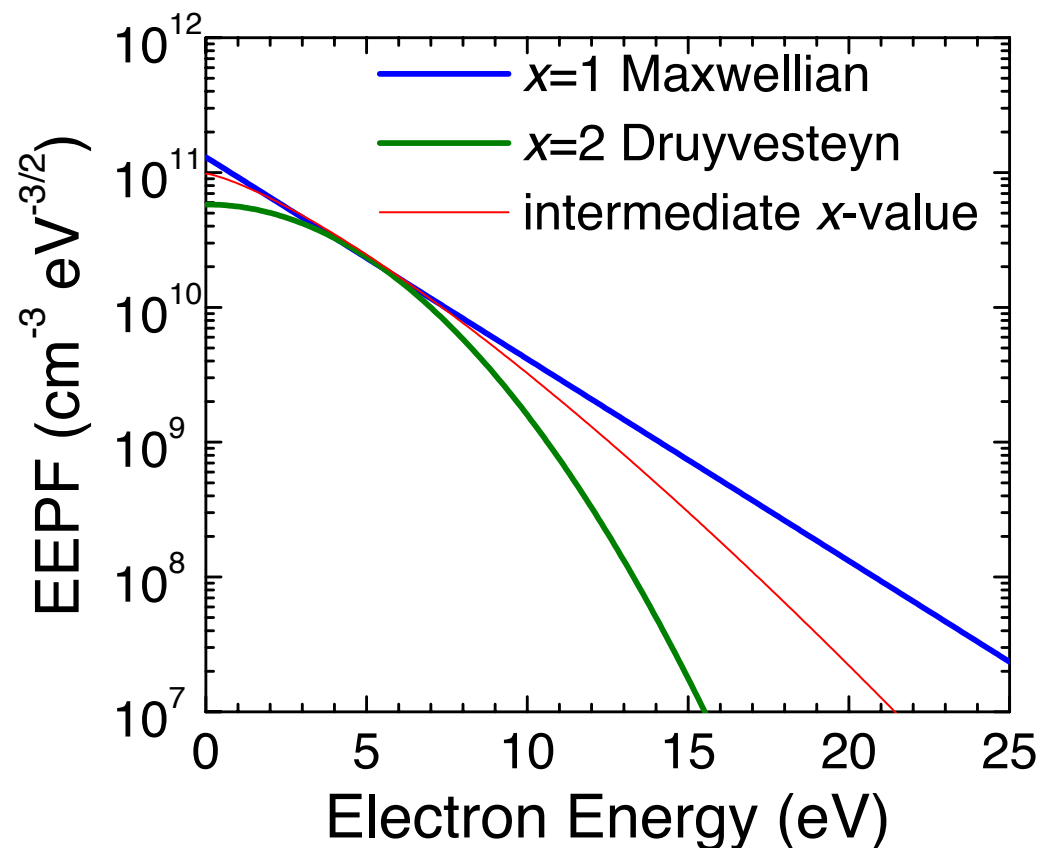
- Langmuir probe measurement in ICP shows “tail” depletion



- Goal: account for non-Maxwellian energy dependence
- use “generalized x-form”

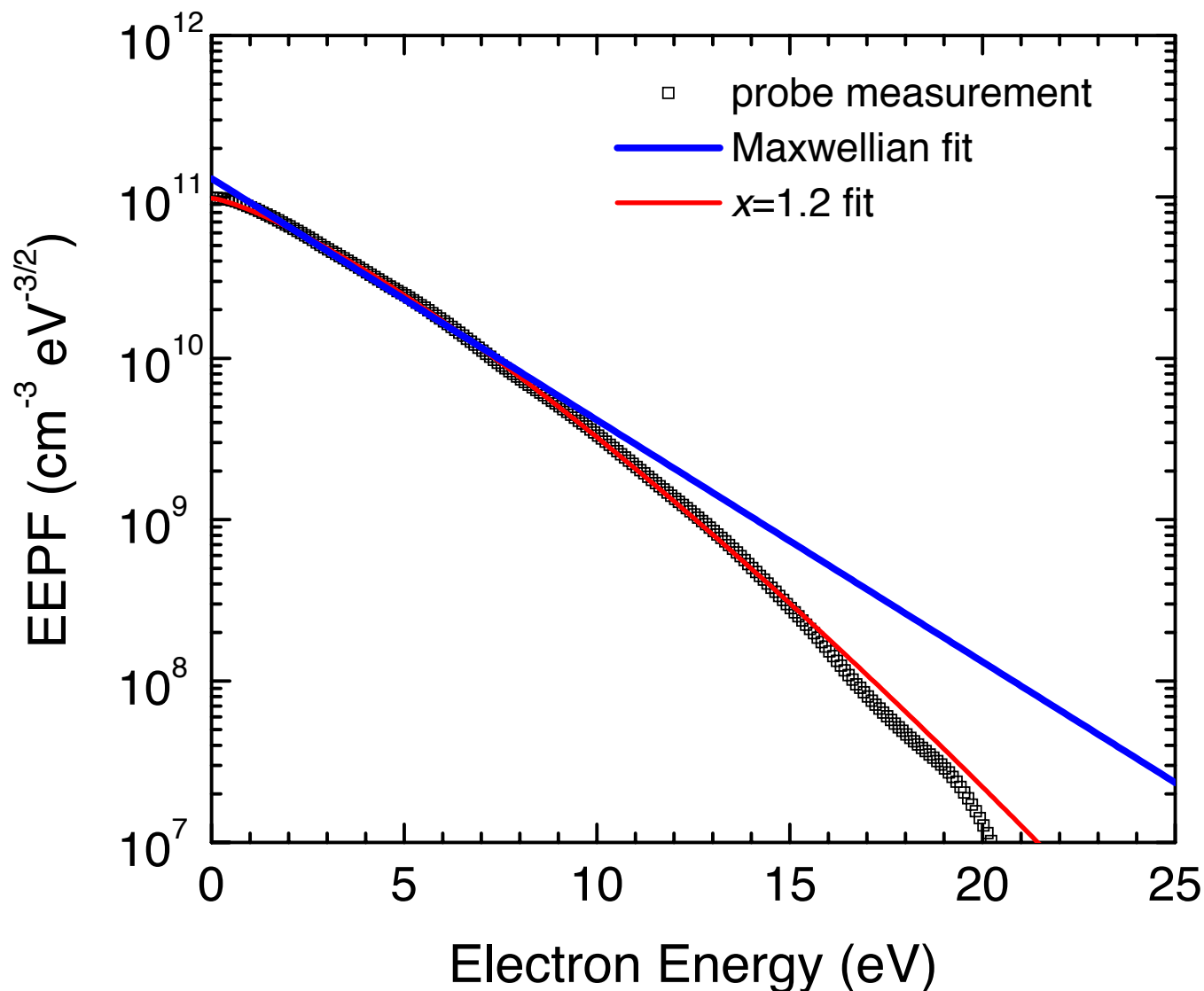
$$f_x(E) = c_1 T_x^{-3/2} \sqrt{E} e^{-c_2 (E/T_x)^x}$$

- T_x is an *effective* electron temperature ($T_x = 2/3 \langle E \rangle$)
- c_1 and c_2 are functions of x
- $x=1$ corresponds to a Maxwellian
- $x=2$ corresponds to a Druyvesteyn
- x -form good fit for most ICP conditions



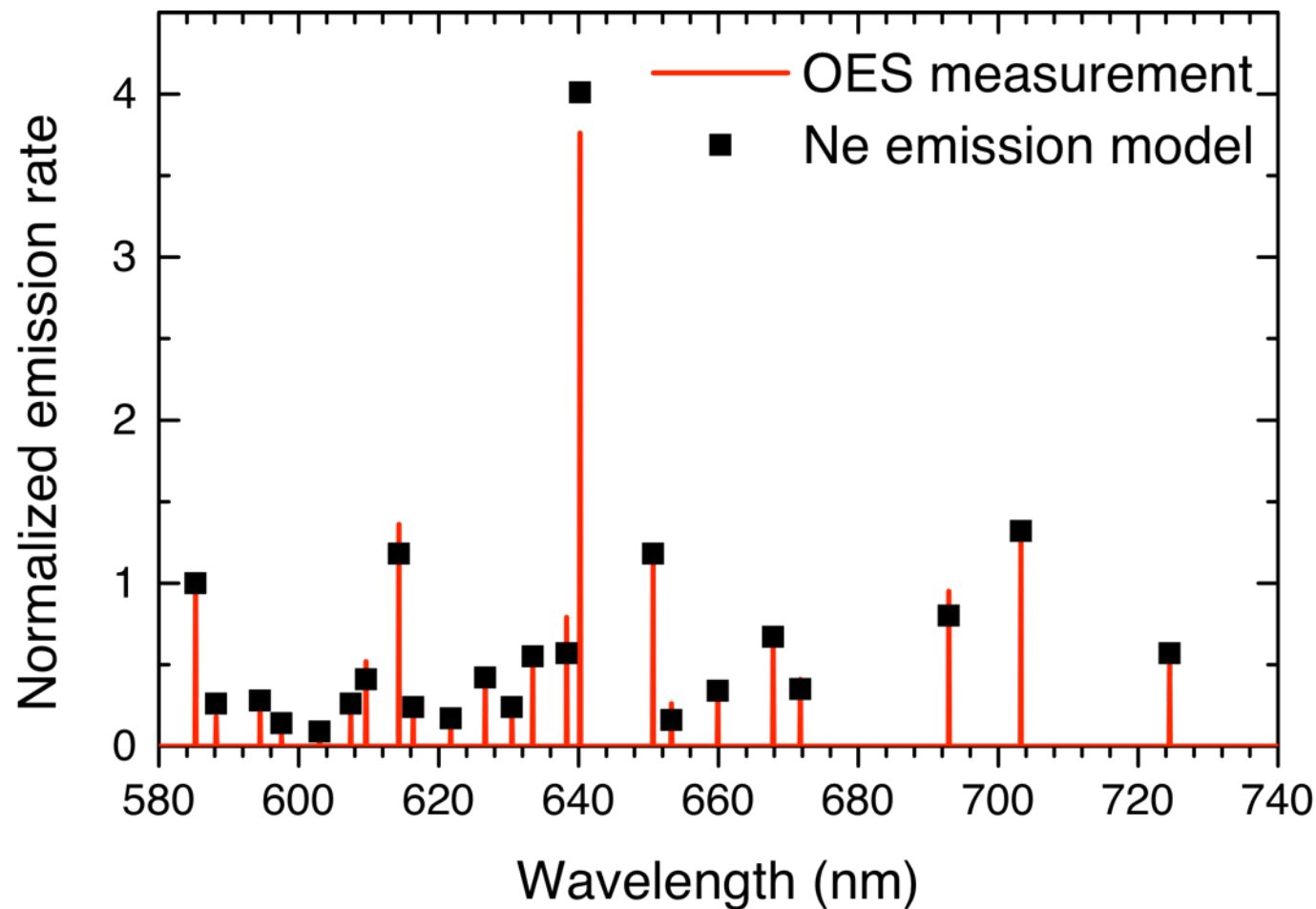
EEDF with $x=1.2$

- example shows improved fit to Ar probe data
- x ranges from $x=1.2$ to $x=1.6$ for Ar ICP



