

INTERACTIONS BETWEEN ATMOSPHERIC PRESSURE PLASMAS AND COMPLEX SURFACES*

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**International Online Plasma Seminar
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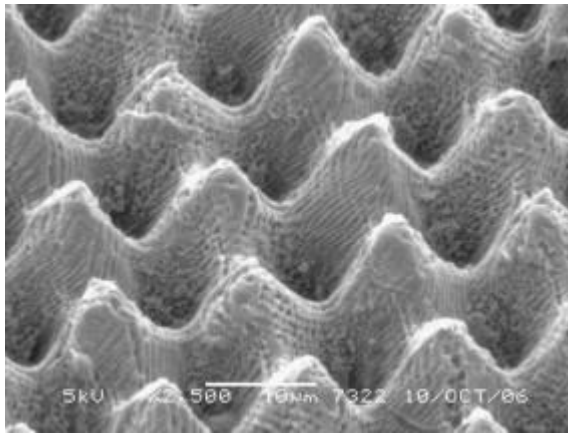
*** Work supported by DOE Fusion Energy Sciences and the National Science Foundation**

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AGENDA

- Synergistic interactions with complex surfaces
- Primer on Surface Ionization Waves
- Description of model
- Surface discharges across wavy and aerosol hosting surfaces
- Surface discharges across porous materials.
- Plasma jets incident onto shallow, wide and deep channels.
- Concluding Remarks

PLASMA SYNERGIES WITH COMPLEX SURFACES



- Atmospheric pressure plasmas (APPs) synergistically interact with complex surfaces.
- In this talk, complex means non-planar, non-uniform, porous dielectrics (wet and dry).
- Just as with metals, non-planar dielectrics having large ϵ_r produce electric field enhancement that intensify the plasma.
- Unlike metals, dielectric have memory through charging of their surfaces.
- Charging (and secondary electron emission) produce synergistic effects.
- As in low pressure plasmas, aspect ratio of features determine plasma penetration and properties.

<https://fineartamerica.com/featured/1-water-drops-on-a-rubber-surface-daniel-sambrausscience-photo-library.html>

www.utwente.nl/en/et/ms3/research-chairs/WAoud_niets_uit_wissen_aub/research/laser/probic/#resultsjects_archive/hydropho

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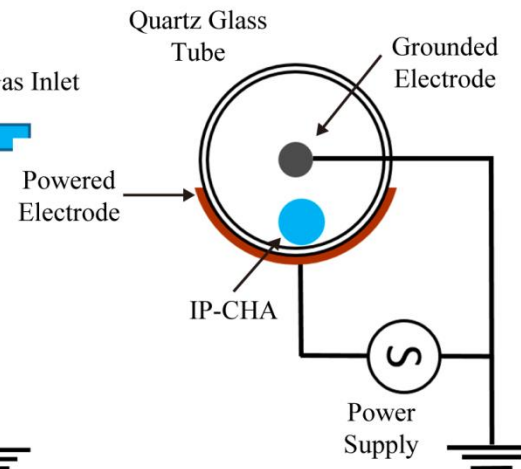
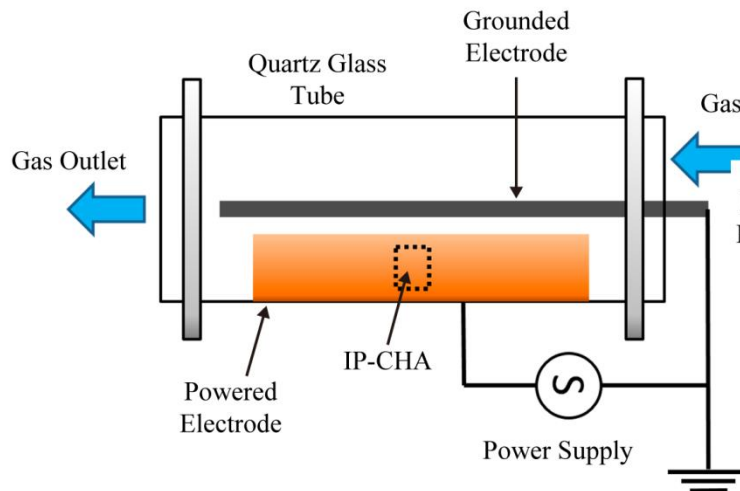
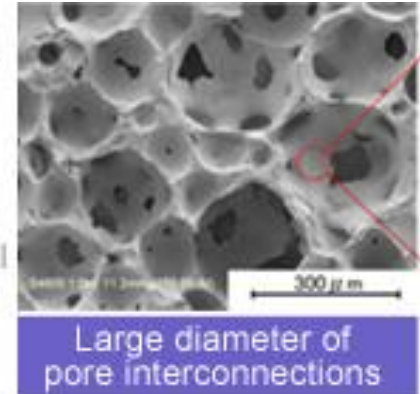
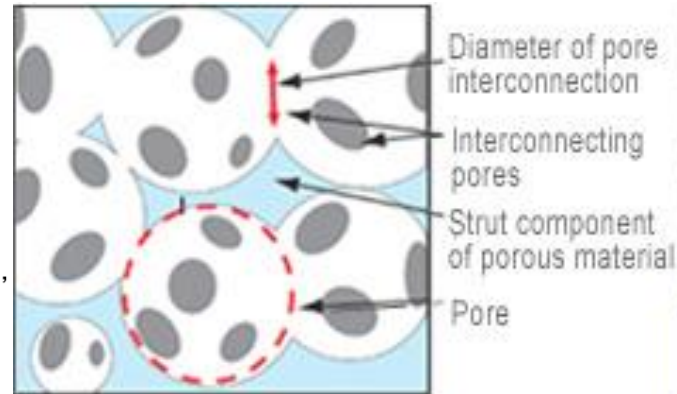
PLASMA TREATMENT OF BONE SCAFFOLDING

- Bone scaffolding – calcium dominated, porous dielectric used as bone replacement. Bone cells grow into pores, eventually replacing scaffold.
- Scaffolding often needs treatment to improve biocompatibility.

Impact of non-thermal plasma surface modification on porous calcium hydroxyapatite ceramics for bone regeneration

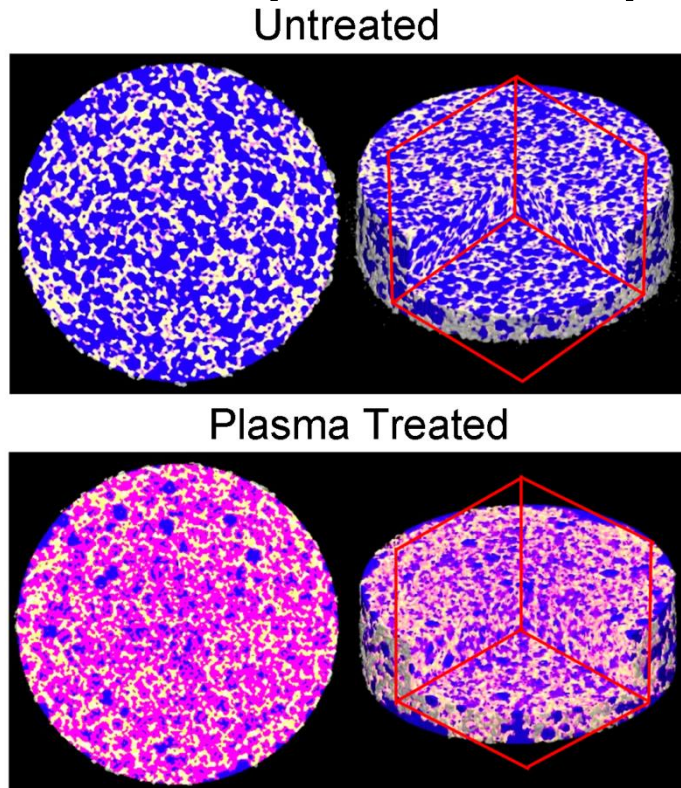
Yu Moriguchi , Dae-Sung Lee , Ryota Chijimatsu , Khair Thamina , Kazuto Masuda , Dai Itsuki , Hideki Yoshikawa , Satoshi Hamaguchi , Akira Myoui

PLOS ONE | <https://doi.org/10.1371/journal.pone.0194303> March 14, 2018



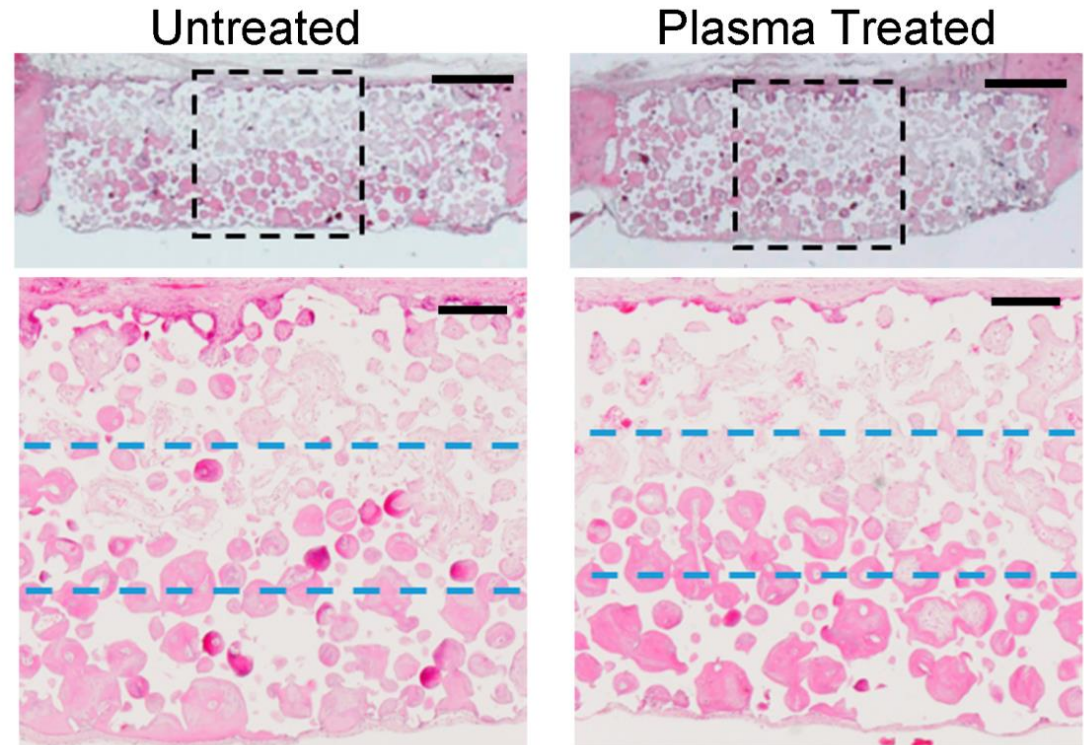
PLASMA TREATMENT OF BONE SCAFFOLDING

- Low pressure oxygen plasma treatment of NeoBone (150 μm pores, 40 μm openings) increases contact angle, wettability and cell proliferation into scaffolding.
- Can such processes be performed faster and cheaper using APPs?



- **Wettability**

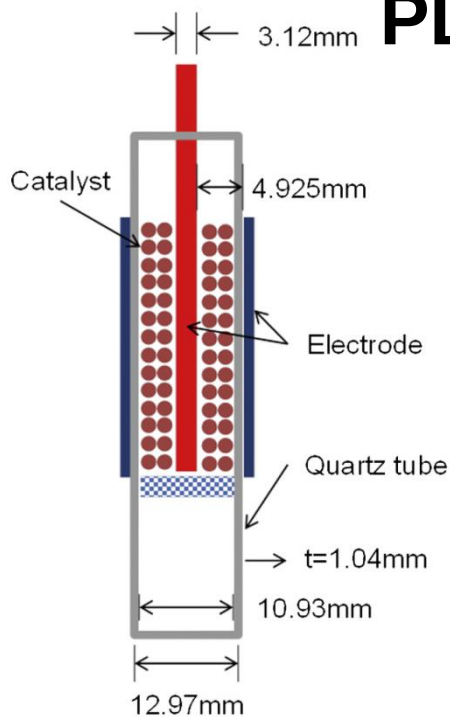
- Y. Moriguchi, et al. PLOS One, 2018



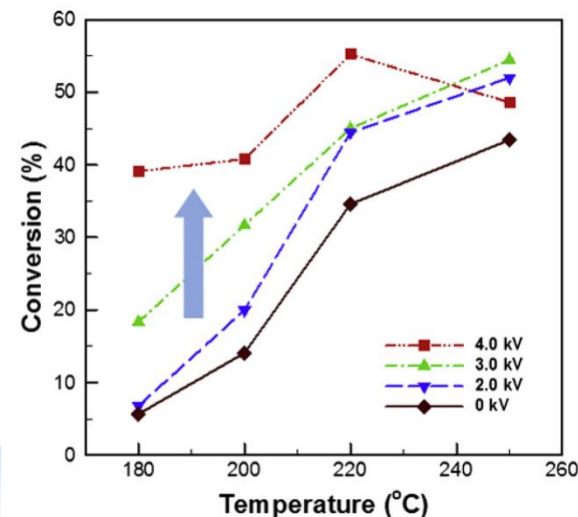
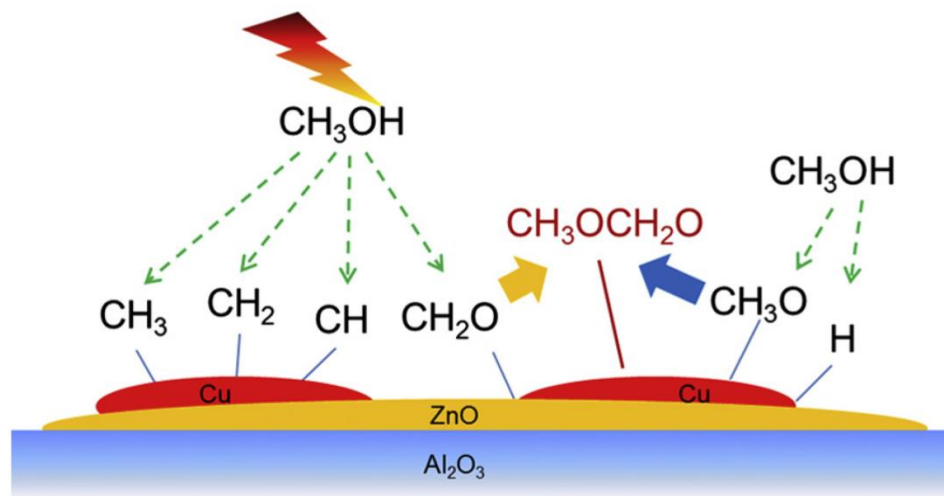
- **Cell proliferation**

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PLASMA CATALYSIS: CH₃OH REFORMING



- Packed bed plasma reactor with Cu/ZnO on Al₂O₃ catalyst improves the efficiency of methanol reforming.
- Higher conversion at lower temperature enabled by plasma formed radicals.
- Poorly understood interactions with porosity of dielectrics.



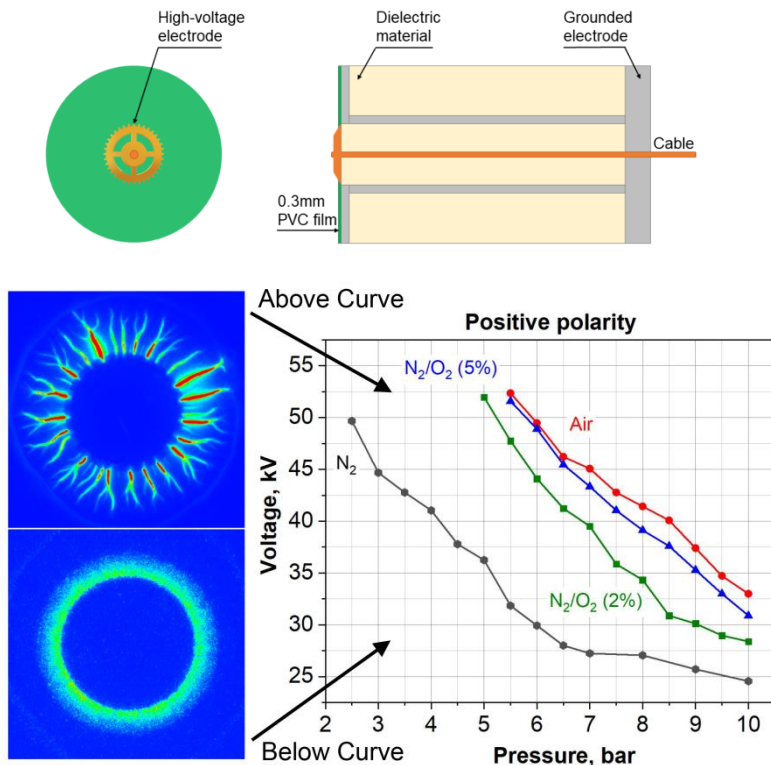
Synergetic mechanism of methanol-steam reforming reaction
in a catalytic reactor with electric discharges

Taegyu Kim, Sungkwon Jo, Young-Hoon Song, Dae Hoon Lee
Applied Energy 113 (2014) 1692–1699

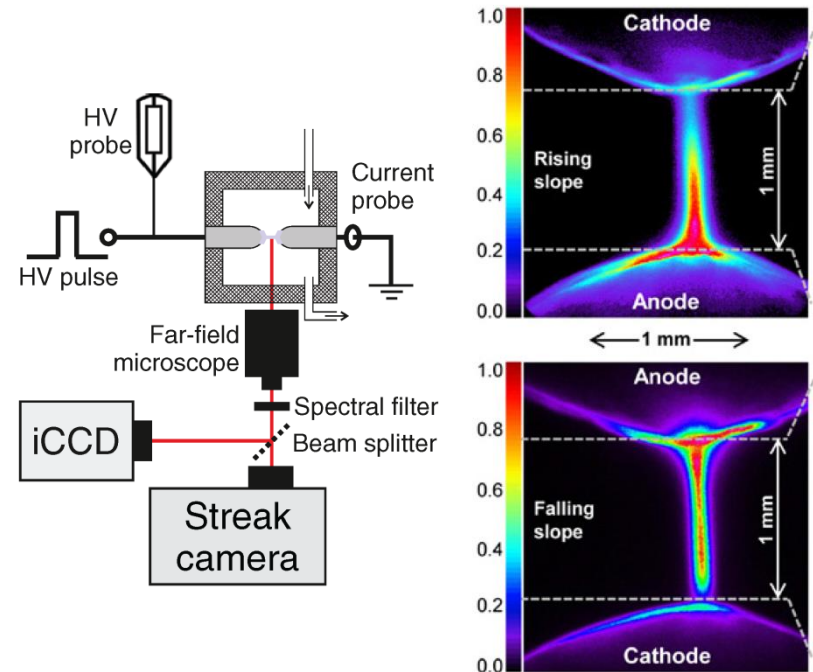
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APPs IN CONTACT WITH SURFACES

- The majority of these applications involve direct contact of APPs with surfaces (solid, soft, liquid).
- One mode of such interactions involve “surface hugging” plasmas or “surface ionization waves” (SIWs) – either deliberate or natural outcome.



- Modes of Deliberate SIW (S. Starikovskaia, AIAA Scitech 2019 Forum)

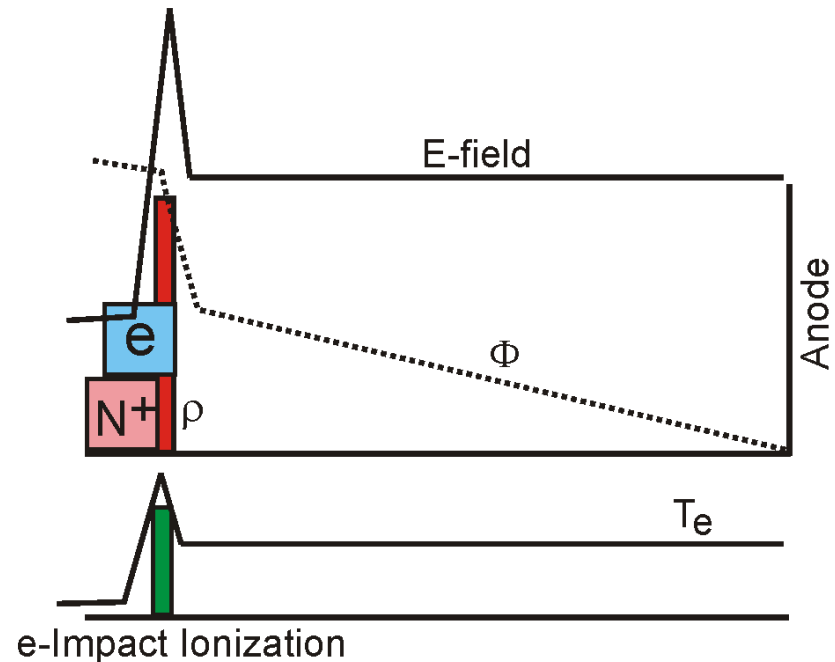
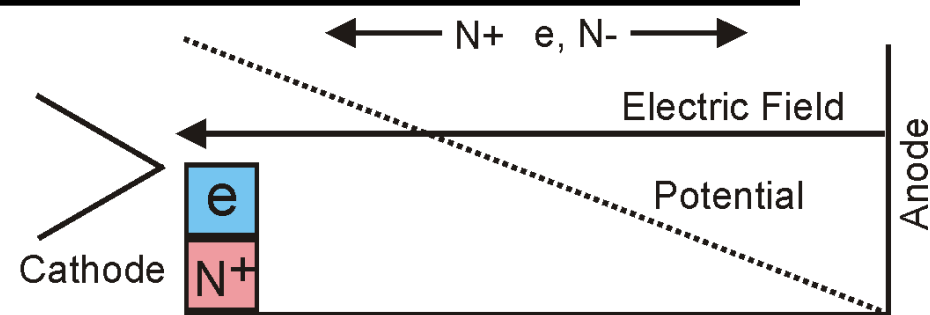


- Natural transition to SIW (Hoft and Brandenburg, JPD 47, 465206 (2014)).

PRIMER ON SURFACE IONIZATION WAVES

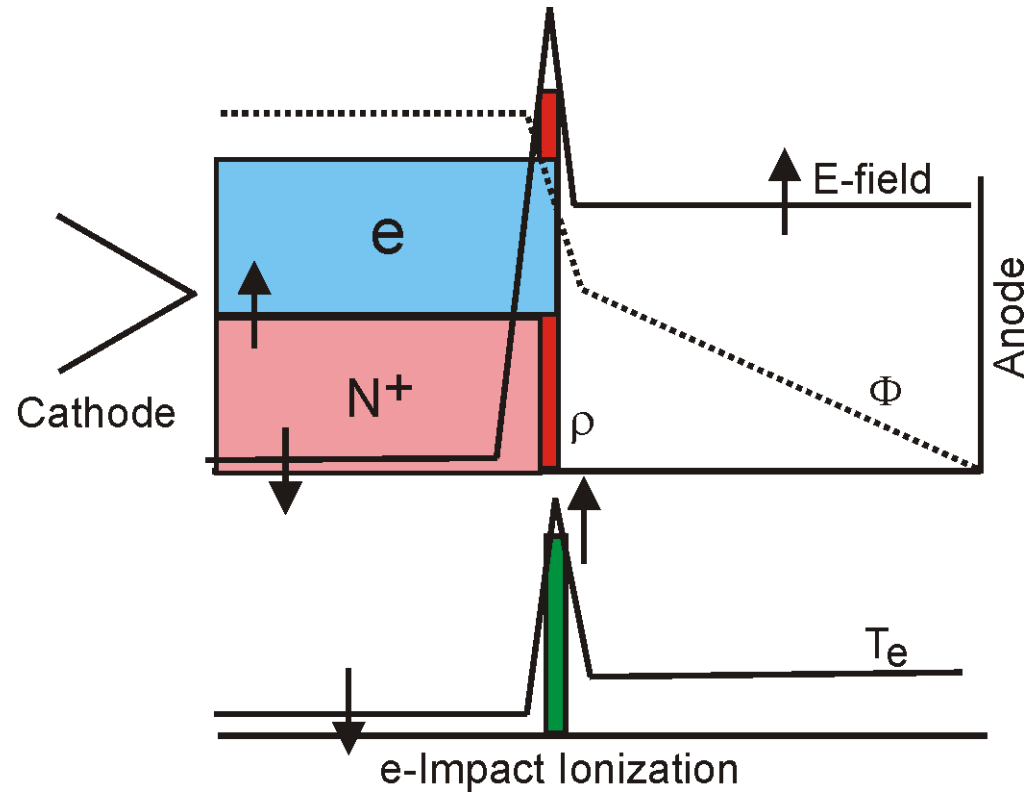
FUNDAMENTALS OF NEGATIVE “BULK” IONIZATION WAVES

- Large negative voltage on cathode – grounded anode.
- Small spot of plasma at tip of cathode (field emission?).
- Electrons drift towards anode – ions towards cathode.
- Negative space charge creates spike in E-field.
- E-field heats electrons producing spike in T_e and ionization.
- Avalanche front develops, increasing plasma density.



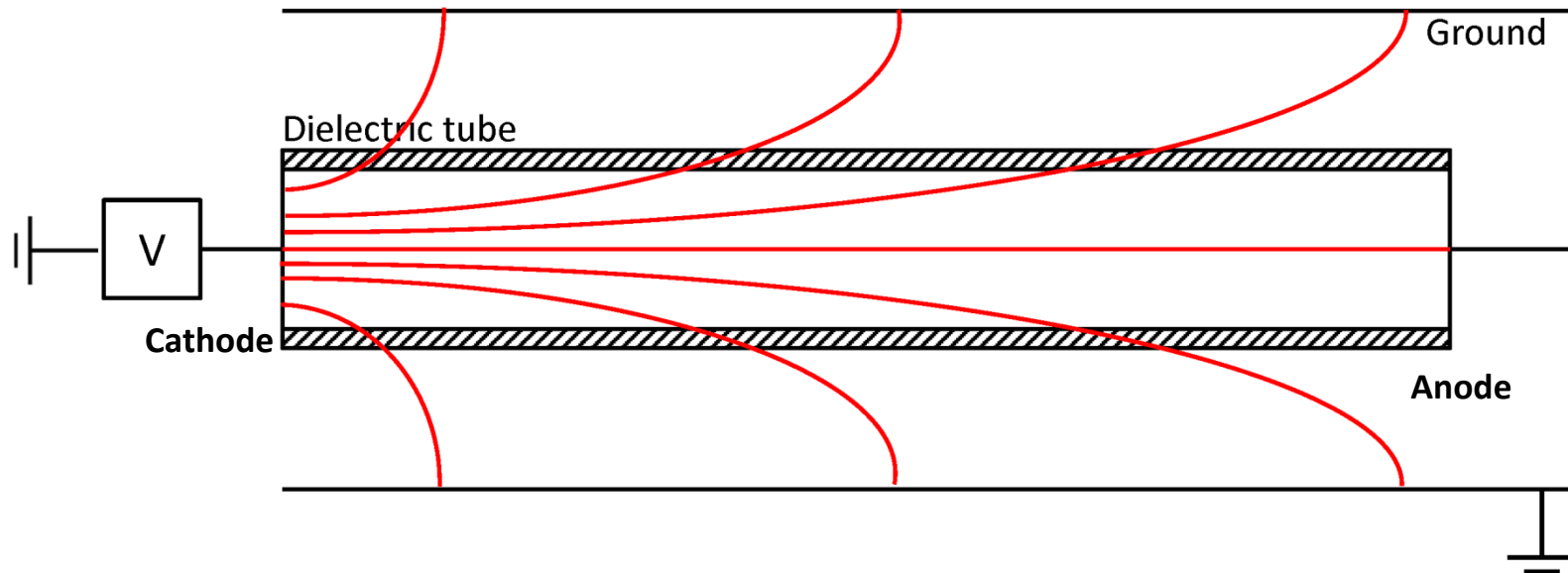
FUNDAMENTALS OF NEGATIVE STREAMERS

- Avalanche front produces a column of high plasma density.
- High conductivity in column reduces E-field, voltage drop and T_e .
- Voltage compressed ahead of streamer, increasing E.
- Electrons drift towards anode – ions towards cathode.
- Negative space charge creates spike in E-field.
- E-field produces spike in T_e and ionization – avalanche front extends plasma column.