OES also produces $x=1.2$

- emission model agrees well with measurement*

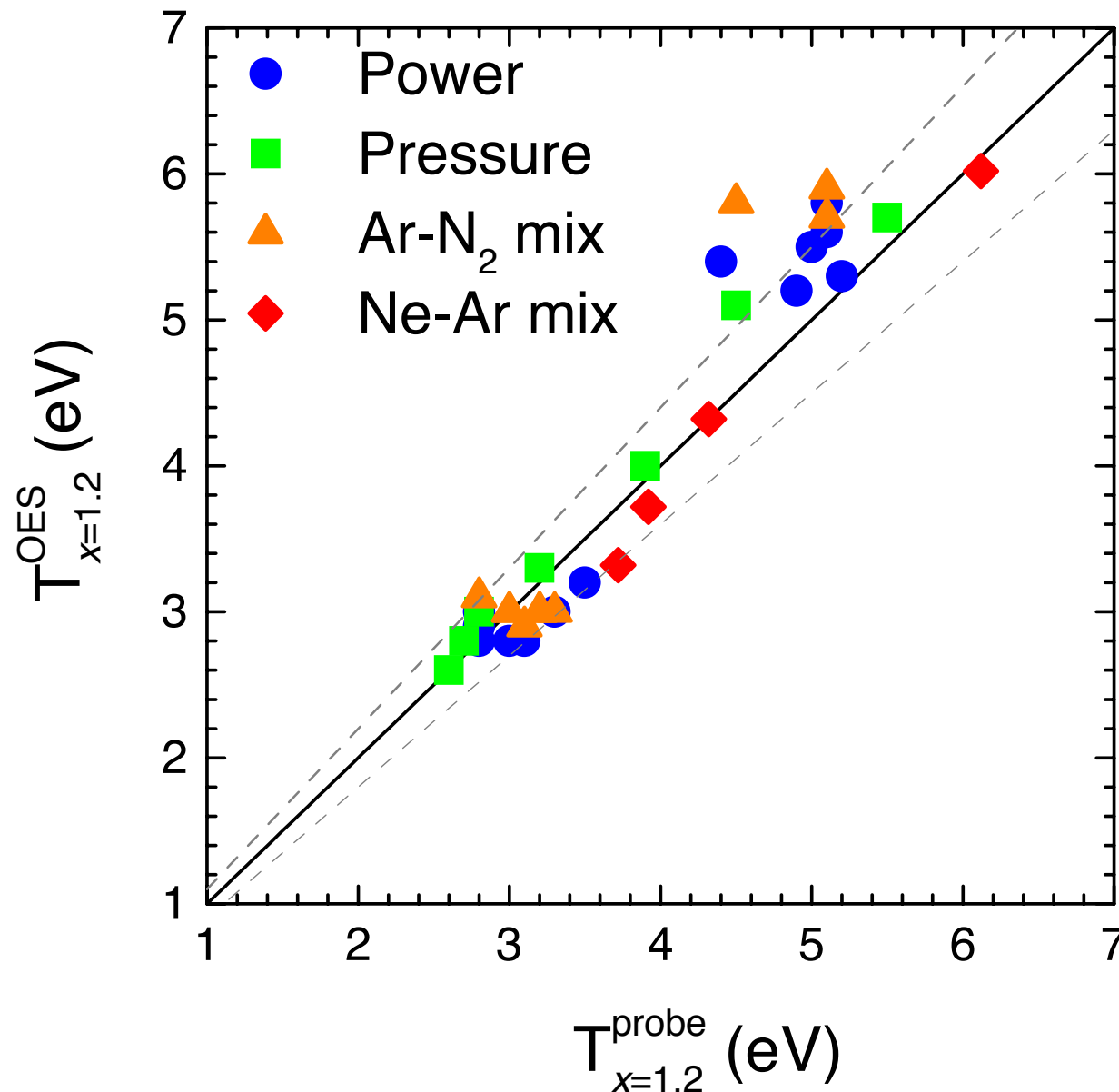


*Plasma Sources Sci. Technol. 19 (2010) 065001.

- $3p_x$ levels have longer radiative lifetimes than $2p_x$ levels
 - $2p_x$ levels: 20-35 ns
 - $3p_x$ levels: 80-160 ns
- Longer radiative lifetime increases chances of electron-atom collision while atom is in excited level
 - electron-quenching: non-radiative decay out of level
 - electron-transfer: excitation from $3p_x$ to another $3p_x'$ level
 - emissions redistributed
 - need modified emission model
- *Bonus*: modified emission model predicts electron density

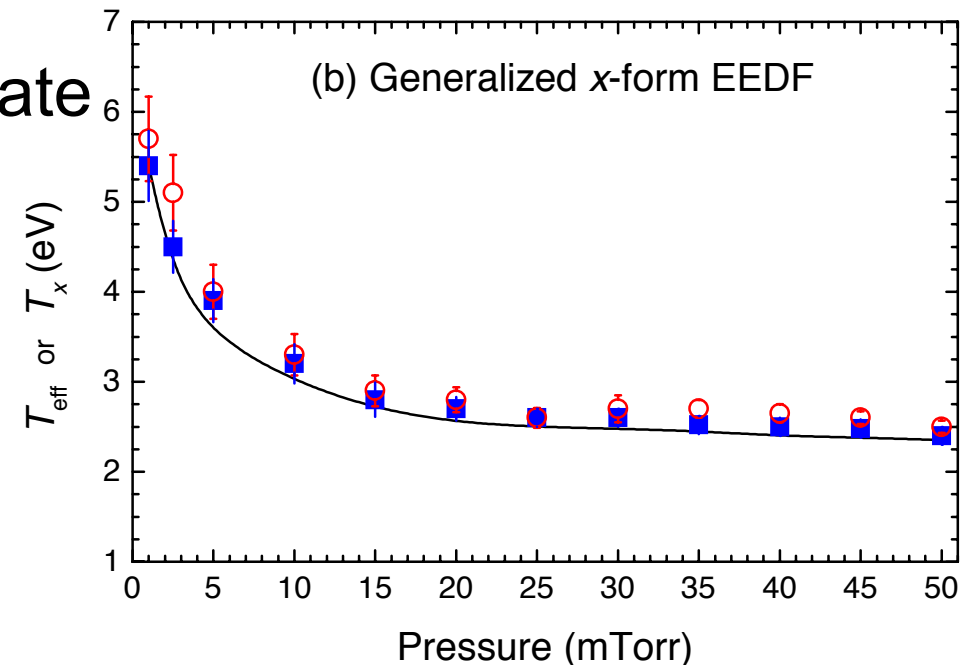
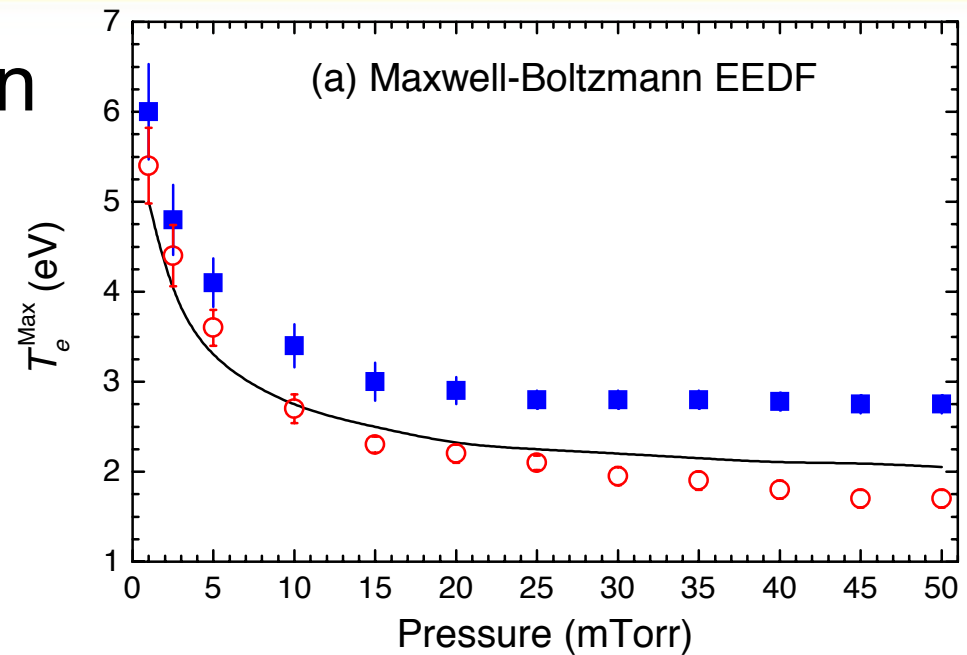
Emission model including 2p and 3p

- agreement with probe: $\pm 7\%$



Ar pressure dependence

- Comparison of a) Maxwellian and b) x-form for
 - Probes (solid)
 - OES (open)
 - 'global' model (line)
- Global model
 - Volume averaged
 - Particle balance – ionization rate must match wall losses

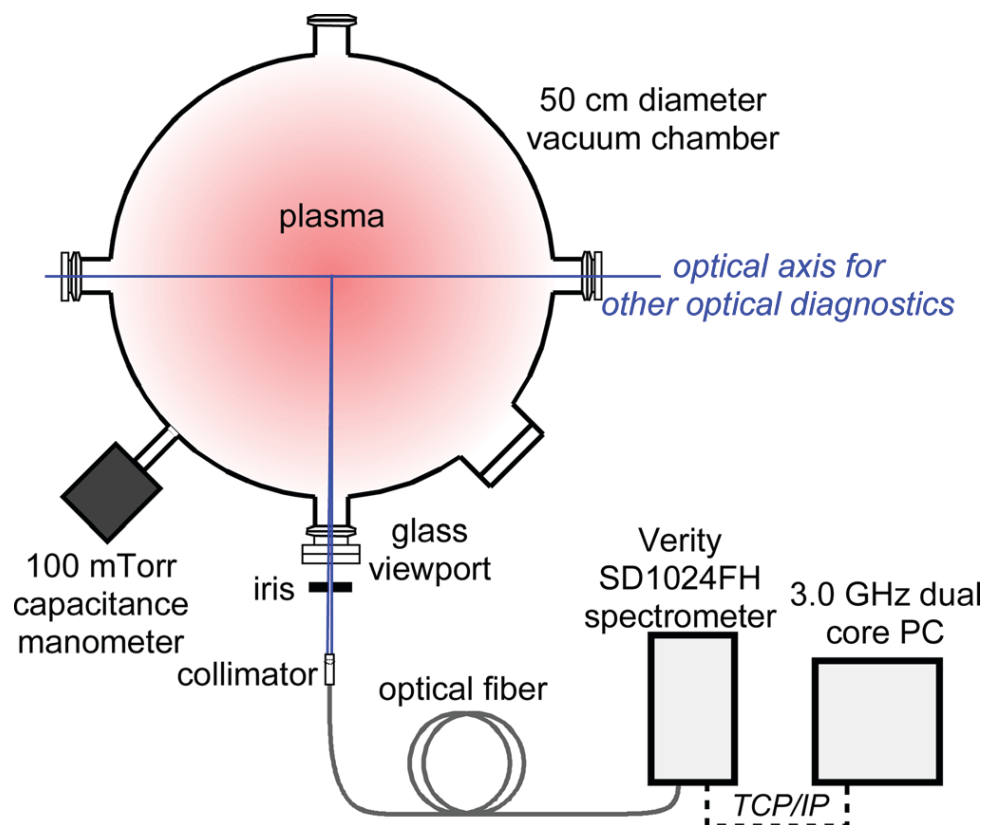


- Non-Maxwellian EEDFs have been observed using Ar emissions for wide range of ICP conditions
 - sensitive to fairly subtle changes in EEDF
 - “x-form” of EEDF well suited for ICP
 - ♦ captures EEDF tail depletion
- OES - BF is a simpler alternative to absorption spectroscopy to determine $1s_y$ concentrations
- emission model must include contributions from:
 - excitation to higher lying states followed by cascade
 - excitation from metastable/resonance levels
 - reabsorption by metastable and resonance levels
 - electron-driven transfer between emitting levels closely spaced in energy
 - ♦ Bonus: allows determination of electron density n_e