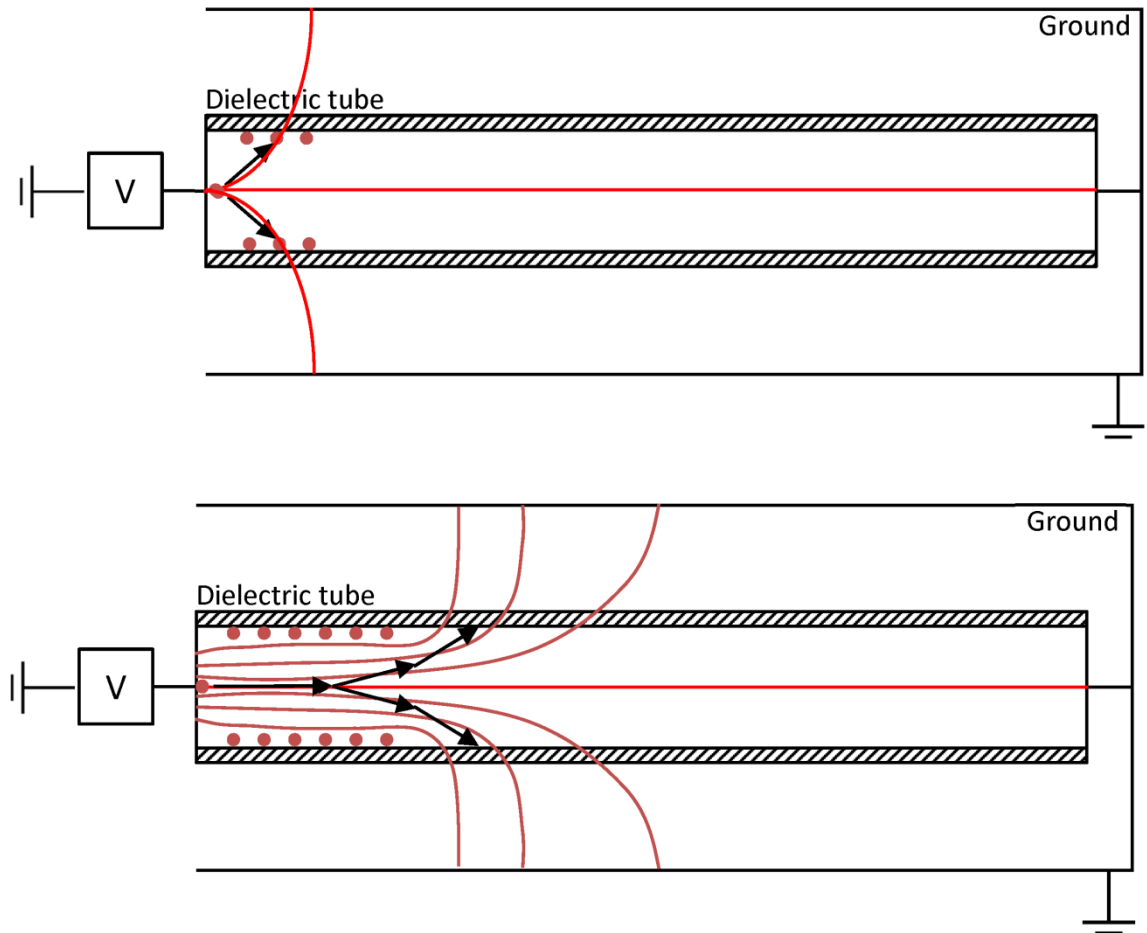
BREAKDOWN PROCESS – INITIAL E-FIELD

- Consider electrical breakdown in a fluorescent lamp-type discharge sustained in a dielectric tube.
- Initially, electric field lines are directed from cathode to *where-ever ground can be found*, here to the coaxial ground.
- The closest ground plane is rarely the anode – only electric field lines near the axis terminate on the anode



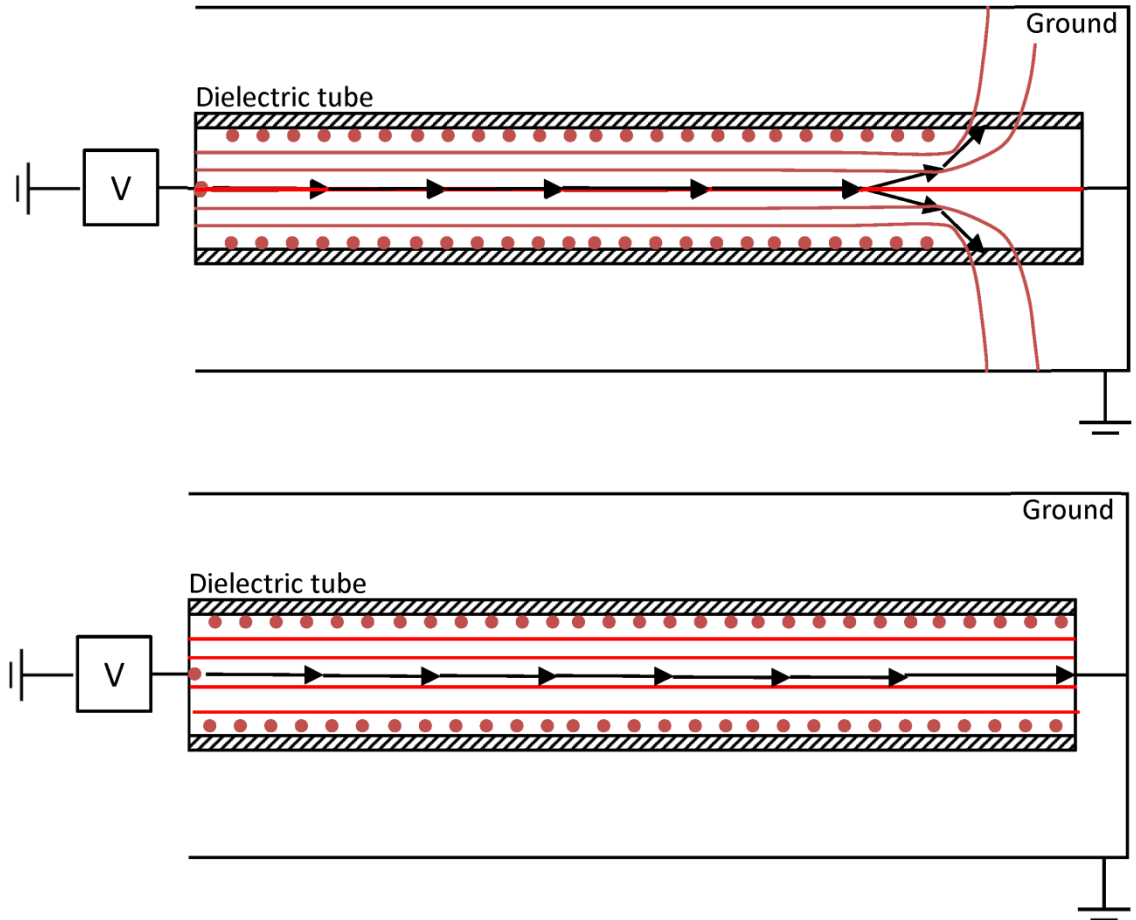
BREAKDOWN PROCESS – REDIRECTING E-FIELD

- Initial electrons from cathode follow vacuum electric field lines which intersect walls.
- Electrons charge wall, which shields potential and redirects E-field parallel to wall.
- Electrons drift axially to end of charged region, where they follow E-field to further charge wall.



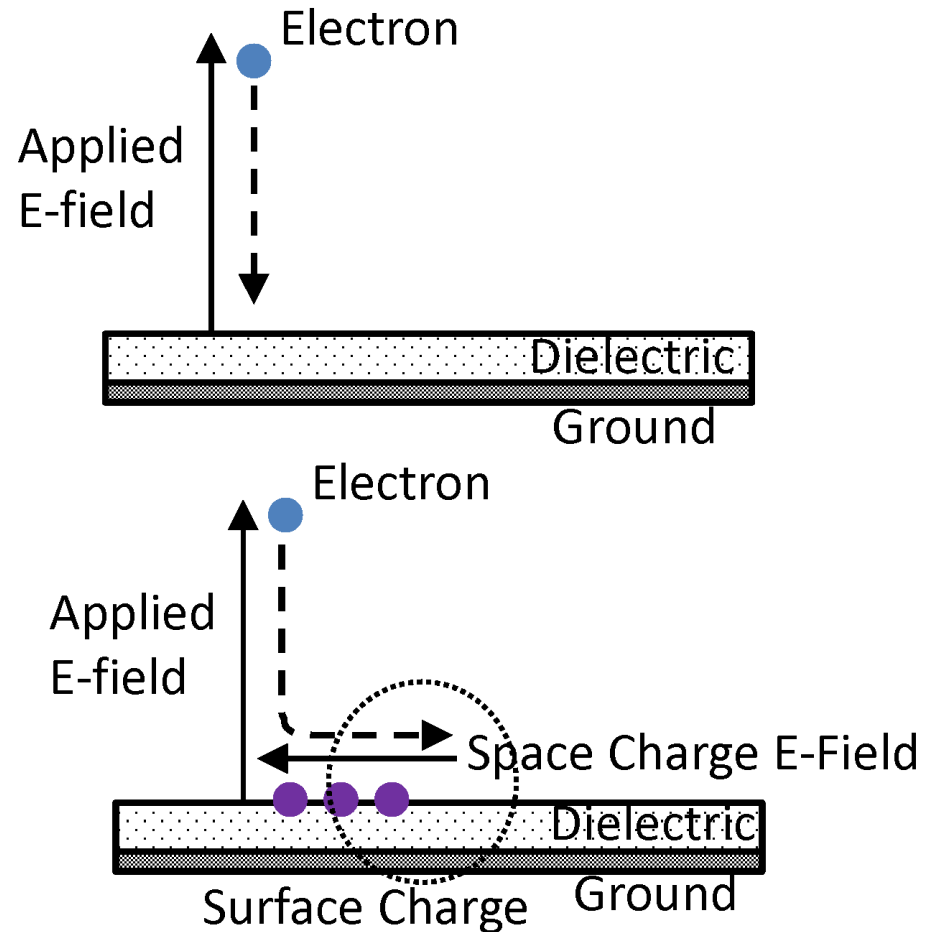
BREAKDOWN PROCESS – CONNECTING TO ANODE

- Process of wall charging, turning E-field parallel to wall and electrons further charging wall continues along length of tube.
- Eventually, E-field lines seek out the cathode is closest ground plane.
- With full wall charging, coaxial ground is shielded from plasma.



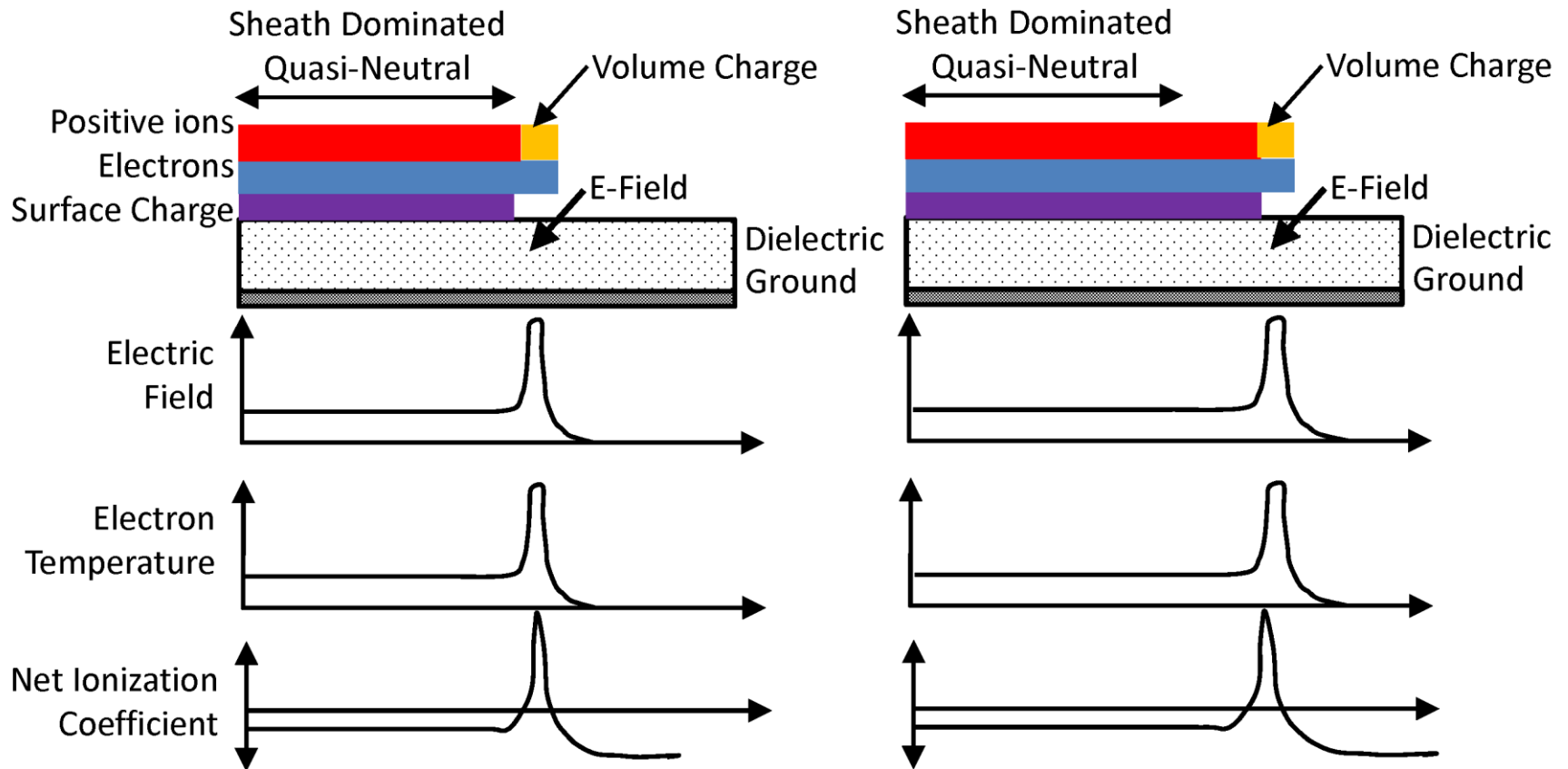
SURFACE IONIZATIONS WAVES

- Surface Ionization Waves (SIWs) are streamer-like extensions of the plasma column that *hug* the dielectric.
- Space charge enhancement at the front of the SIW produces electric fields that avalanches the gas.
- Surface charging produces a sheath that shields underlying ground.
- Volume charge in head of stream produces peak in T_e .
- For negative SIW, process starts with external electric field driving electrons into surface,



SURFACE IONIZATIONS WAVES

- Electric field in surface-hugging plasma consumes voltage.
- SIWs will continue to propagate as long as there is sufficient voltage remaining.