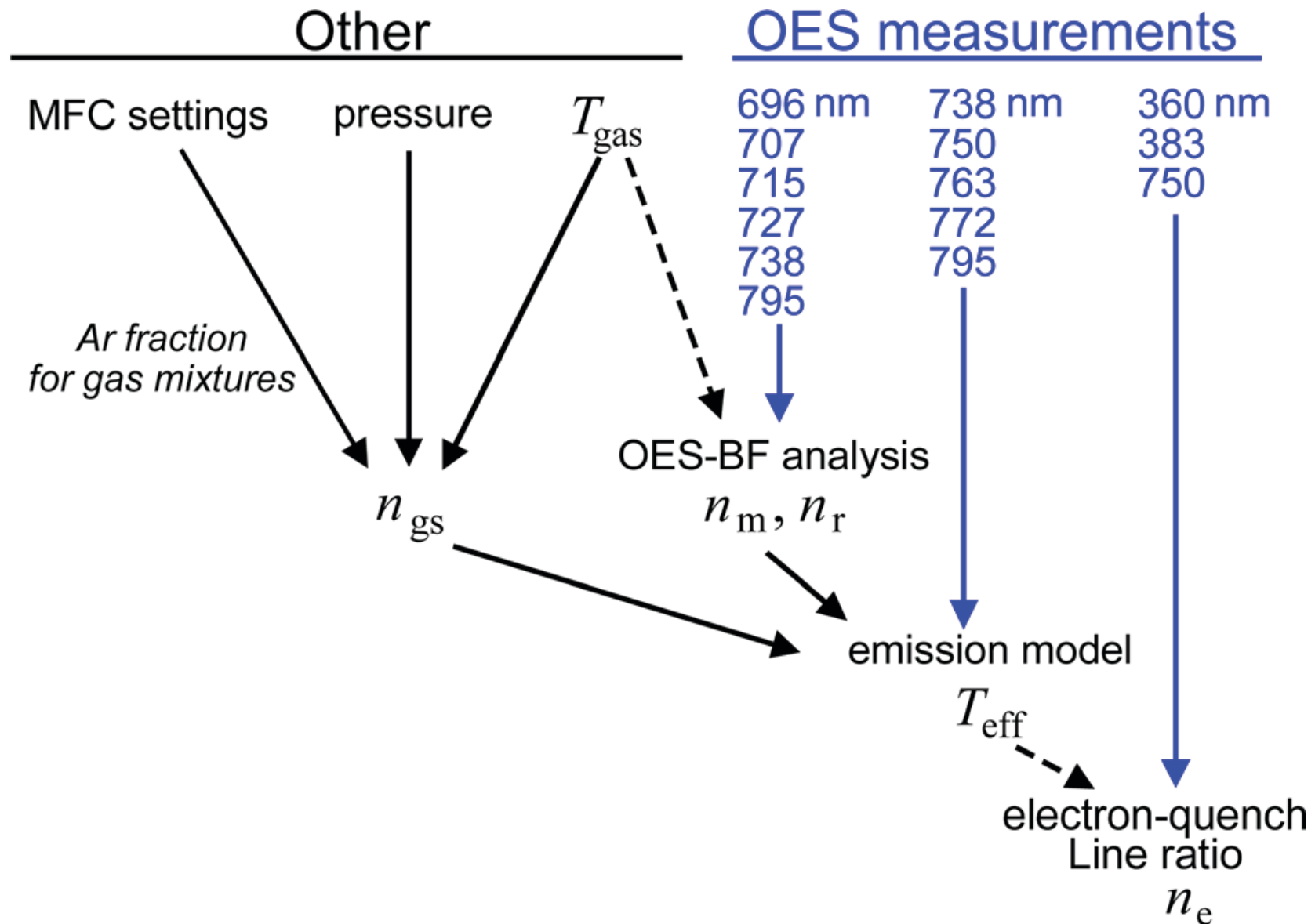
Real-time monitor of plasma conditions: time-varying and pulse-modulated ICP

OES plasma diagnostics in real time

- Verity spectrograph
 - fast spectra acquisition
 - compact
 - low-resolution (~ 1 nm) spectra (200-800 nm)
 - similar to systems often included on plasma process equipment

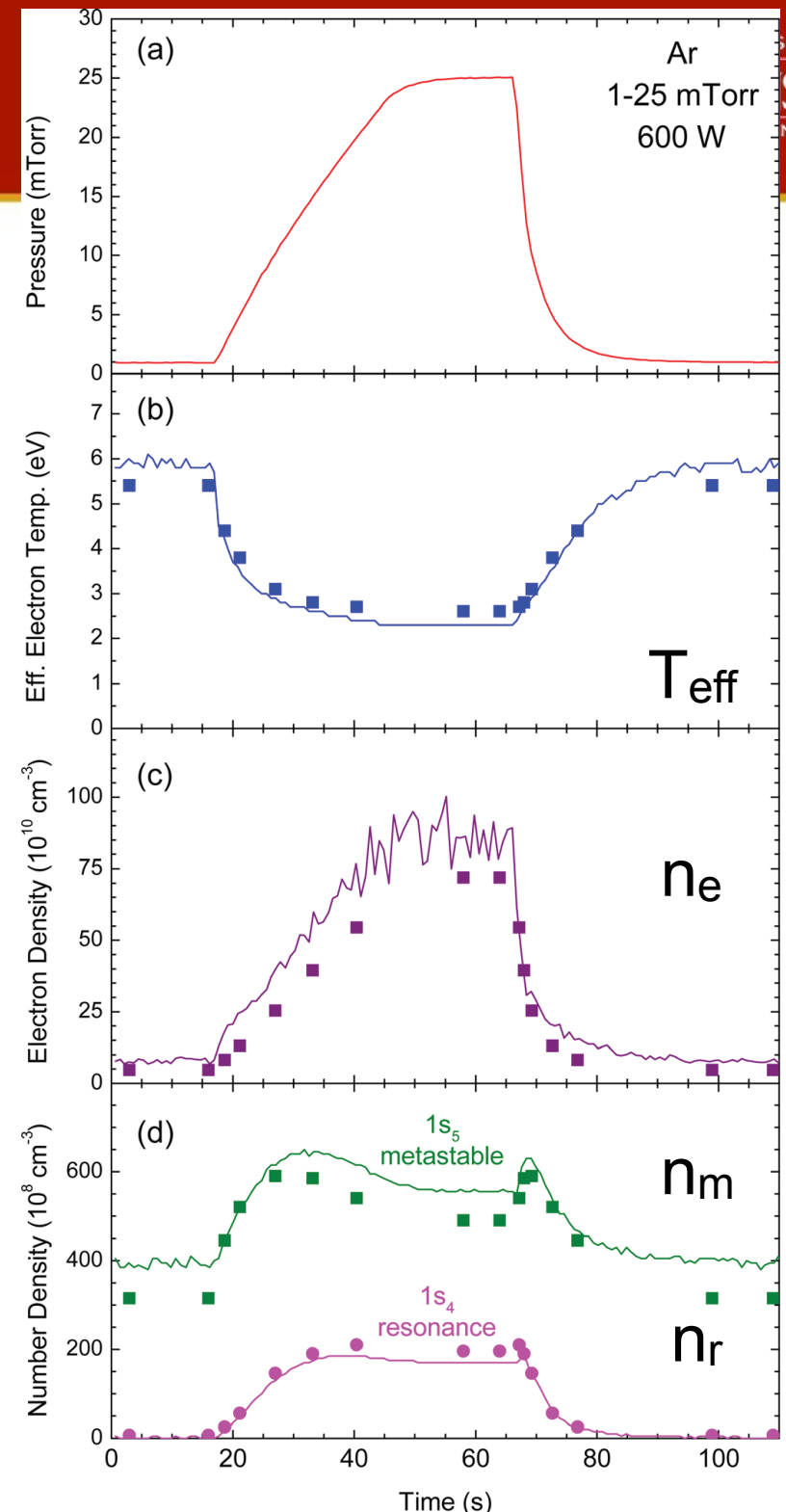


Flow chart for real time measurements



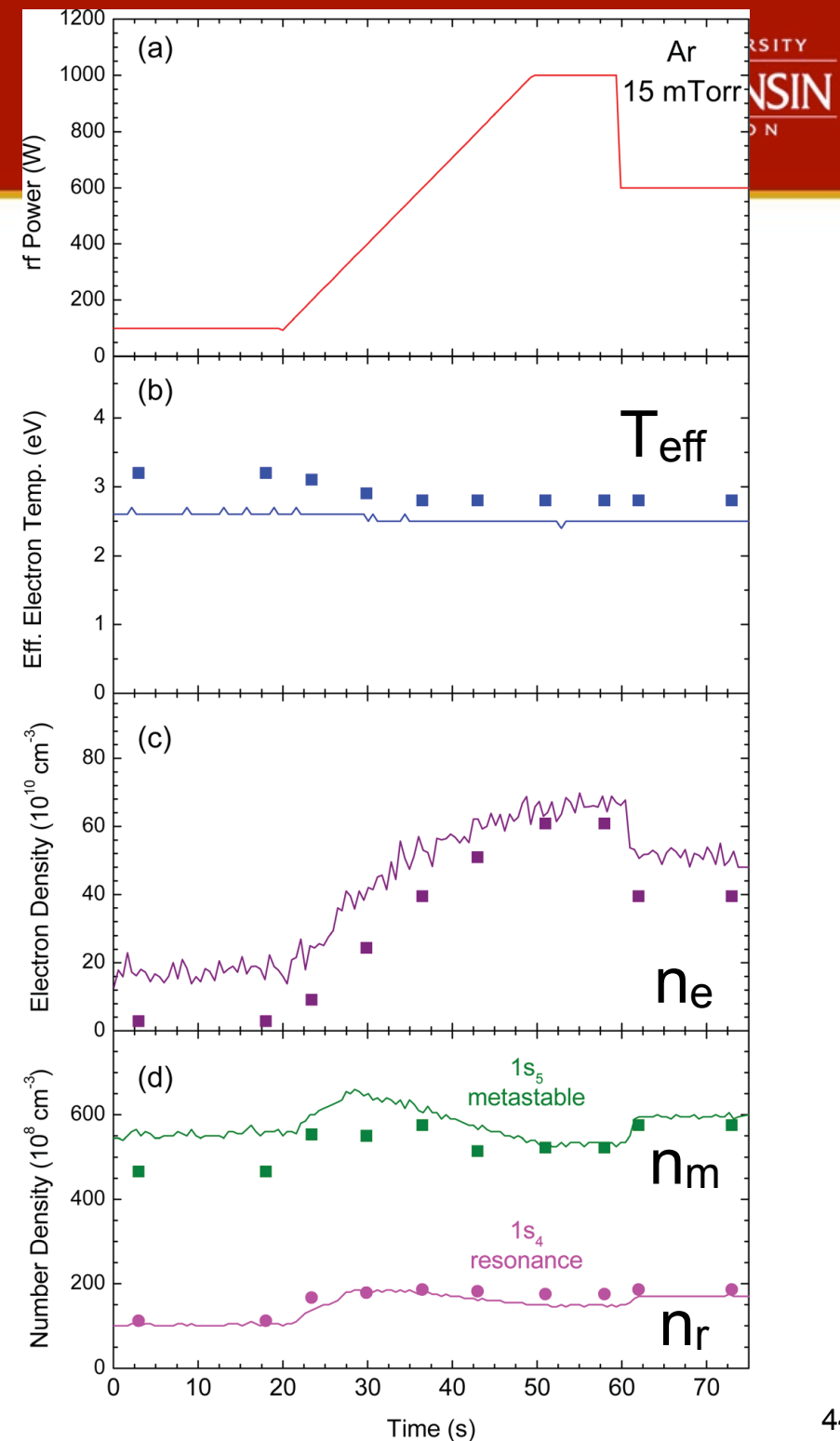
Pure Ar: real time/ pressure varied

- Computer controlled time-varying Ar *pressure*; power held fixed at 600 watts
- plasma parameters evolve in time in response to pressure change
- solid lines show instantaneous measurements vs. time
- dots are static measurements made with the same parameter settings