E-FIELD AND POTENTIAL: DIELECTRICS

- At surface of a dielectric and a gas:

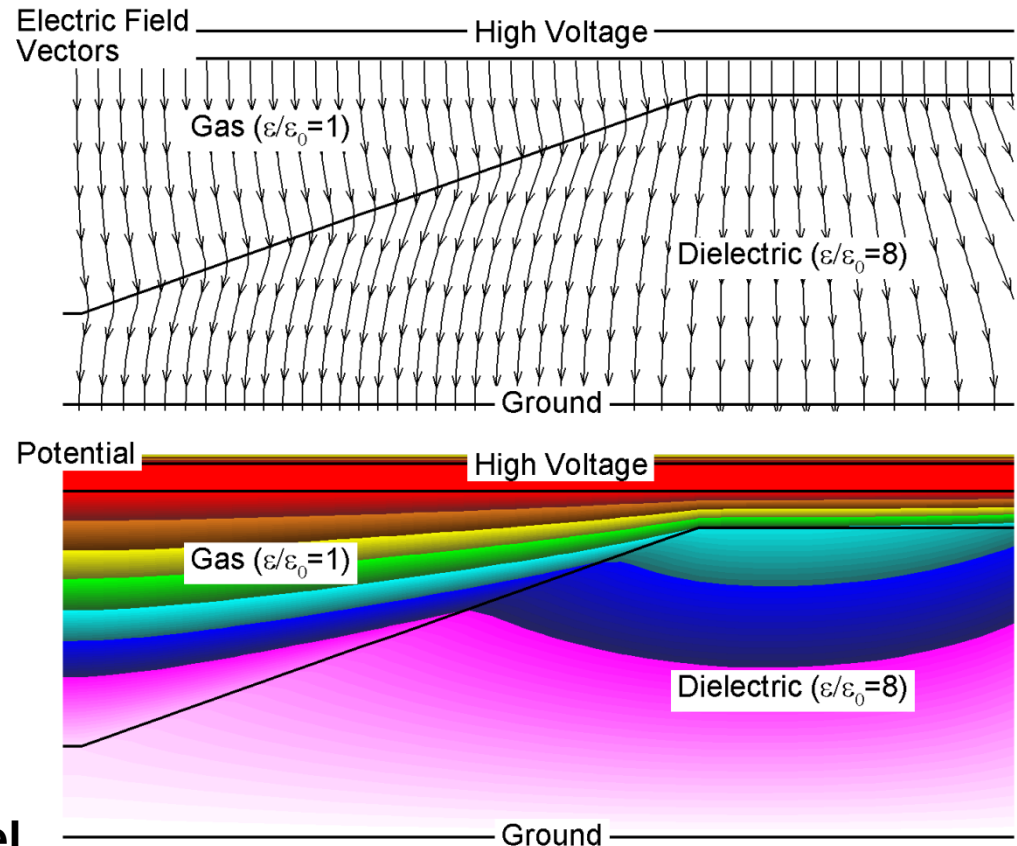
$$D_{\perp gas} - D_{\perp dielectric} = \rho_{surf}$$

$$\epsilon_g E_{\perp gas} = \epsilon_d E_{\perp dielectric} \quad (\rho_{surf} = 0)$$

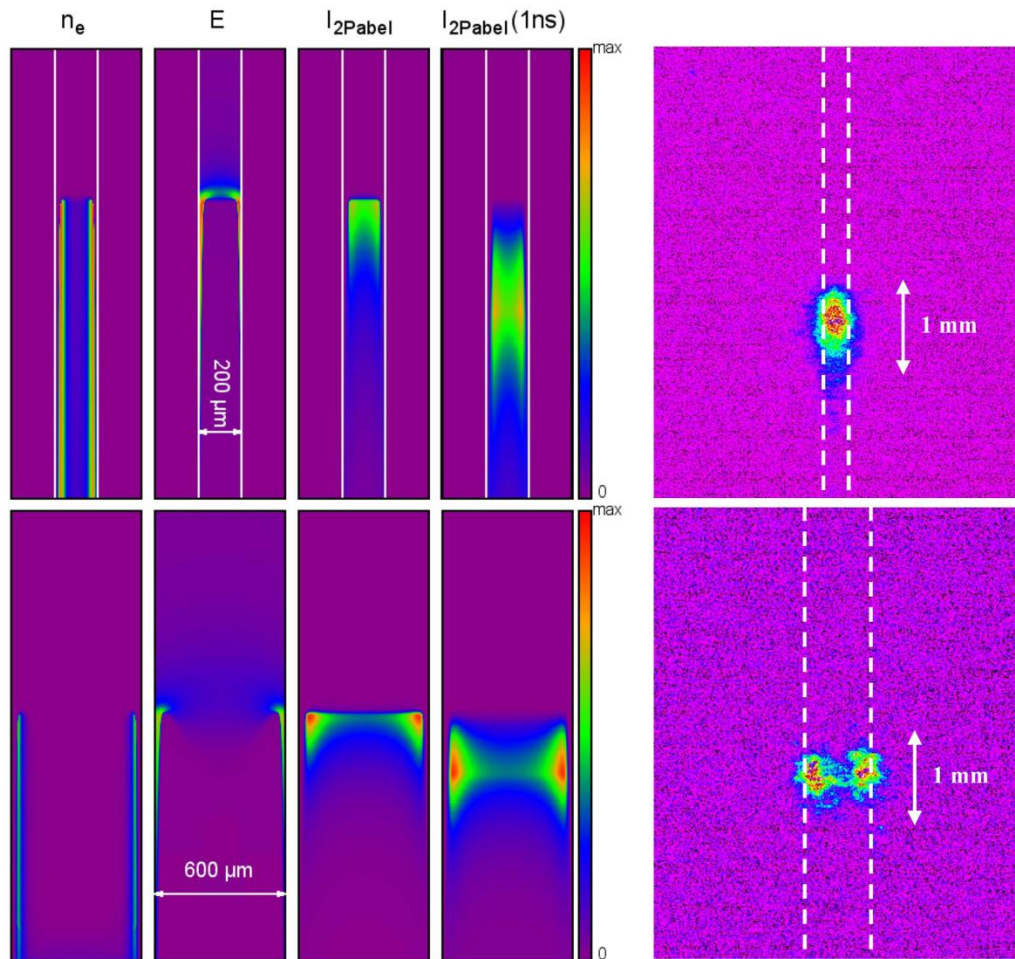
$$E_{\parallel gas} = E_{\parallel dielectric}$$

$$\nabla \cdot \vec{E} = \rho_v$$

- Electric potential is expelled from high ϵ materials (solids) into low ϵ materials (gas).
- Electric field lines are bent at surfaces producing electric field enhancement.
- Surface charging produces parallel components of electric field.
- Result is *surface hugging ionization waves*



TIGHTLY BOUNDED APP STREAMERS



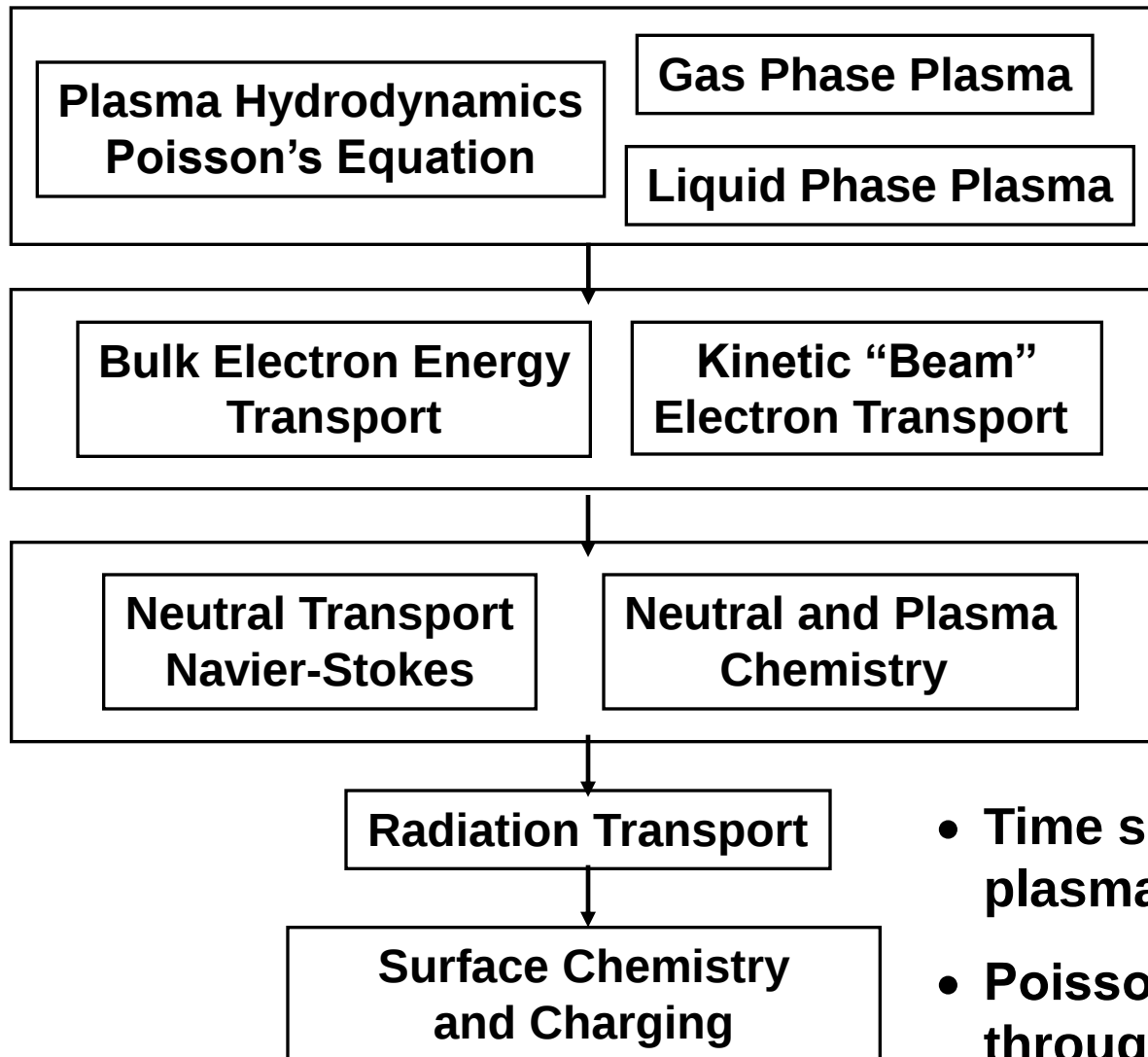
- APPs also interact with their boundaries
- An extreme case is capillary discharges, where the channel width of the channel is commensurate with the streamer.
- Streamers often hug the boundaries as surface discharges.