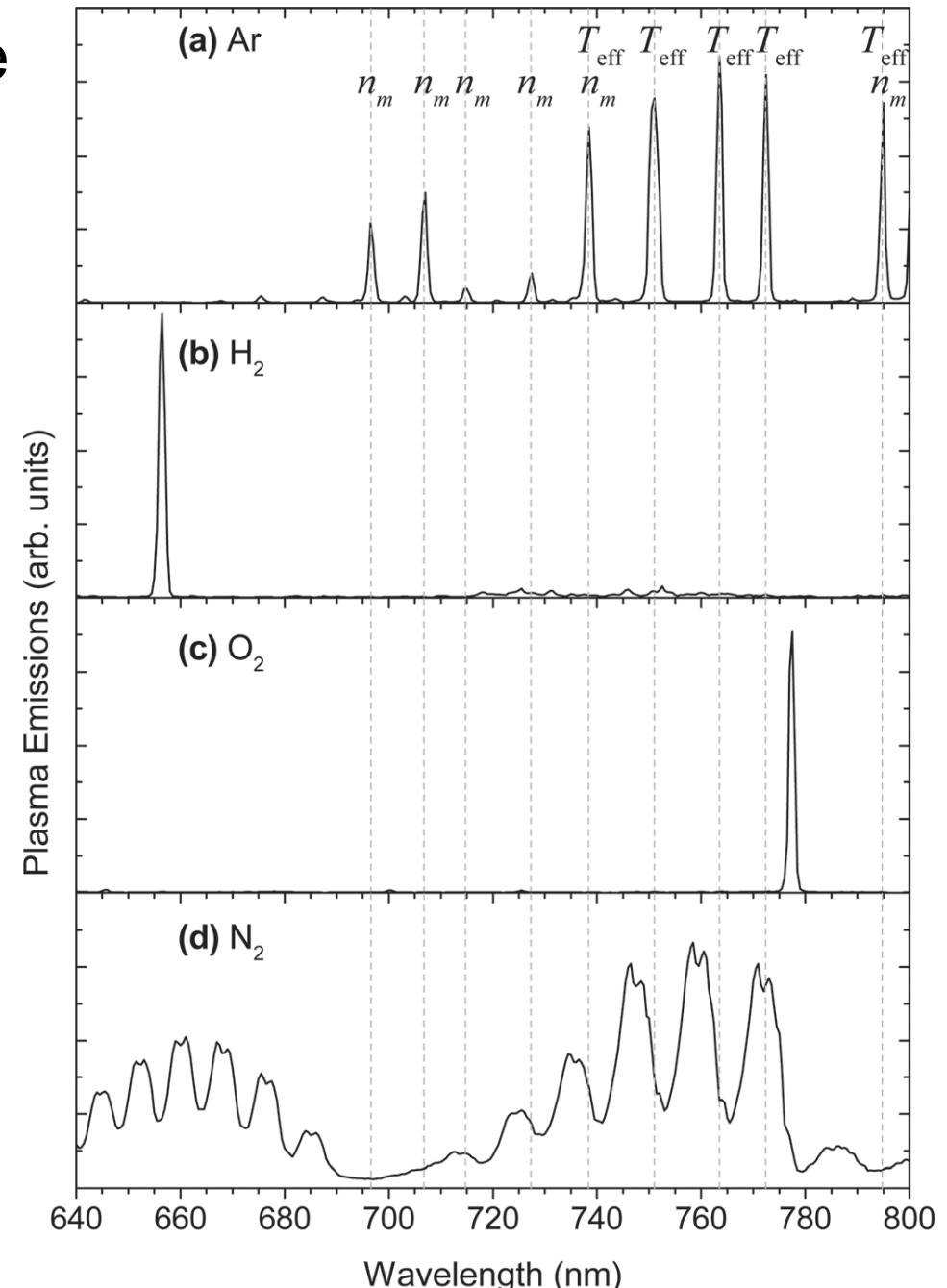
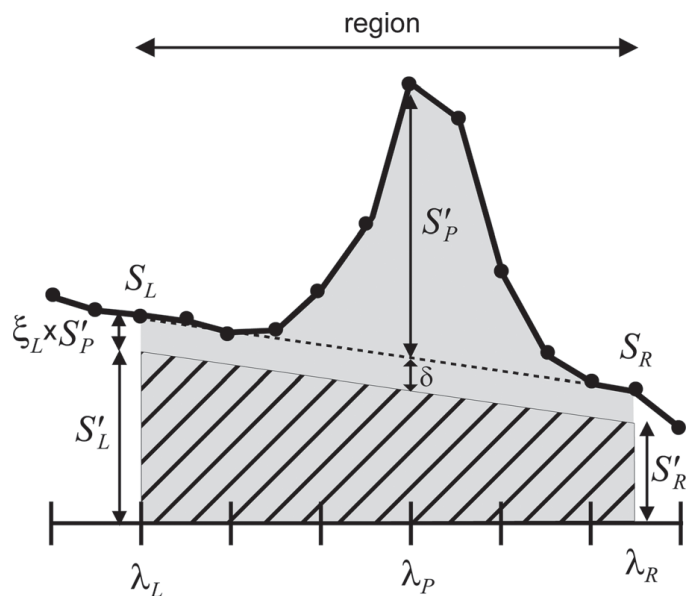


Pure Ar: real-time/ power varied

- Computer controlled time-varying Ar *power*; pressure held fixed at 15 mTorr
- plasma parameters evolve in time in response to pressure change

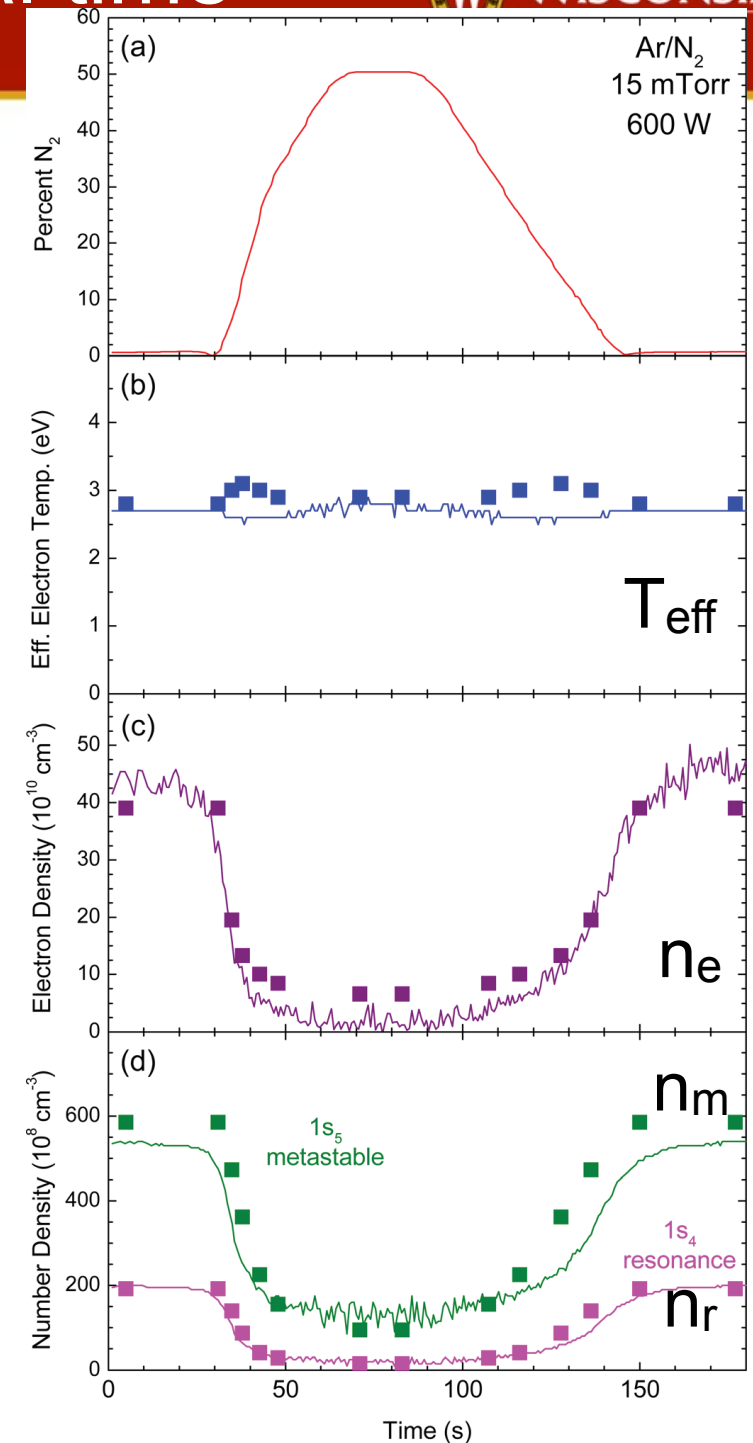


- Challenge: difficult to separate OES lines
 - from neighboring Ar lines
 - from molecular OES features



OES plasma diagnostics in real time

- Computer controlled Ar and N₂ flow rates; pressure held fixed at 15 mTorr, power at 600 W

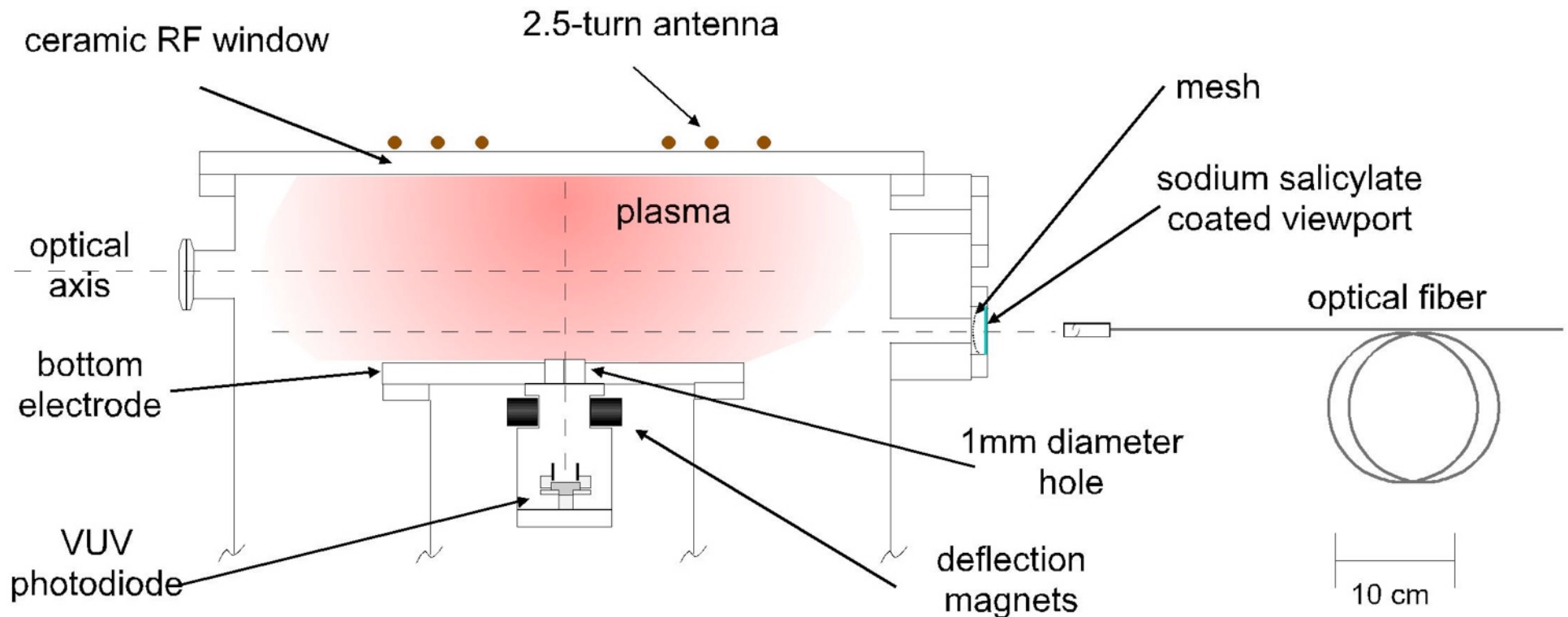


Detection of VUV in argon-containing ICP

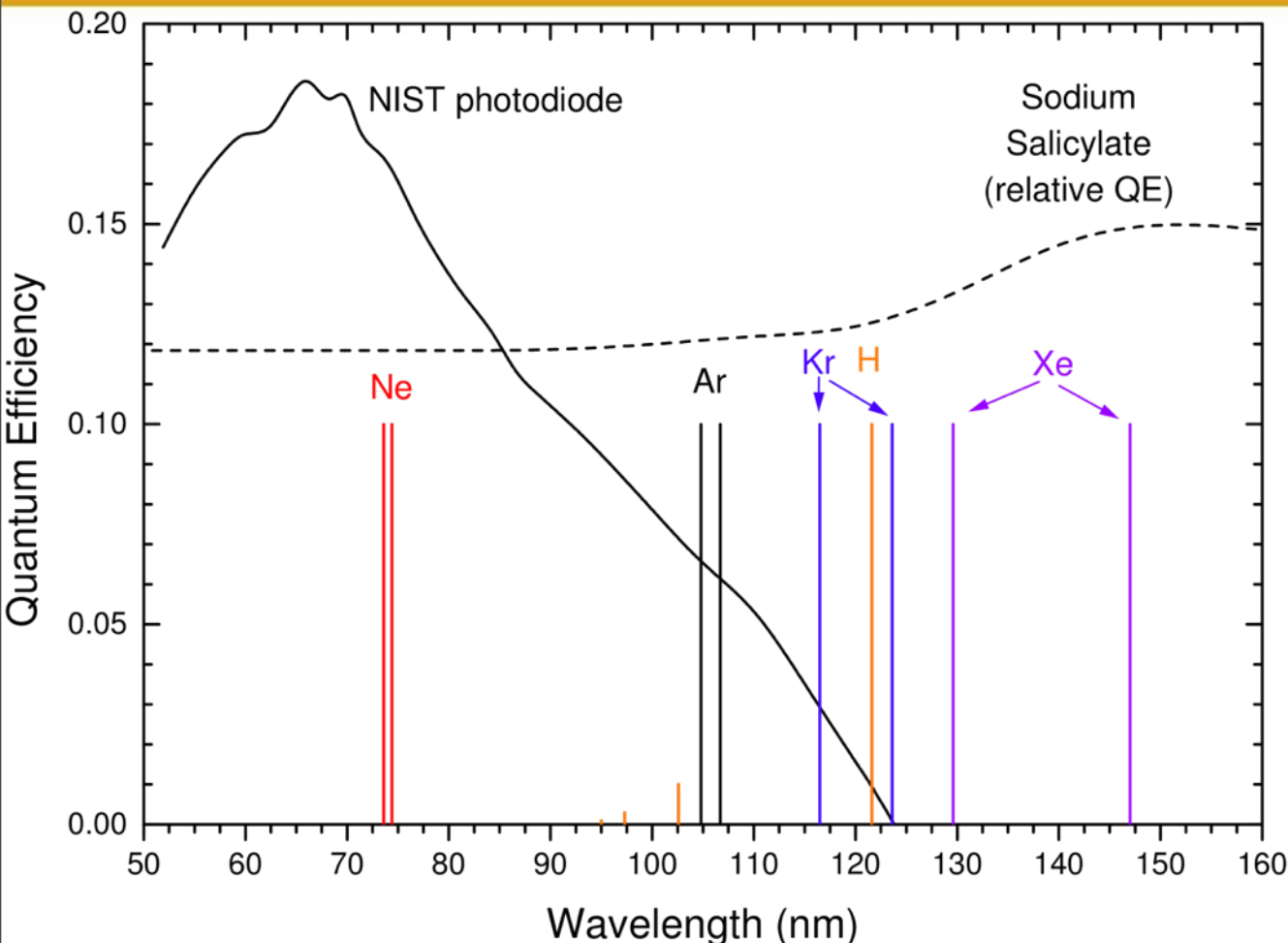
- VUV flux to wafer matters in plasma processing for integrated circuit manufacturing (and other applications!)
- Diagnostic challenge: no way to non-invasively measure VUV flux directly
 - VUV is *not* transmitted by windows or air
- In Ar containing plasmas, VUV is emitted by resonance level atoms. We posed the following question: can non-invasive measurements of n_r serve as a proxy for VUV?
- Experimental evidence indicates the answer is a *qualified* “yes”...

VUV Photon Flux Detection – 2 methods

- Windowless NIST photodiode installed below ground electrode
- Relative VUV flux detected through visible fluorescence from sodium salicylate coated viewport

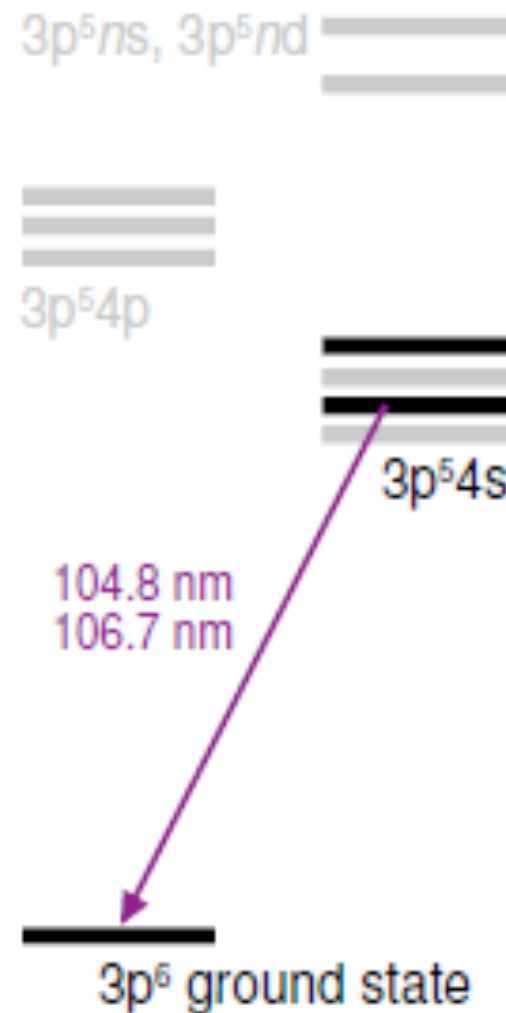


Sodium salicylate signal: sensitive at wavelengths beyond photodiode limit



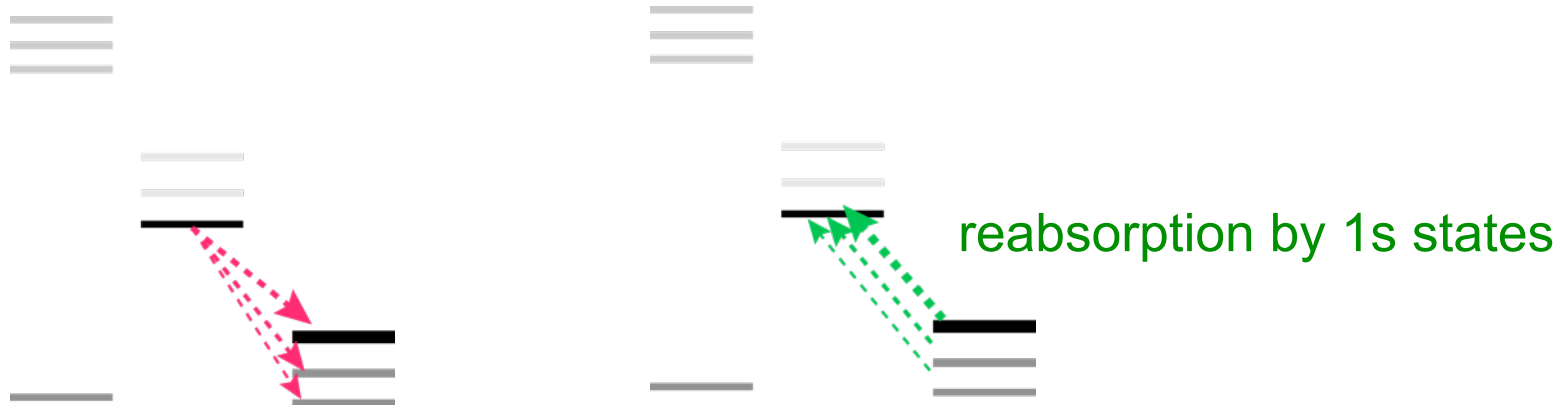
Due to *different* wavelength dependent sensitivities, relative signal strengths for photodiode and sodium salicylate vary with gas type

Ar energy levels

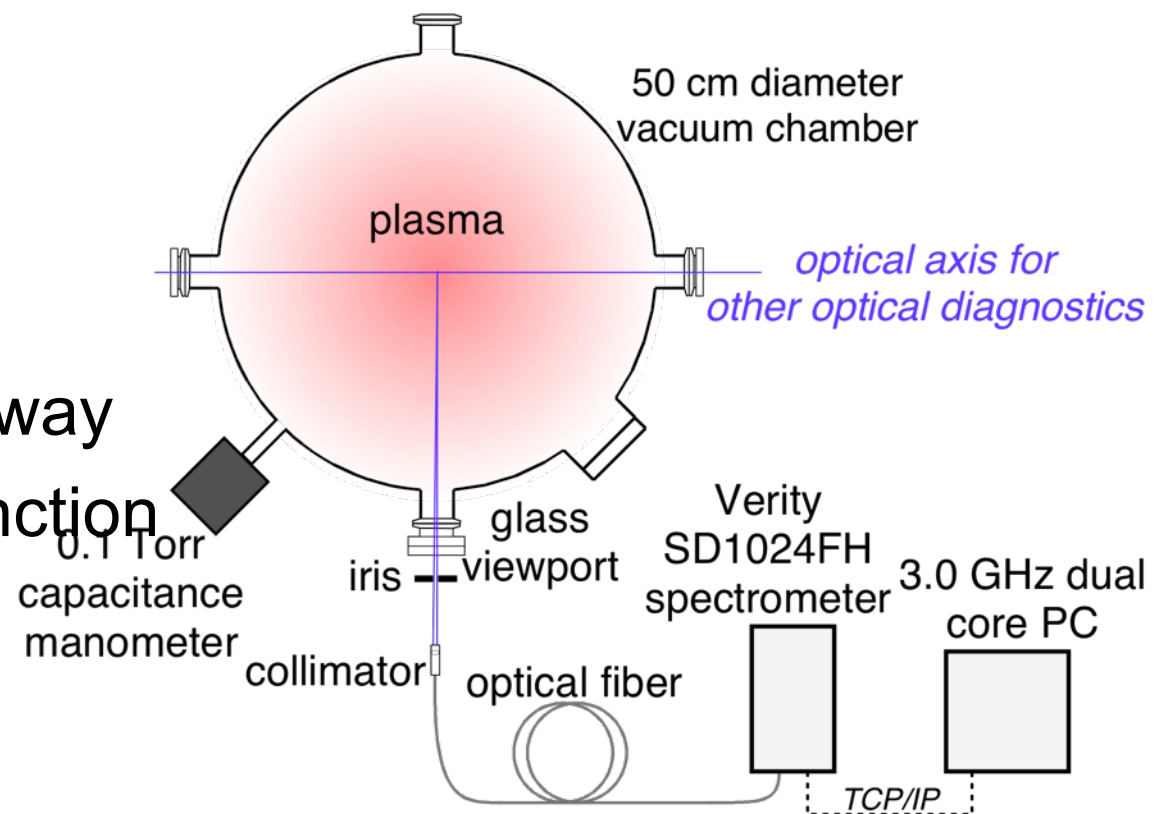


Non-invasive measurement of resonance level density

Branching Fraction



- Intensity ratios *altered* by reabsorption in quantifiable way
- Reabsorption is a known function of resonant level density