Propagation of an Air Discharge at Atmospheric Pressure in a Capillary Glass Tube

Jaroslav Janský, Pierre Le Delliou, Fabien Tholin, Zdeněk Bonaventura, Pierre Tardiveau, Anne Bourdon, and Stéphane Pasquiers

IEEE TRANSACTIONS ON PLASMA SCIENCE

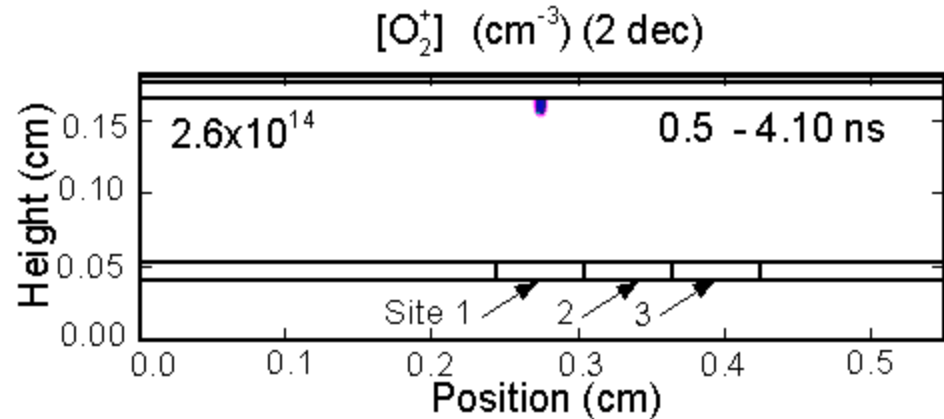
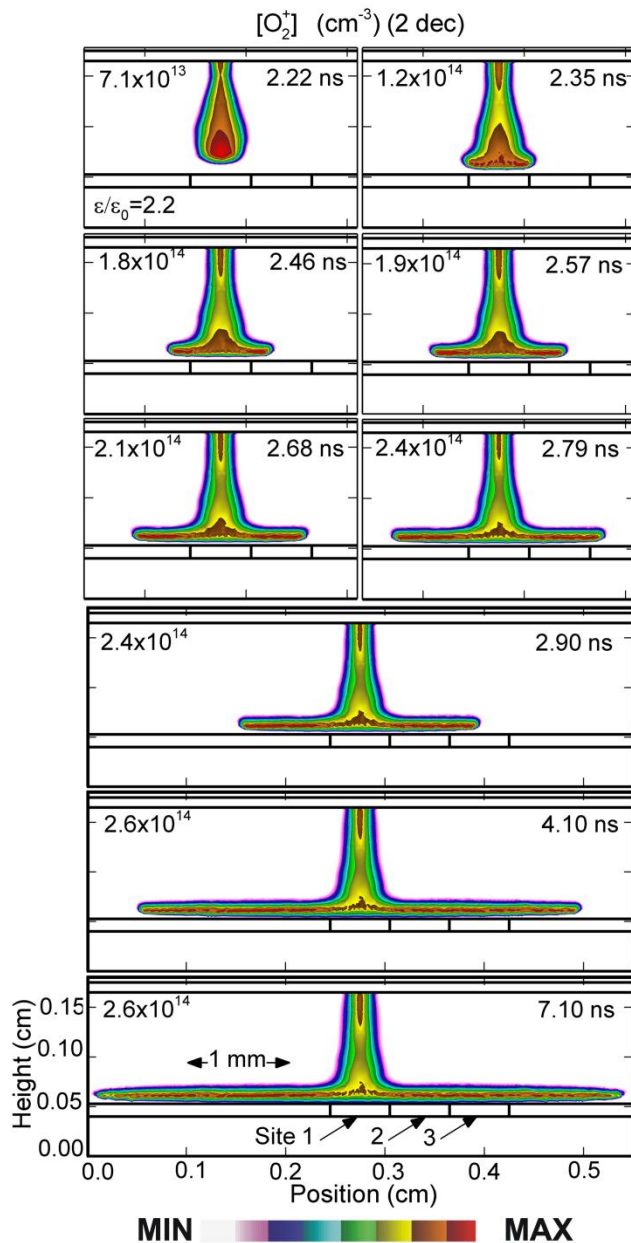
DESCRIPTION OF MODEL – AND TRANSITION TO SIW



MODEL: *nonPDPSIM*

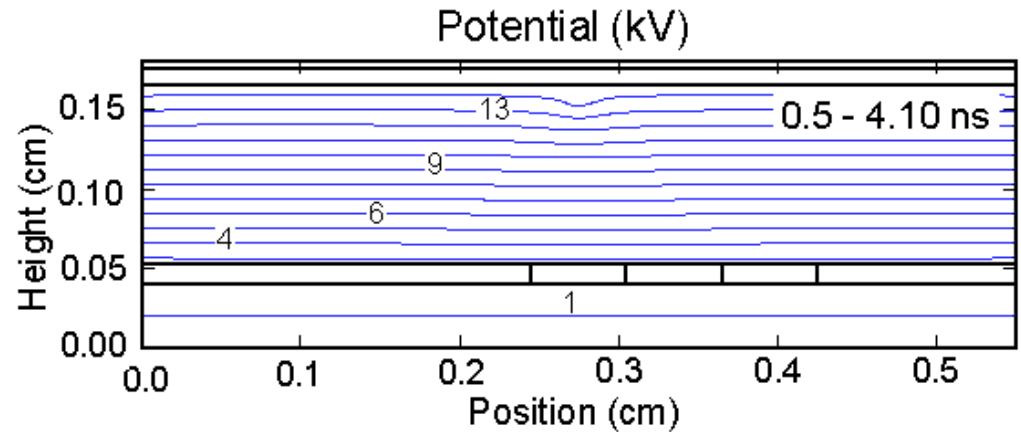
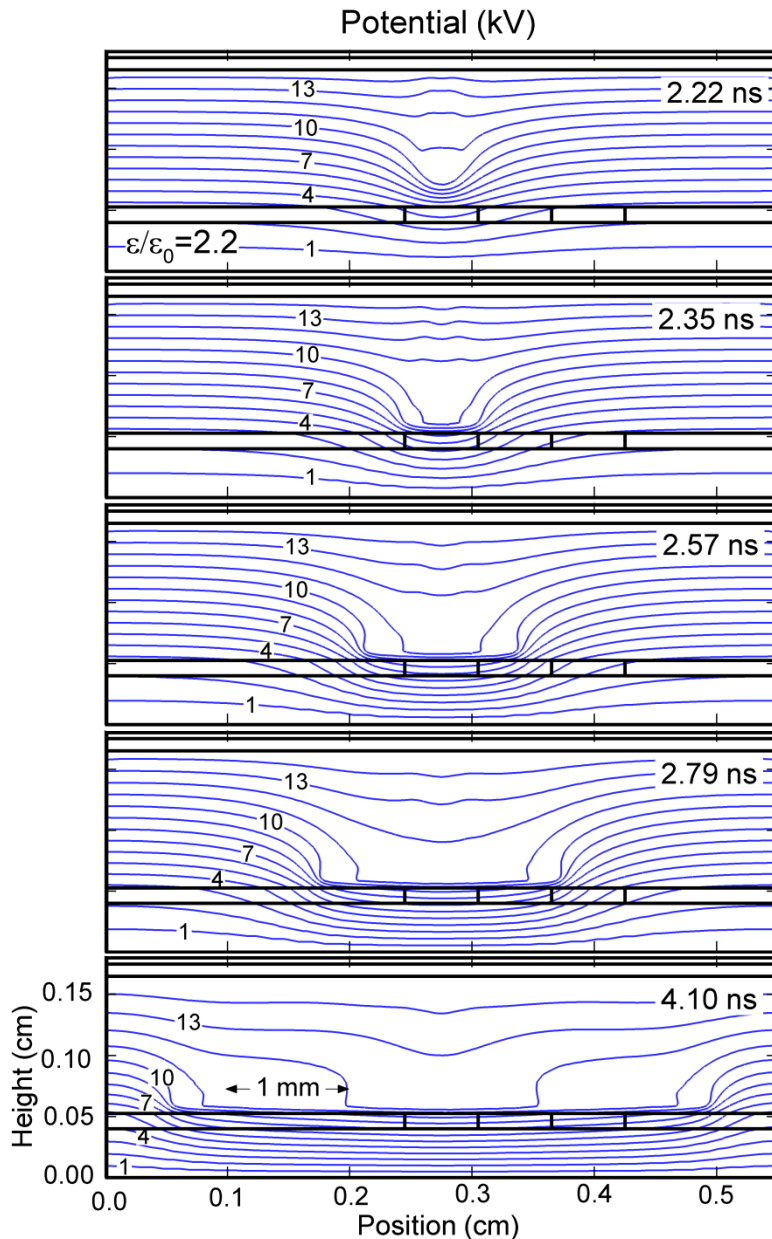
- 2-D unstructured mesh
- Liquid computationally treated identically to gas phase albeit different reactions.
- Time slicing algorithms between plasma and fluid timescales.
- Poisson's equation is solved throughout the computational domain including liquid.

STREAMER DYNAMICS: OXYGEN IONS



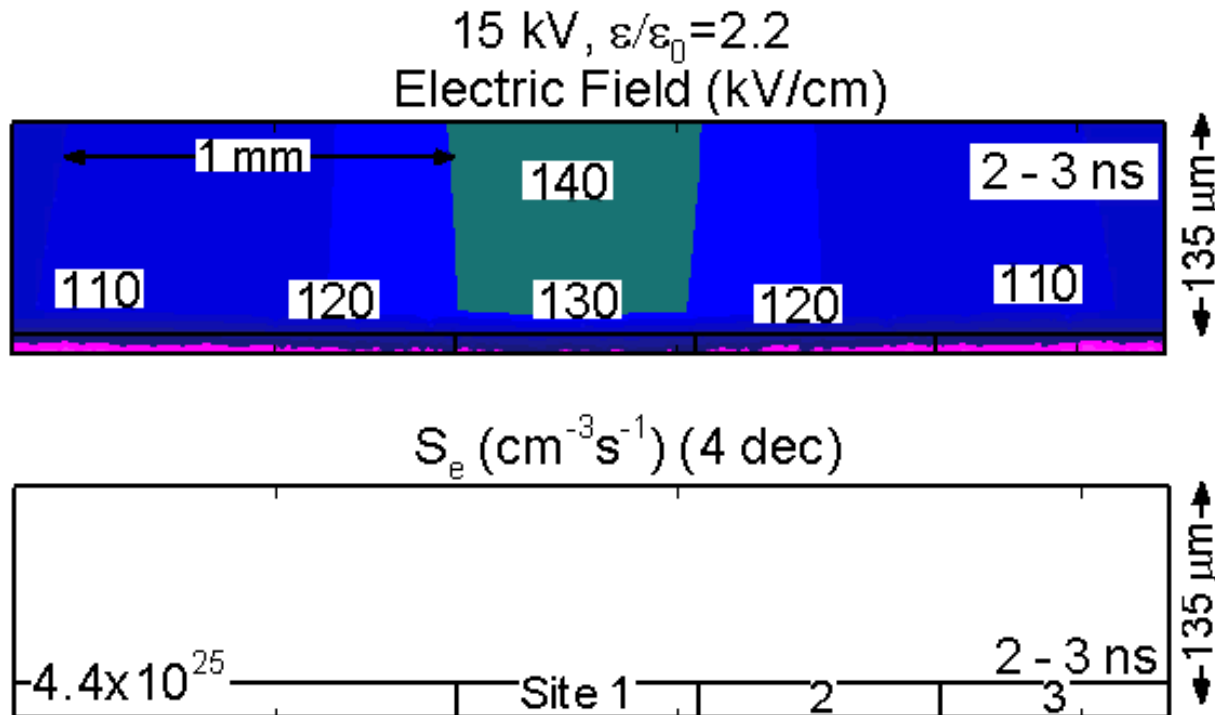
- Electron and ion densities in the streamer channel approach 10^{15} cm^{-3} .
- Upon intersection with dielectric, the streamer charges and spreads along the surface. A wave propagates outward
- Positive corona, $N_2/O_2/H_2O = 79.5/19.5/1.0$, 1 atm, $\epsilon_{r1} = \epsilon_{r2} = 2.2$
- N. Babaeva, PSST 20, 035017 (2011)

STREAMER DYNAMICS: VOLTAGE



- As the streamer crosses the gap, voltage is shorted by the plasma column, transferring much of the applied potential to the streamer head.
- Upon intersection with the surface, much of this voltage is transferred first to the sheath and then to the capacitance of the dielectric.
- Positive corona, $N_2/O_2/H_2O = 79.5/19.5/1.0$, 1 atm, $\epsilon_{r1} = \epsilon_{r2} = 2.2$

STREAMER DYNAMICS: ELECTRIC FIELD



- In the streamer head, E-fields are 200 kV/cm when the streamer is far from surfaces.
- Upon intersection of the streamer with the surface, the electric fields in the resulting sheath can exceed 500-800 kV/cm.
- $800 \text{ kV/cm} \times \lambda (0.5 \mu\text{m}) = 40 \text{ eV}$

- Spreading of streamer on surface produces a wave of sheath electric field away from point of intersection with E-fields of 300-500 kV/cm.
- Fields are radially outward at first, then oriented towards surface.
- Positive corona, $\text{N}_2/\text{O}_2/\text{H}_2\text{O} = 79.5/19.5/1.0$, 1 atm, $\epsilon_{r1} = \epsilon_{r2} = 2.2$.

Animation Slide

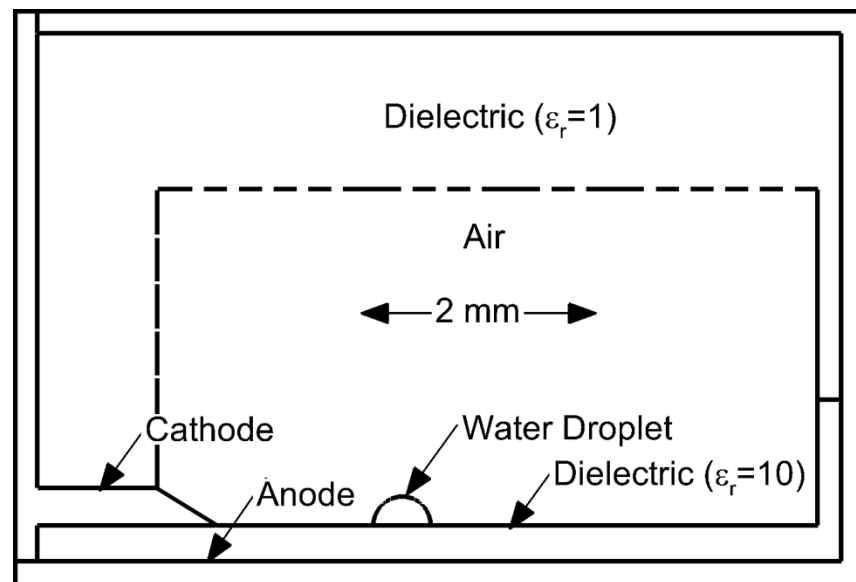
MIN  MAX

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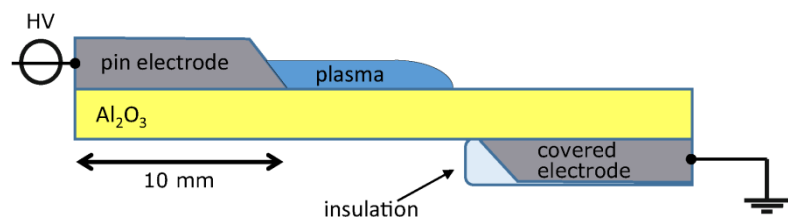
DELIBERATE SURFACE DISCHARGES OVER COMPLEX SURFACES

APP SURFACE DISCHARGES

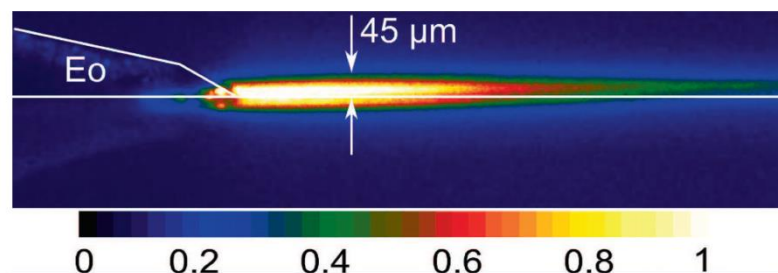
- Modeling of humid air surface discharges over nonplanar surfaces and droplets on dielectric.
- Deliberate launching of SIWs in a coaxial geometry.
- Deliberately over-voltaged to maintain SIW over long distances.
- Droplet study motivated by virus containing aerosol contamination.



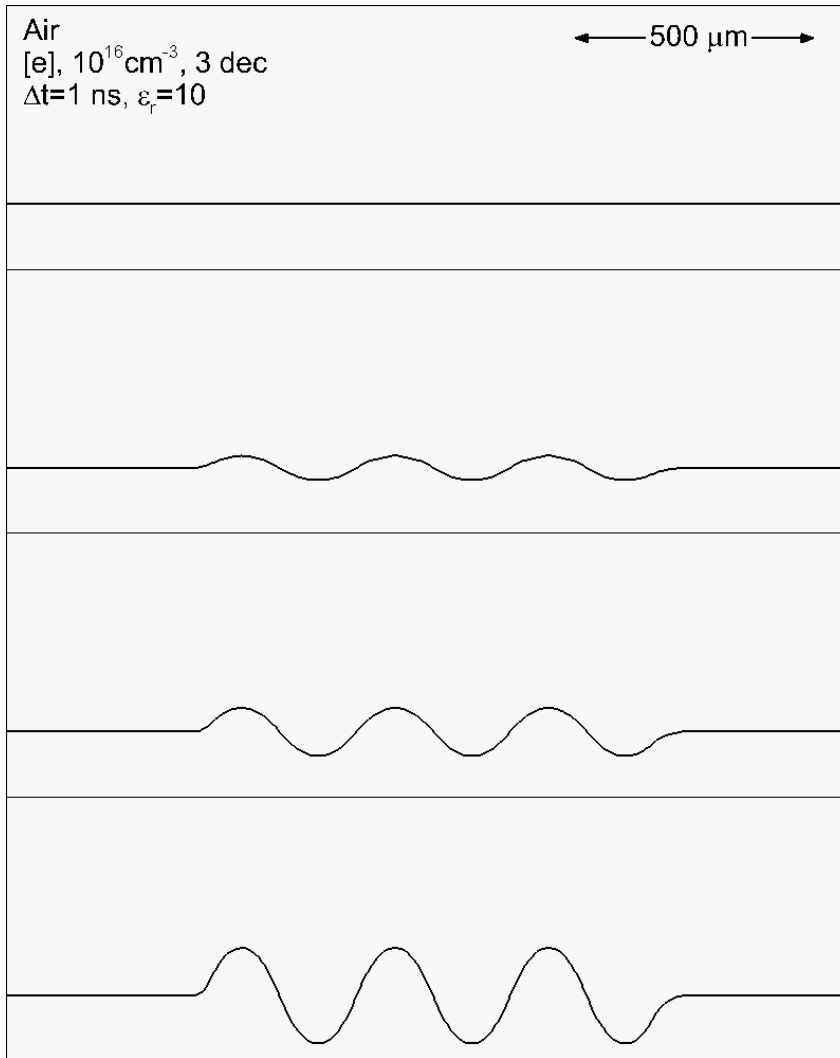
• Surface Discharge Geometry



- Motivating experiments (Hoft and Brandenburg, EPJD 74, 110 (2020)).

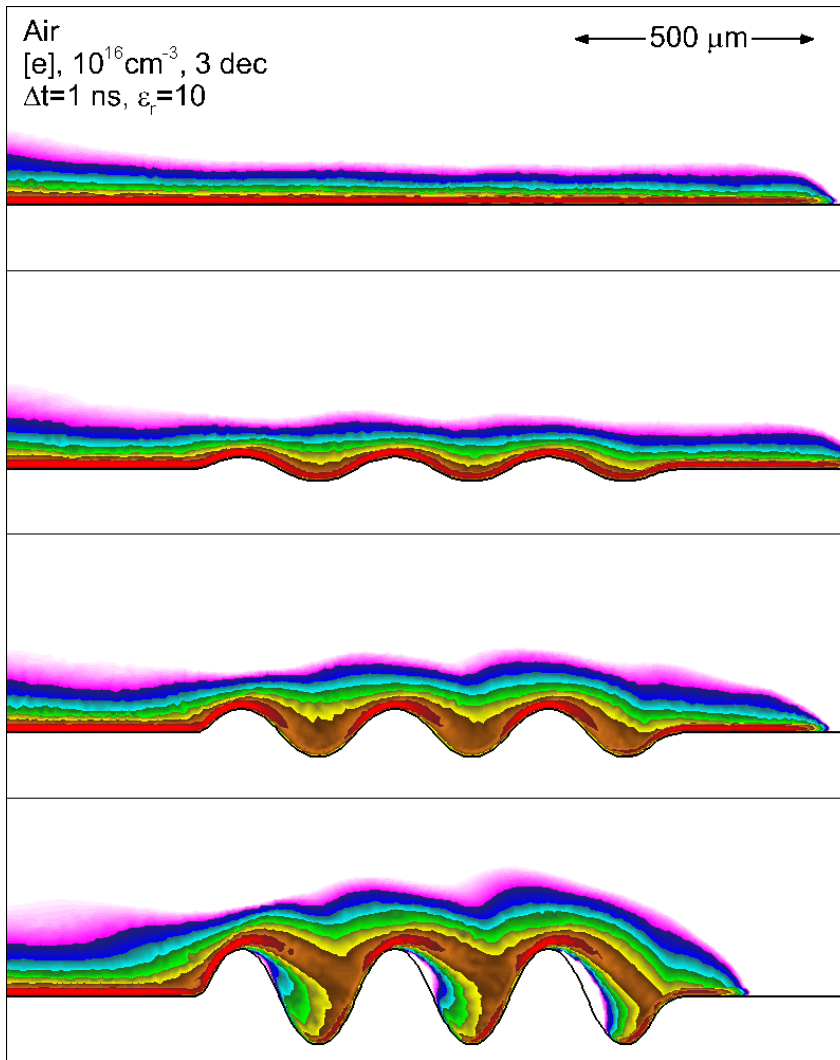


HUMID AIR WAVY SURFACE DISCHARGE: [e]



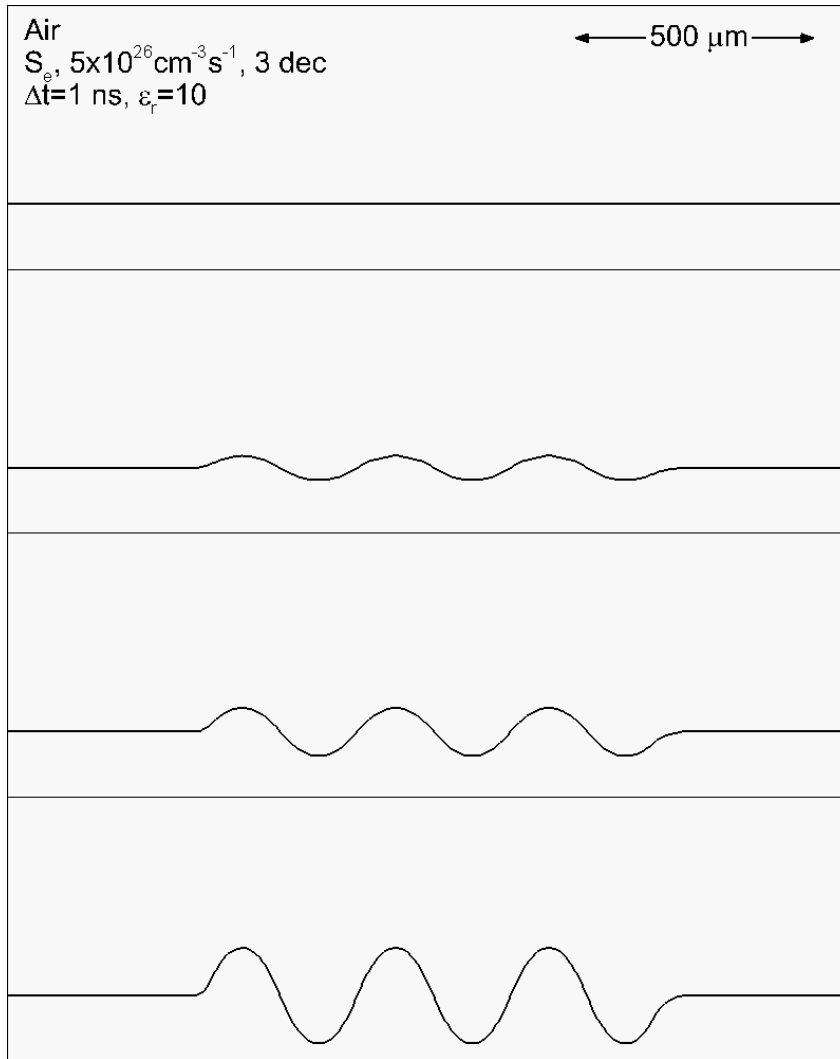
- Humid air, -20 kV (step), $\epsilon_r = 10$
- Electron density (10^{16} cm^{-3} , 3-dec)
- Surface ionization wave thickness 100 μm on flat surface.
- SIW adheres to surface for wave height $< 300 \mu\text{m}$.
- Separation of SIW from surface for deeper ridges.
- Ridges produce electric field enhancement.
- Ridges block photoelectron emission and photoionization ahead of SIW.

HUMID AIR WAVY SURFACE DISCHARGE: [e]



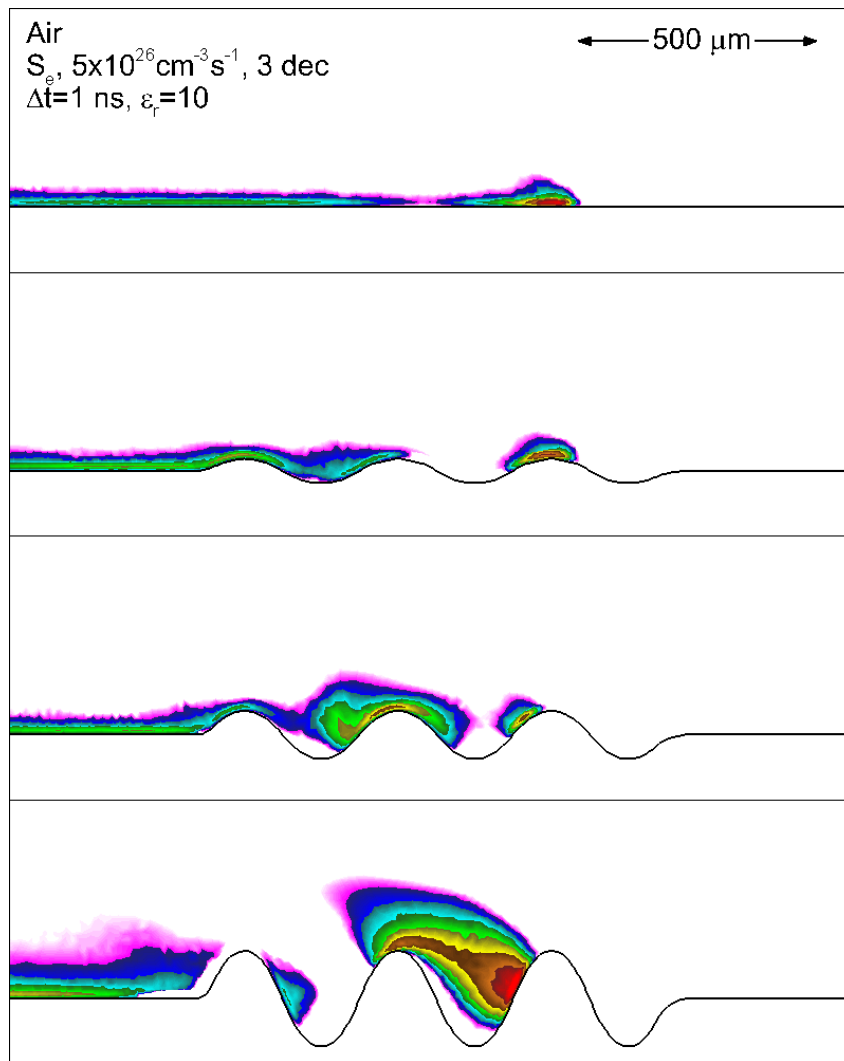
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WAVY SURFACE IONIZATION SOURCE