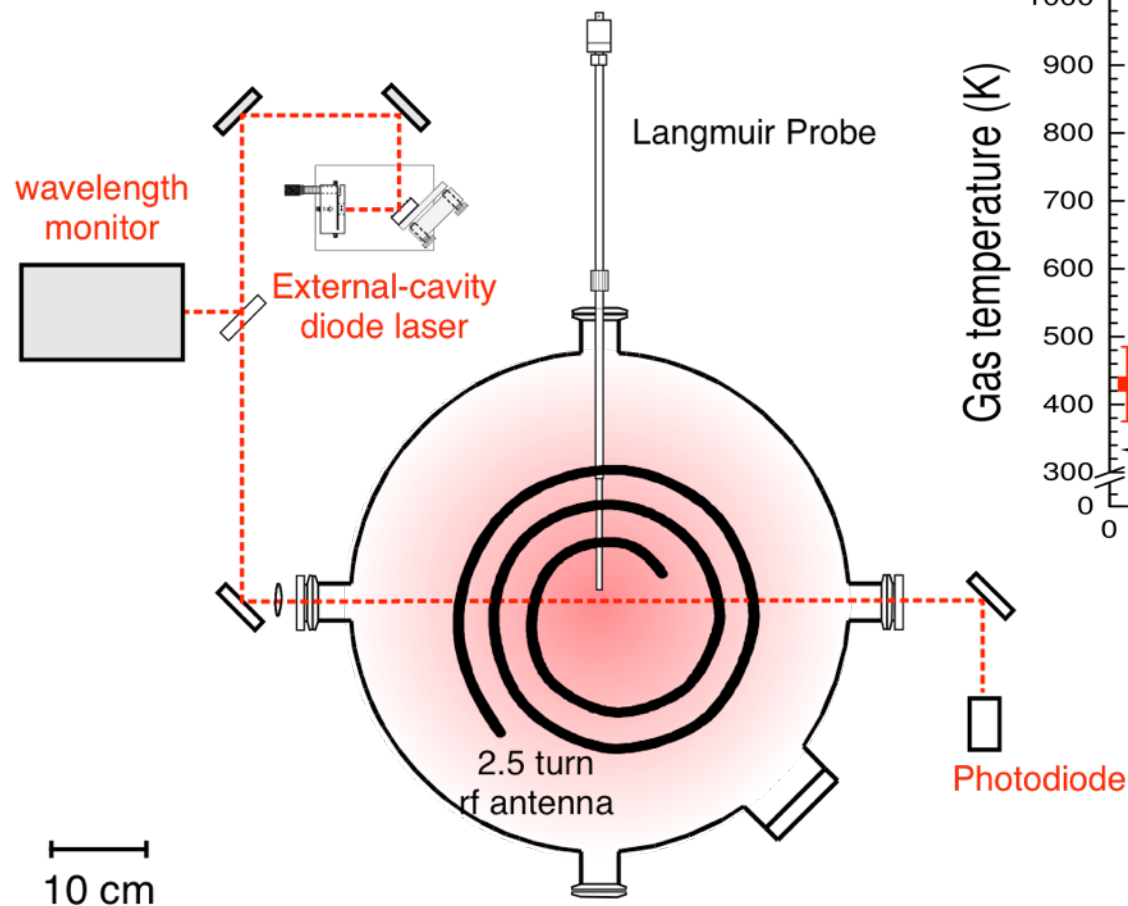


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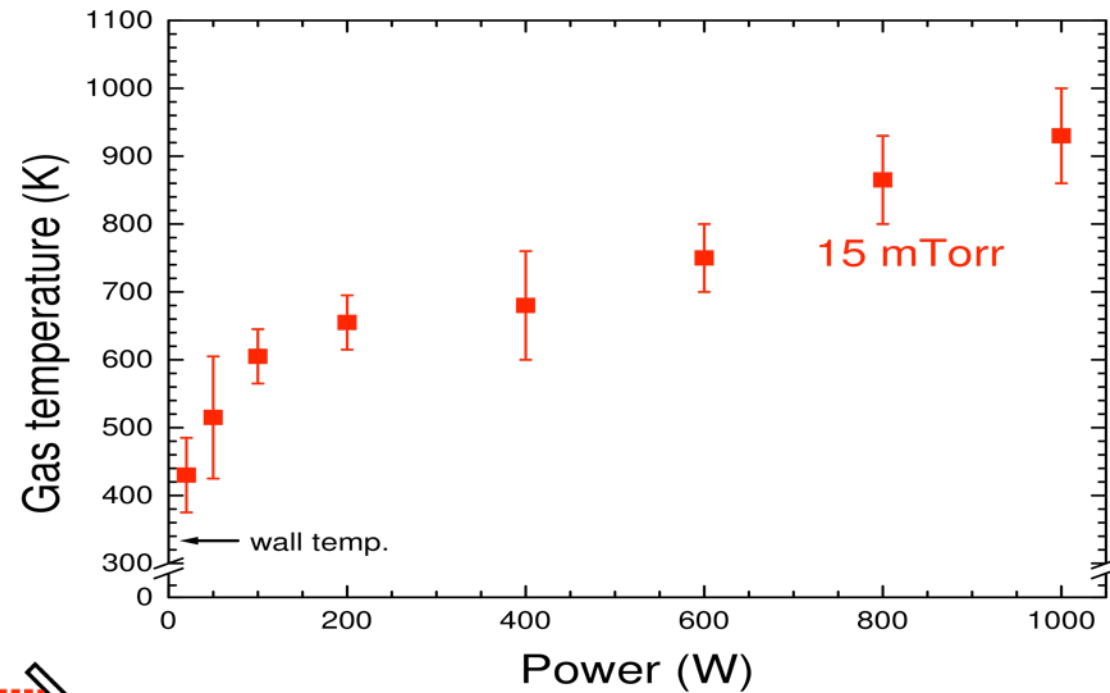
- Model elements
 - Local VUV emission is proportional to resonance density
 - Chance of emitted photon reaching VUV photodiode depends on:
 - ♦ Geometry – line of sight through pinhole
 - ♦ Reabsorption – “escape factor” depends on $n_{\text{gas}} = p/kT_{\text{gas}}$, thus:
 - Gas pressure, p
 - Gas temperature, T_{gas}
- Both photodiode signal and VUV flux to walls calculated

Non-invasive measurement of gas temperature

Diode laser absorption



T_{gas} varies significantly

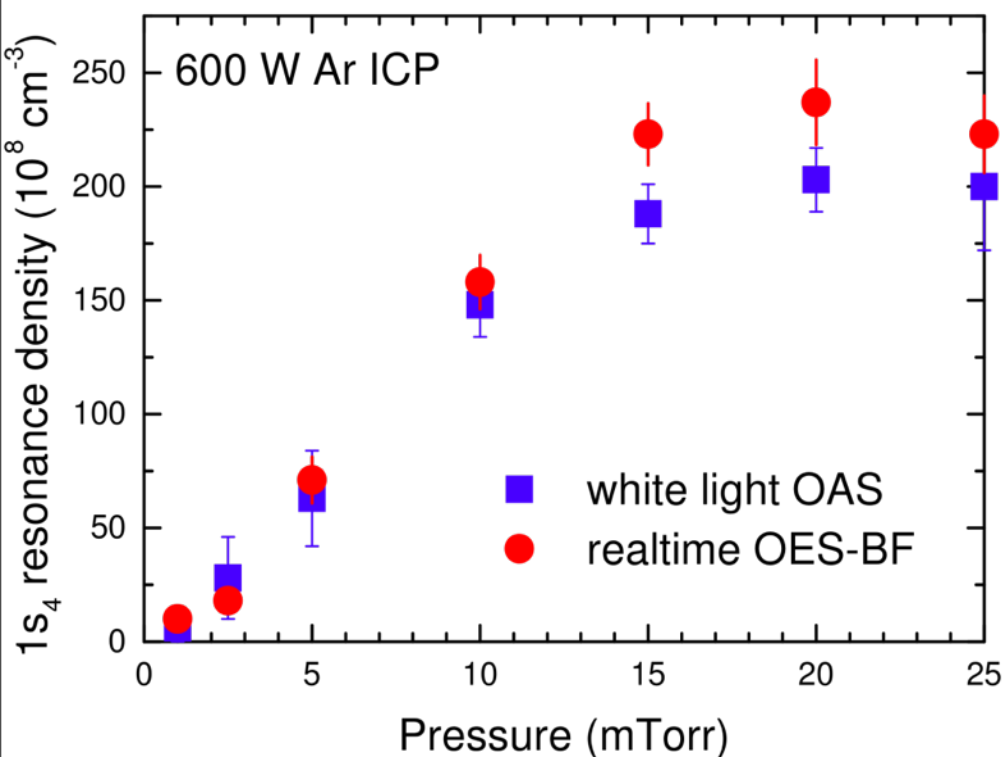


Ground state atoms reabsorb VUV

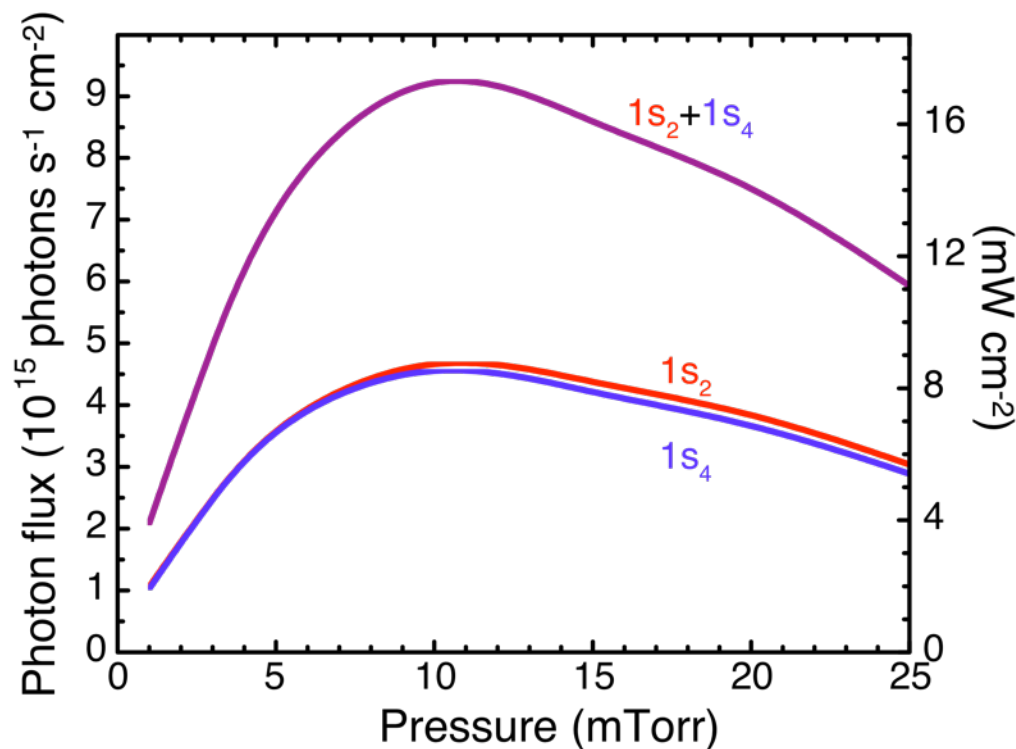
Gas density $n = p/kT_{\text{gas}}$ decreases as T_{gas} increases

Pressure scan - Ar

Resonance density

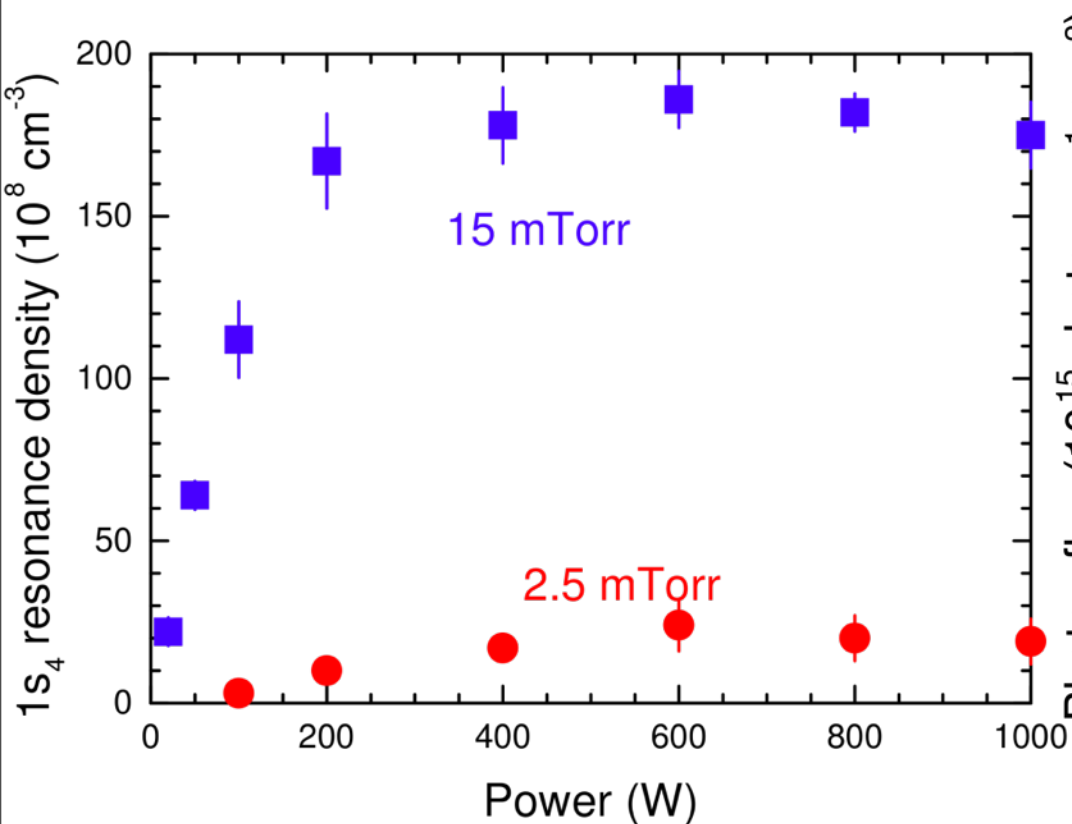


VUV (predicted by model)