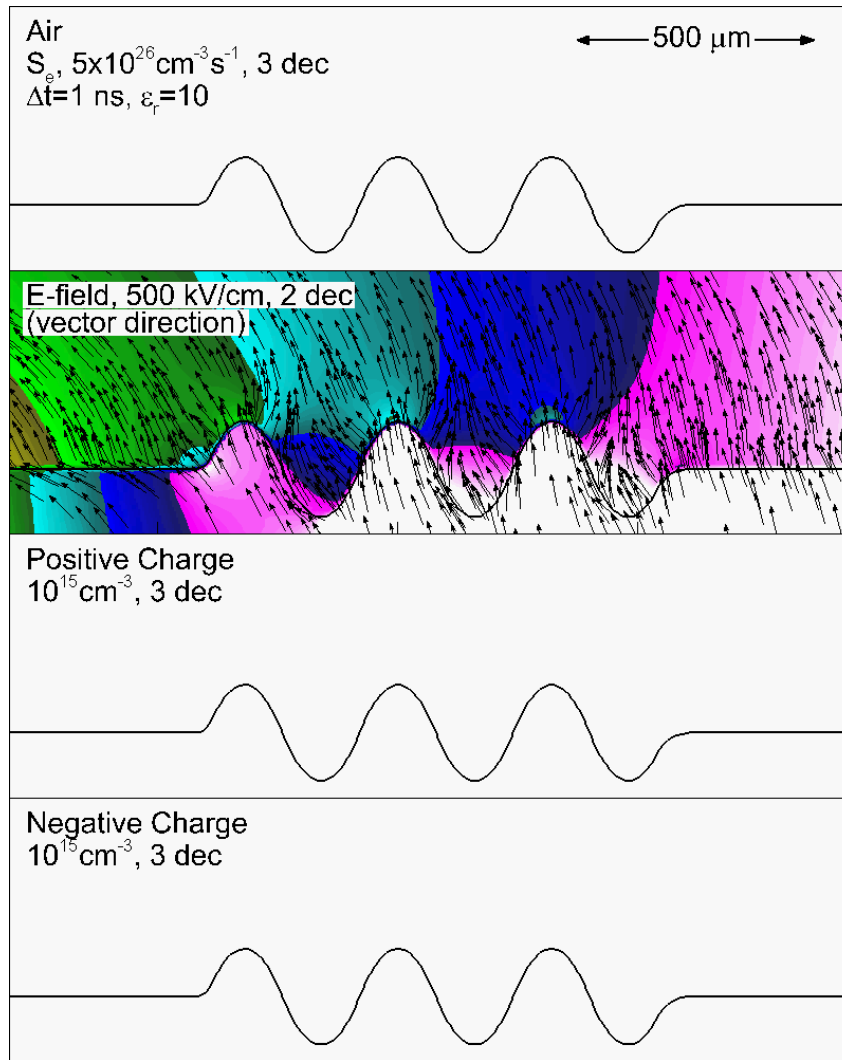
- Humid air, -20 kV (step), $\epsilon_r = 10$
- Ionization Rate ($5 \times 10^{26} \text{ cm}^{-3} \text{ s}^{-1}$ 3 dec)
- Even shallow waves produce backwards traveling surface ionization waves.
- Process is similar to conventional dielectric barrier discharge where surface charging launches reverse filament.
- Apex of wave with electric field enhancement intensifies ionization, relaunching streamer.
- Valleys have low electric field and slow (or quench) SIW.

WAVY SURFACE IONIZATION SOURCE



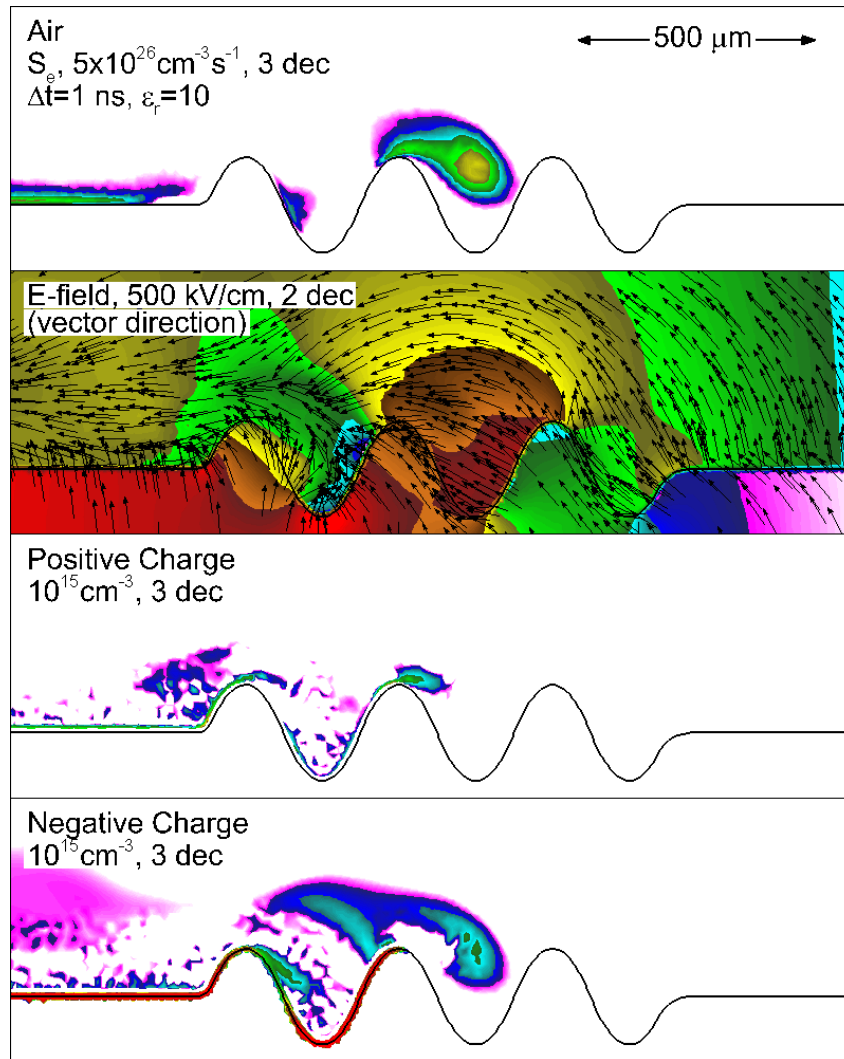
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WAVY SURFACE CHARGING



- Humid air, -20 kV (step), $\epsilon_r = 10$
- Surface charging is in large part responsible for SIW hopping and restrike.
- Negative SIW charges surface ahead and beneath negative, which largely shields voltage.
- Positive sheath forms directly above surface.
- Shielded voltage facilitates field enhancement at next peak, and provides needed polarity for positive restrike.

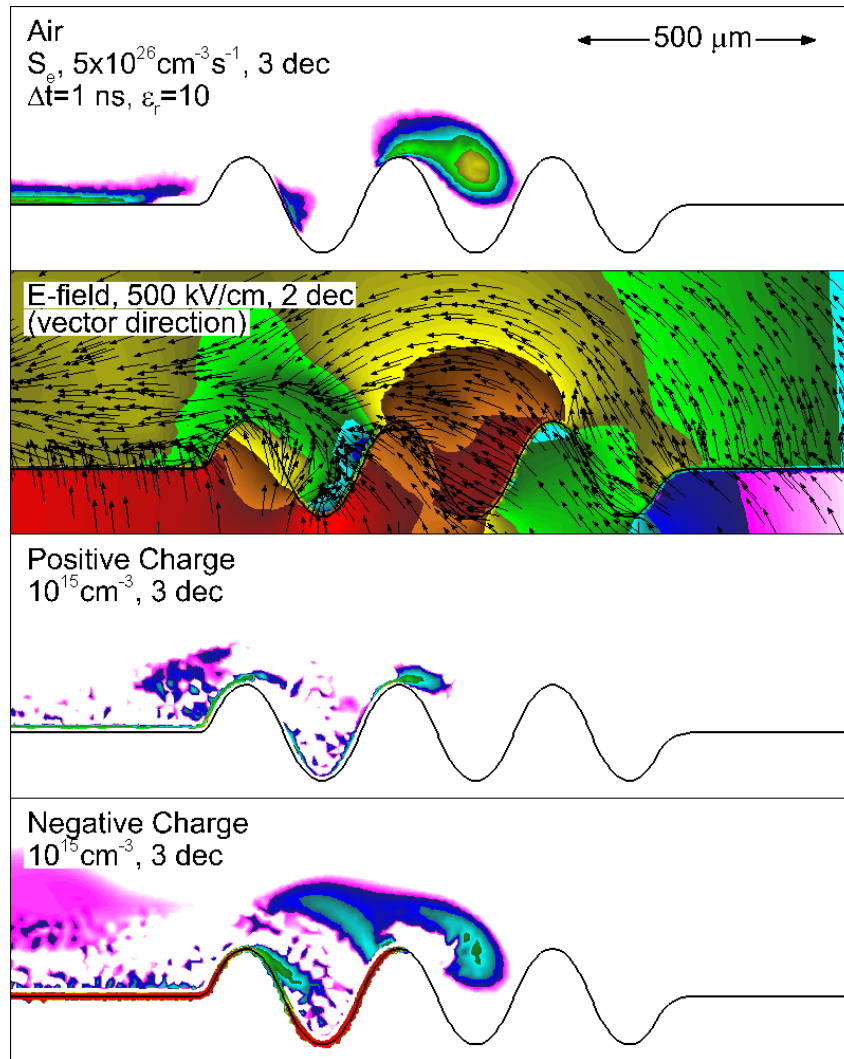
WAVY SURFACE CHARGING



Min  Max

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DIELECTRIC CONSTANT OF WAVY SURFACE

[e] $2 \times 10^{16} \text{ cm}^{-3}$ (2.5 dec)
0 - 1.6 ns

[e]

$\epsilon_r = 4$

250 μm

$\epsilon_r = 20$

$S_e 7 \times 10^{26} \text{ cm}^{-3} \text{ s}^{-1}$ (3 dec)
0 - 1.6 ns

S_e

$\epsilon_r = 4$

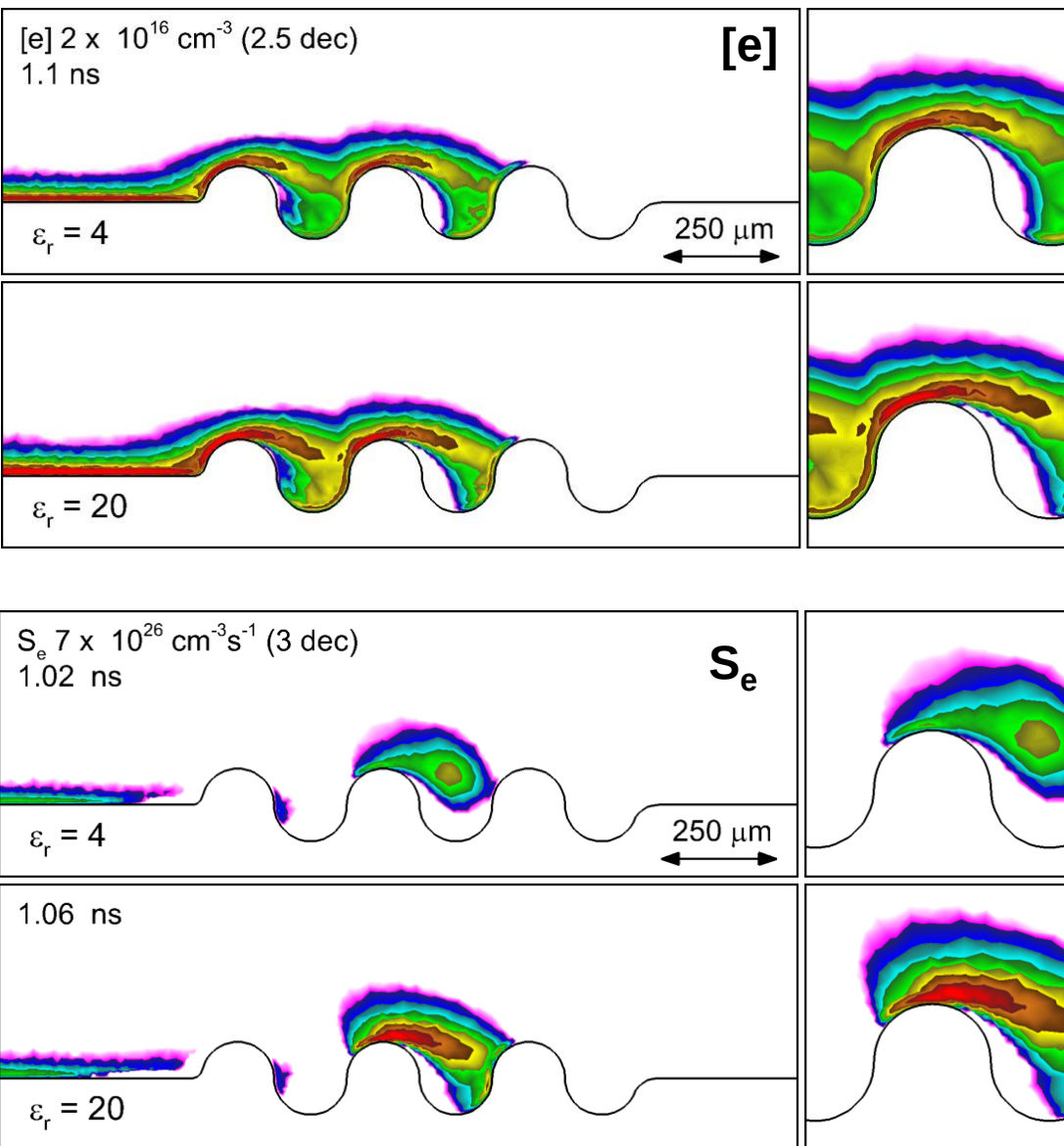
250 μm

$\epsilon_r = 20$

- Humid air, -20 kV (step)
- Dielectric constant ($\epsilon_r = 4, 20$) should affect propagation of SIW by:
- Capacitance is proportional to ϵ_r - larger capacitance, larger RC constant, slower propagation.
- Larger ϵ_r , more electric field enhancement at apex, producing larger ionization rates.
- With large over-voltage, effects are detectable, but not larger.