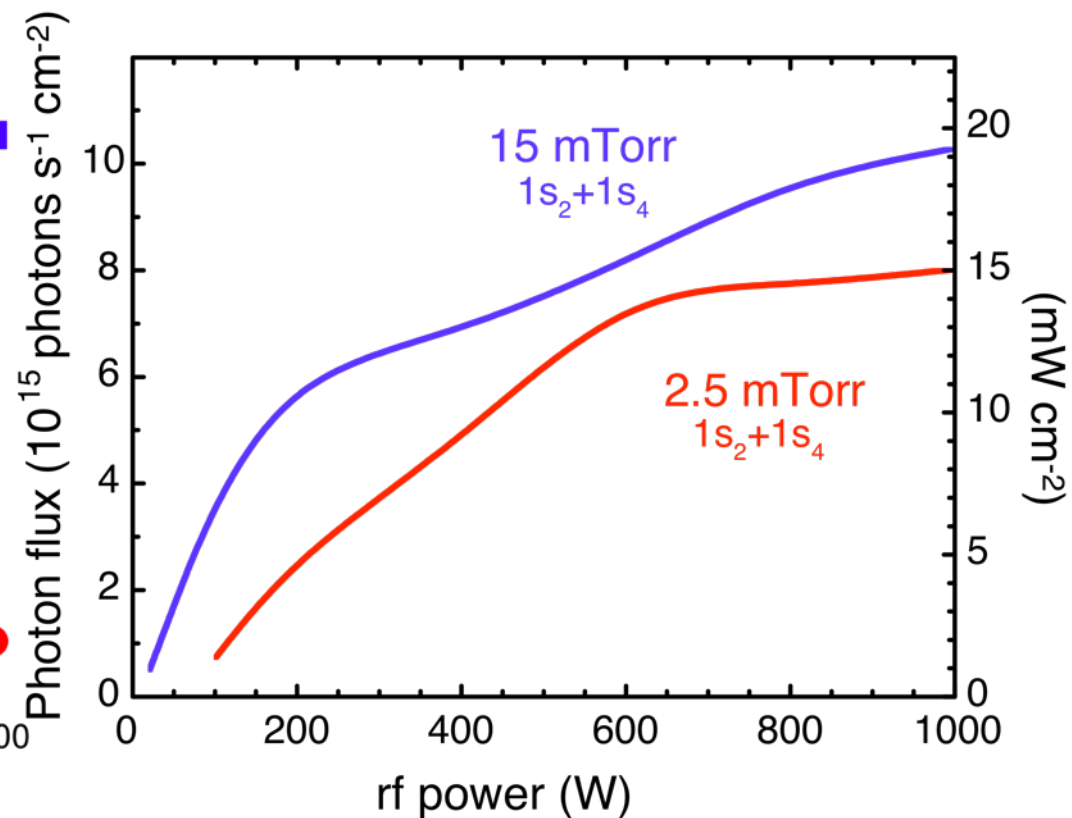


VUV flux decreases with pressure due to enhanced reabsorption.

Resonance density

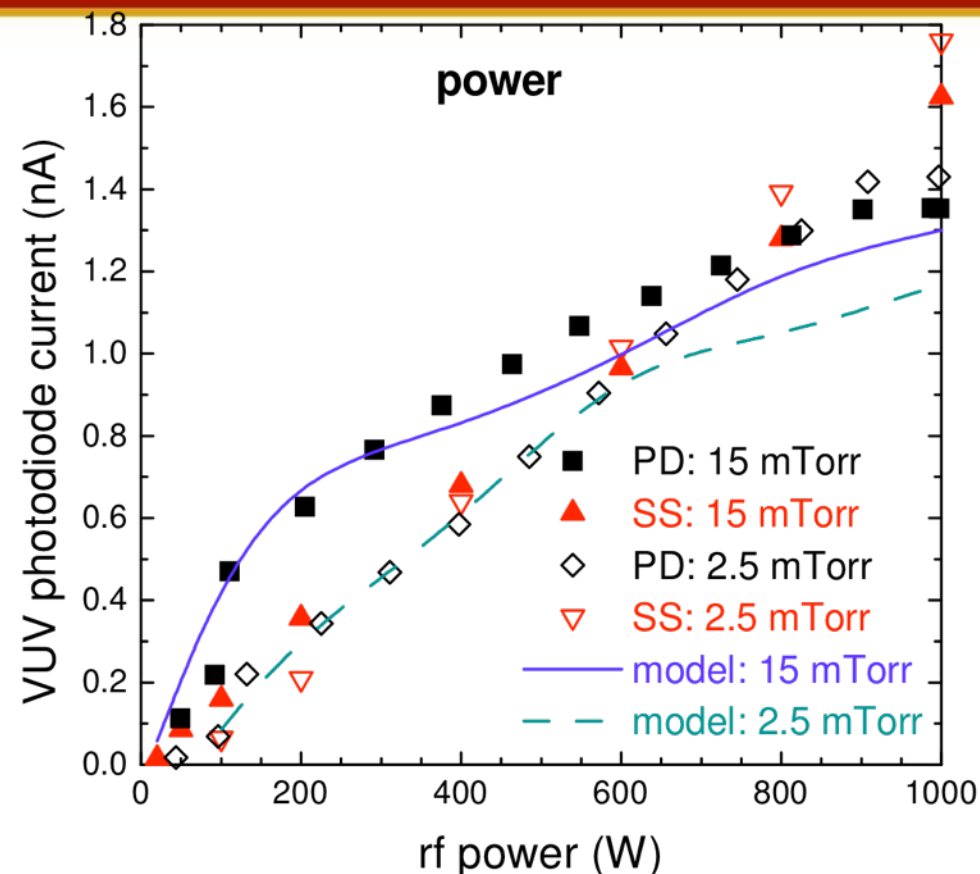
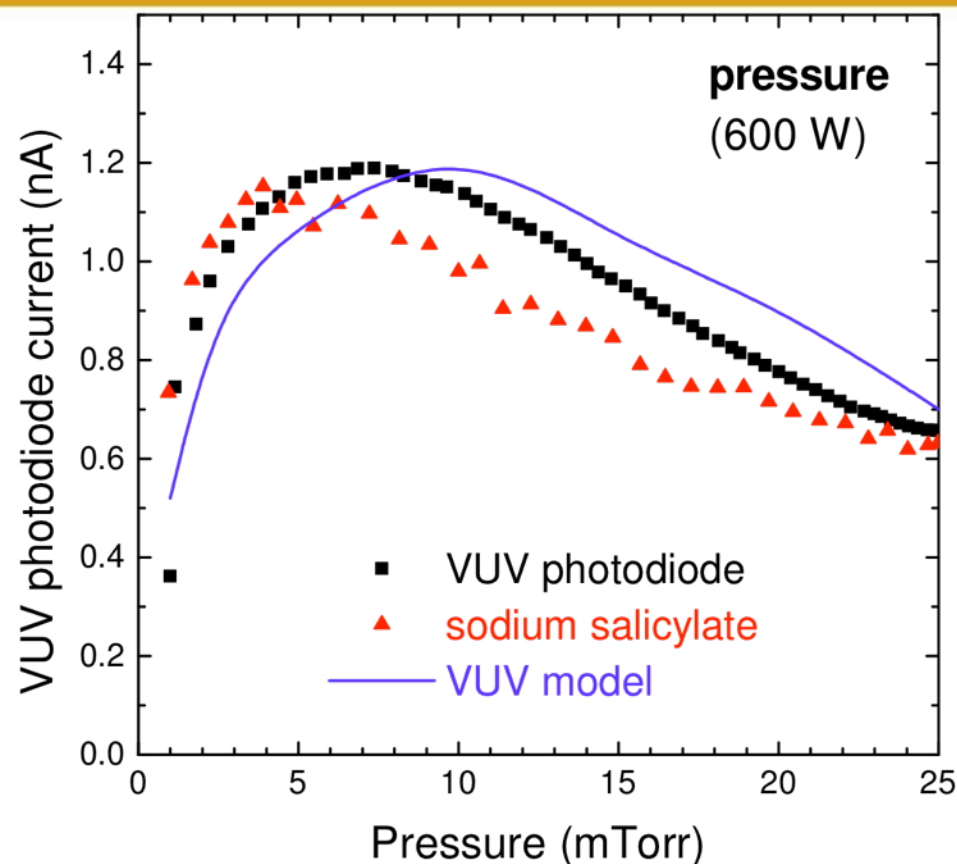


VUV prediction



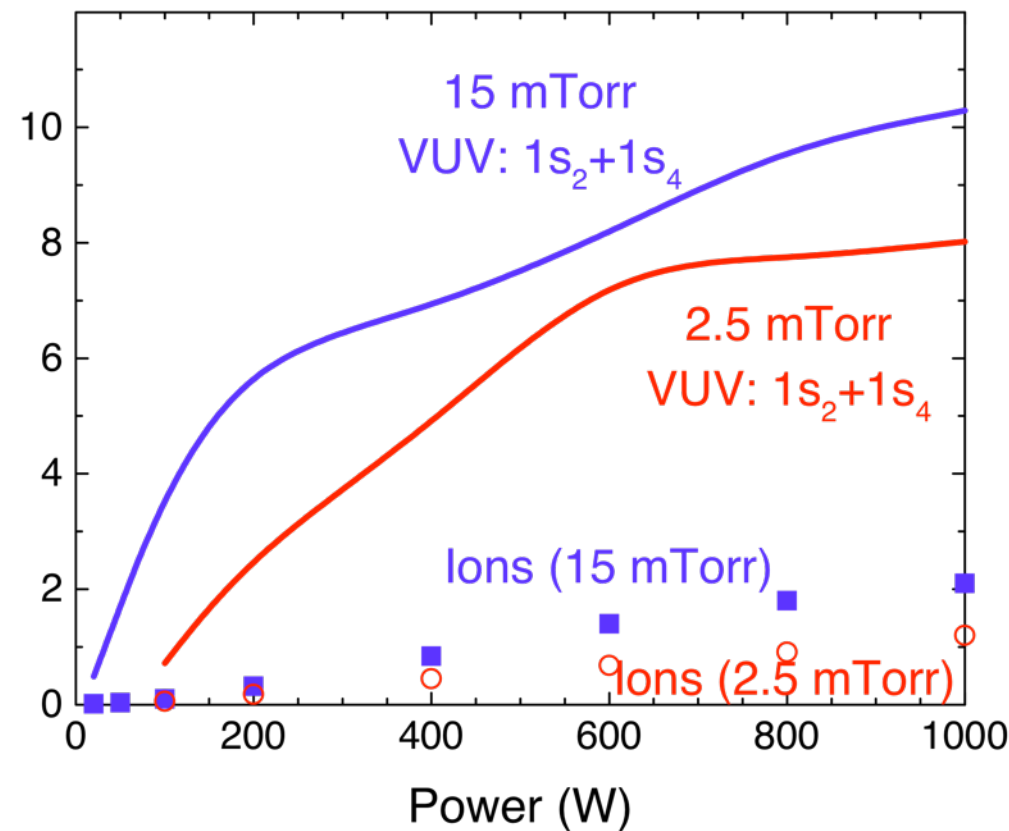
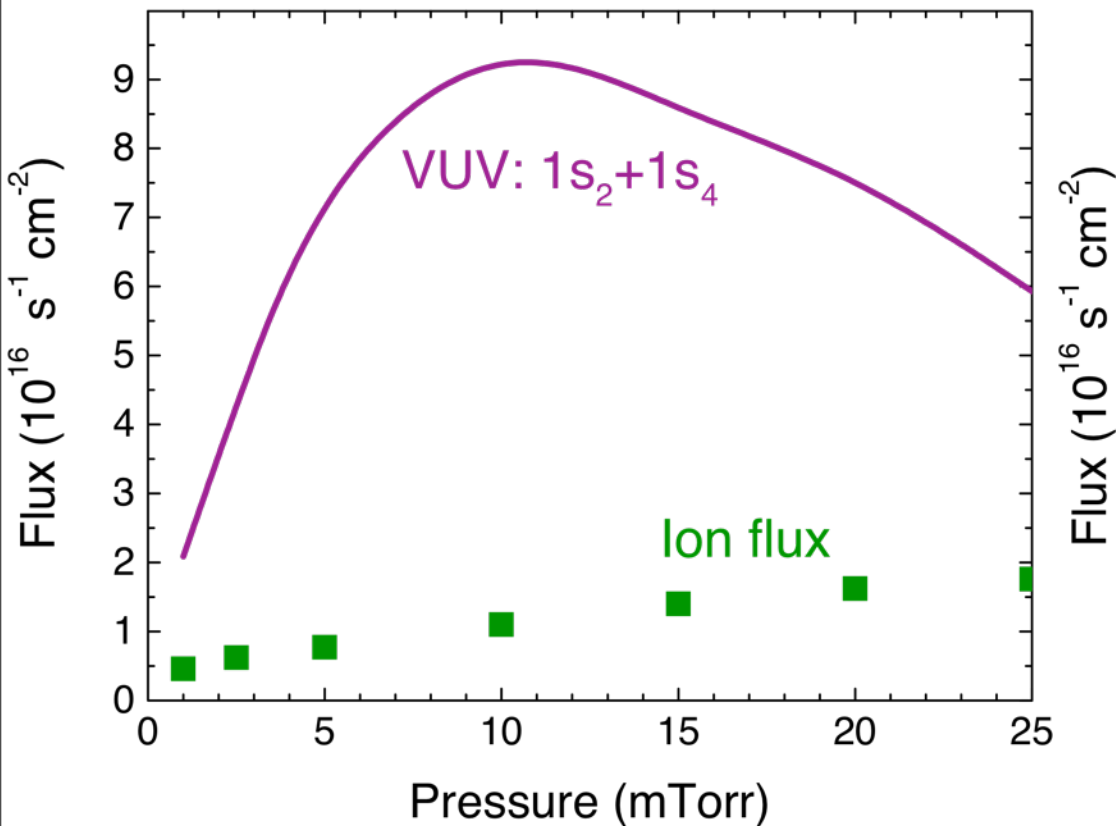
VUV flux increases with increasing power – rarefaction leads to diminished reabsorption.

VUV photon flux in Ar ICP from photodiode and sodium salicylate agree with model



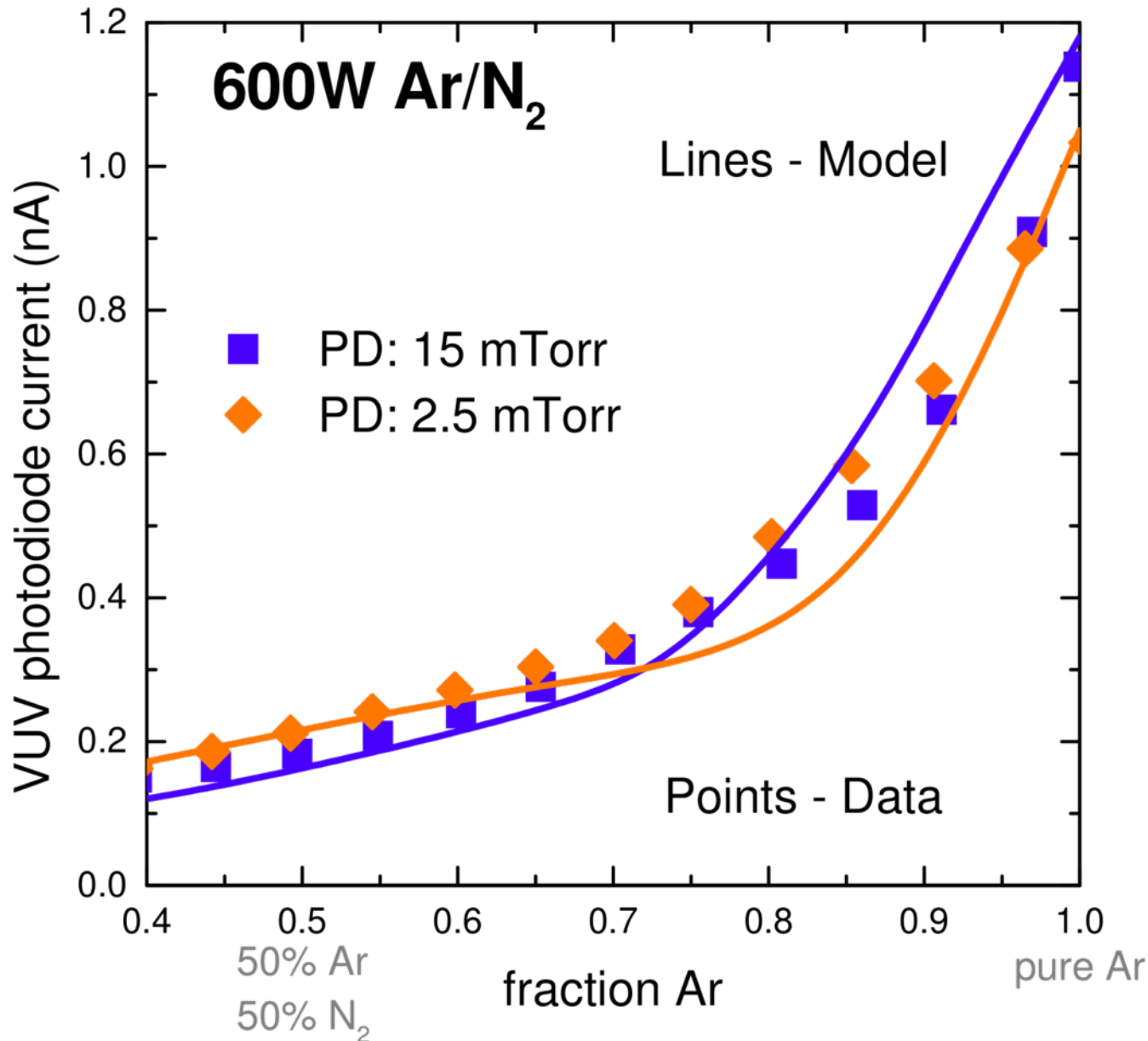
- At higher pressures, radiation reabsorption increases due to increased ground state atom concentration.
- At constant pressure, VUV photon flux increases with power when excited state densities increase.
- Model is in accord with photodiode and sodium salicylate data

VUV flux at electrode location

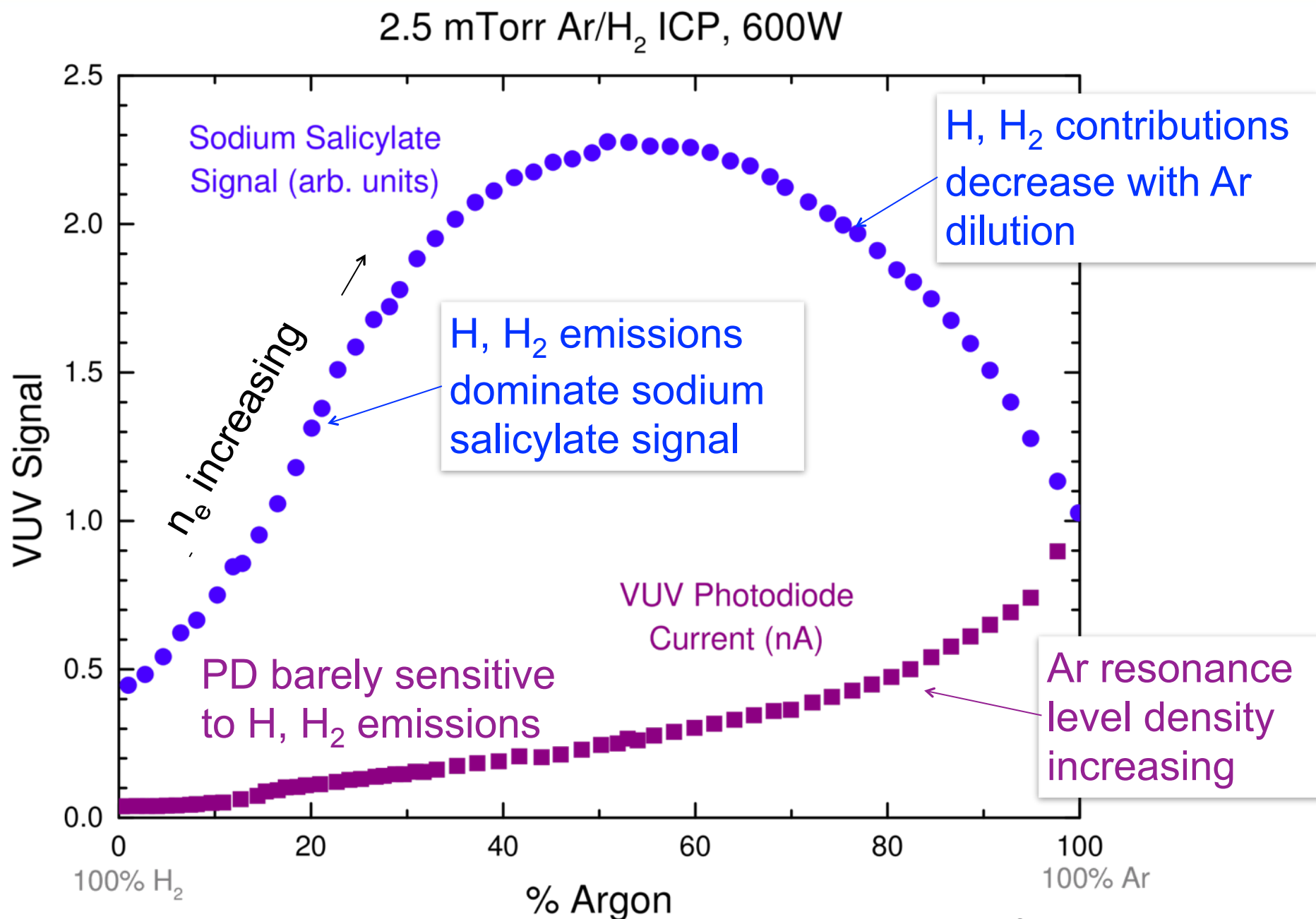


VUV photon flux is greater than ion flux for all conditions

Ar/N₂ plasmas - Ar model works due to no N₂ emissions in VUV range



For Ar/H₂, H, H₂ VUV (90-180 nm) emissions *add* to Ar signals at 104.8 & 106.7 nm



VUV: concluding remarks

- Non-invasive VUV detection has been demonstrated
- A radiation model has been developed to infer VUV fluxes:
 - Inputs: measured resonance level density, Ar pressure, gas temperature
 - The model accounts for VUV reabsorption in plasma
- VUV fluxes from model in Ar agree with direct measurements using
 - VUV photodiode
 - Sodium salicylate fluorescence
- For Ar/N₂ mixture, the model works well because only Ar contributes to VUV emissions
- For Ar/H₂ mixture, H and H₂ contribute to VUV emissions – relative contributions inferred by combining Ar model and sodium salicylate measurements

- OES offers a promising approach to plasma characterization
 - Non-Maxwellian EEDFs, electron density and effective temperature
 - metastable concentrations
 - real-time plasma parameters
 - VUV
 - method is
 - ♦ non-invasive
 - ♦ applicable in rare gas containing plasmas
 - ♦ applicable in conditions where Langmuir probes impractical
- Successfully applied to low-pressure ICP - what about higher pressure?