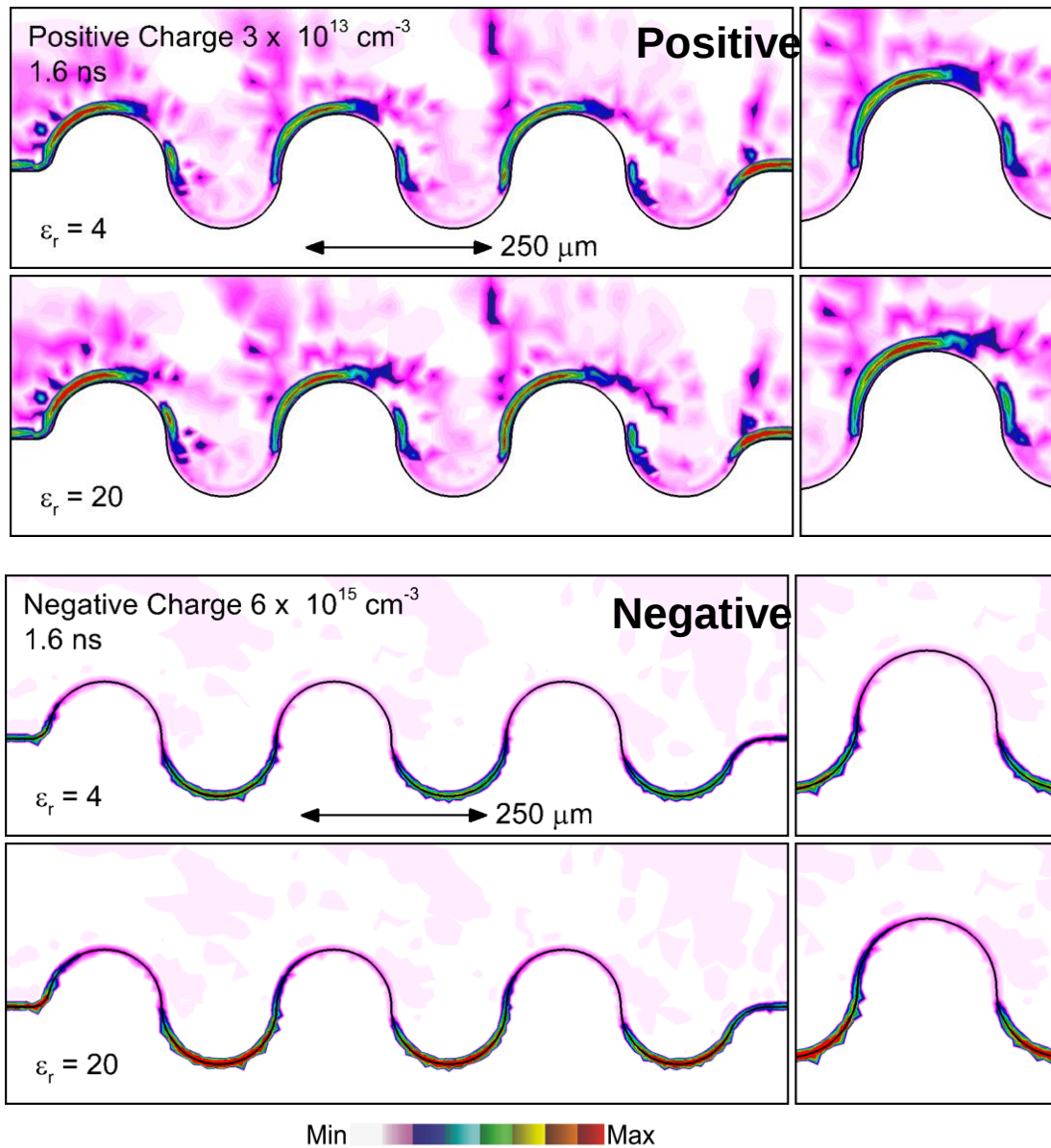
Animation Slide

DIELECTRIC CONSTANT OF WAVY SURFACE



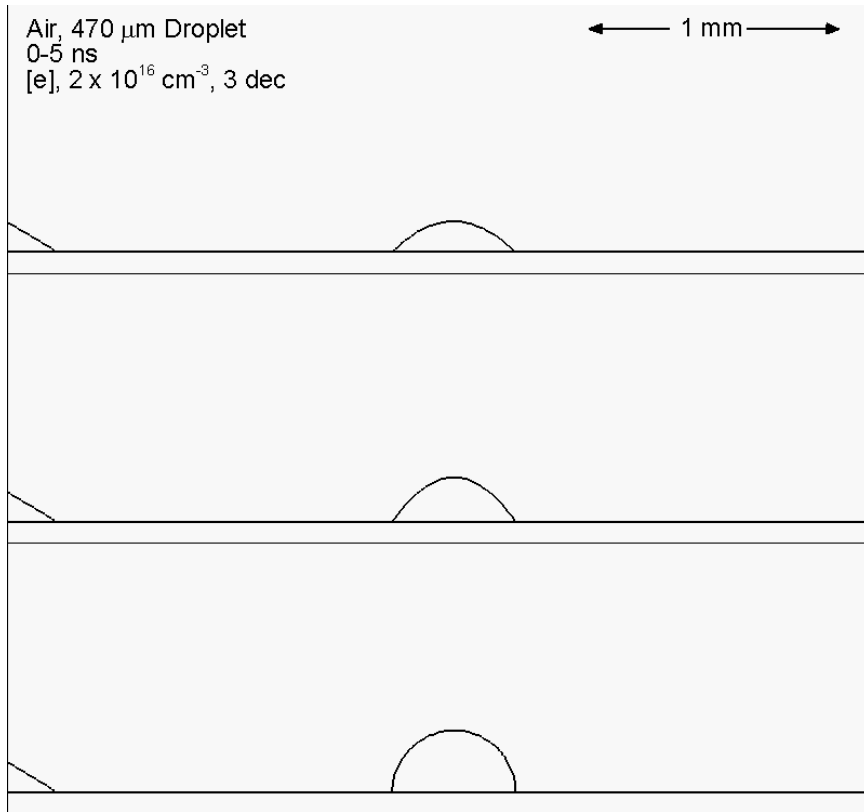
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ϵ_r OF WAVY SURFACE: CHARGING



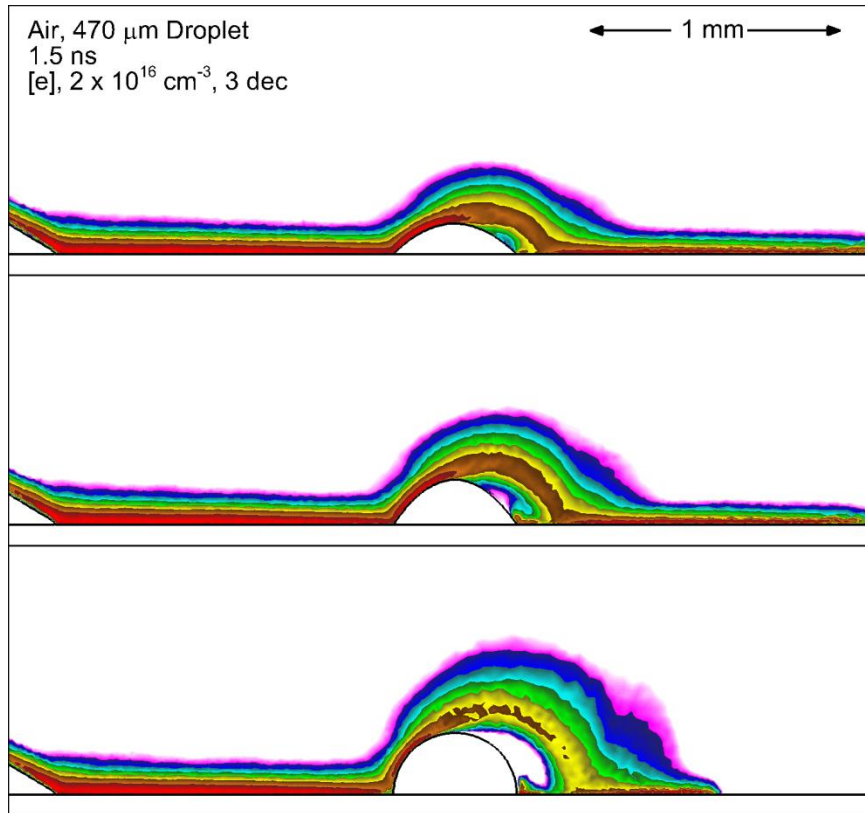
- Humid air, -20 kV, $\epsilon_r = 4, 20$
- Surface capacitance determines rate of charging, but not final state.
- Once surface is charged and potential established, gas phase does not see difference in ϵ_r (E-field enhancement aside).
- Sheath (and restrike remnants) not terribly sensitive to ϵ_r .
- Negative surface charging scales with ϵ_r to charge capacitance.
- Much less sensitive than packed bed reactor.

WATER DROPLET ON DIELECTRIC: [e]



- Humid air, -24 kV (2.5 ns), $\epsilon_r=10$, electron density (10^{16} cm^{-3} , 3 dec)
- Motivated by plasma deactivation of virus containing aspirated aerosols landing on surfaces.
- 450 μm diameter water droplet with varying contact angle reflecting different materials.
- Droplets behave similar to wavy surface – plasma detaches when contact angle is small.
- Significant electric field enhancement at plasma-water-dielectric triple point.

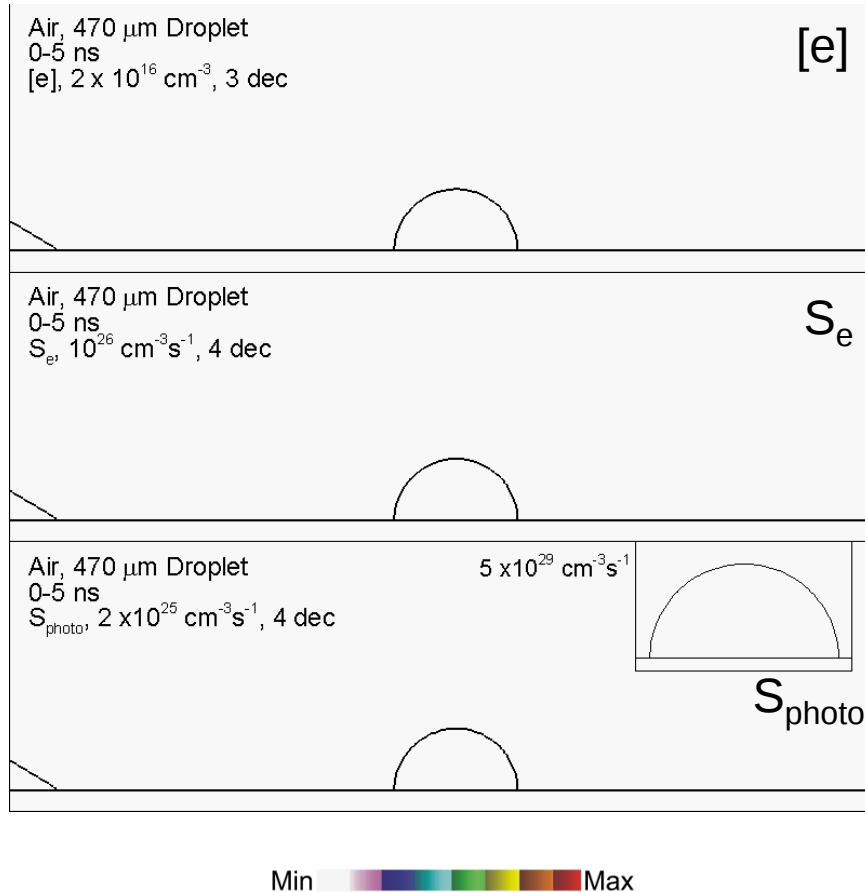
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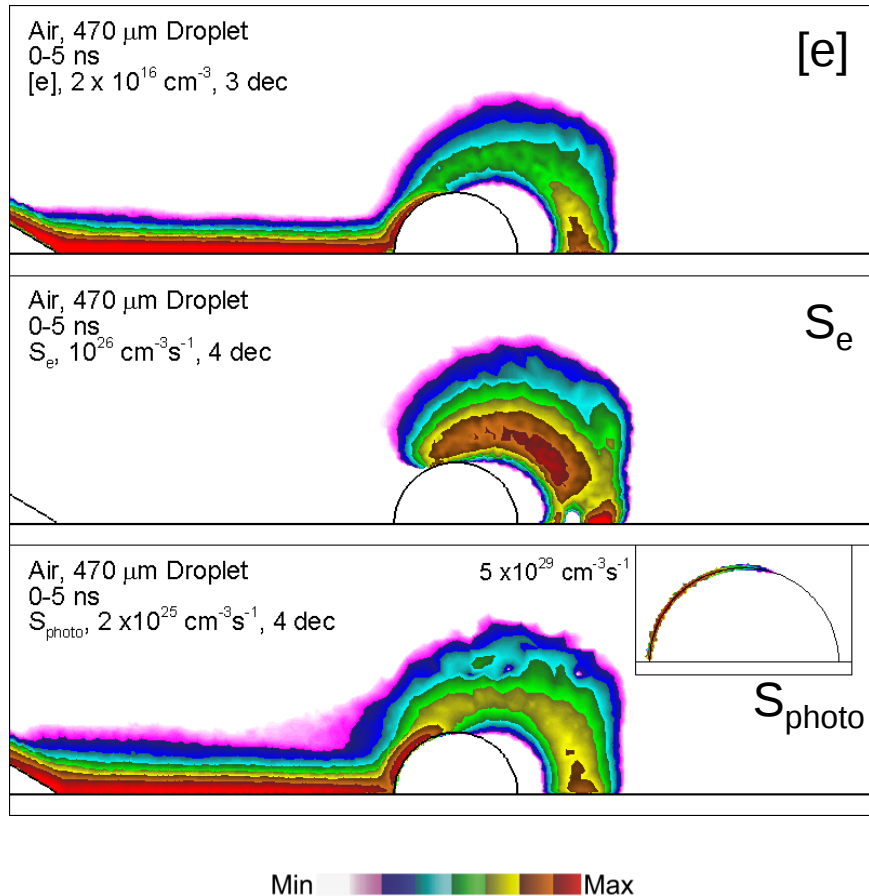
WATER DROPLET ON DIELECTRIC: IONIZATION



- Humid air, -24 kV (2.5 ns), $\epsilon_r=10$
- On time scale of ns pulse, water acts as a dielectric whose surface charges.
- Electrons immediately solvate but do not convect far from surface.
- Electric field enhancement at leading corner attracts SIW.
- Surface charging launches ionization wave over droplet.
- Photoionization at surface of droplet under SIW is 10^3 - 10^4 times more intense than in gas phase.
- Majority of immediate activation is VUV dominated.

Animation Slide

WATER DROPLET ON DIELECTRIC: IONIZATION



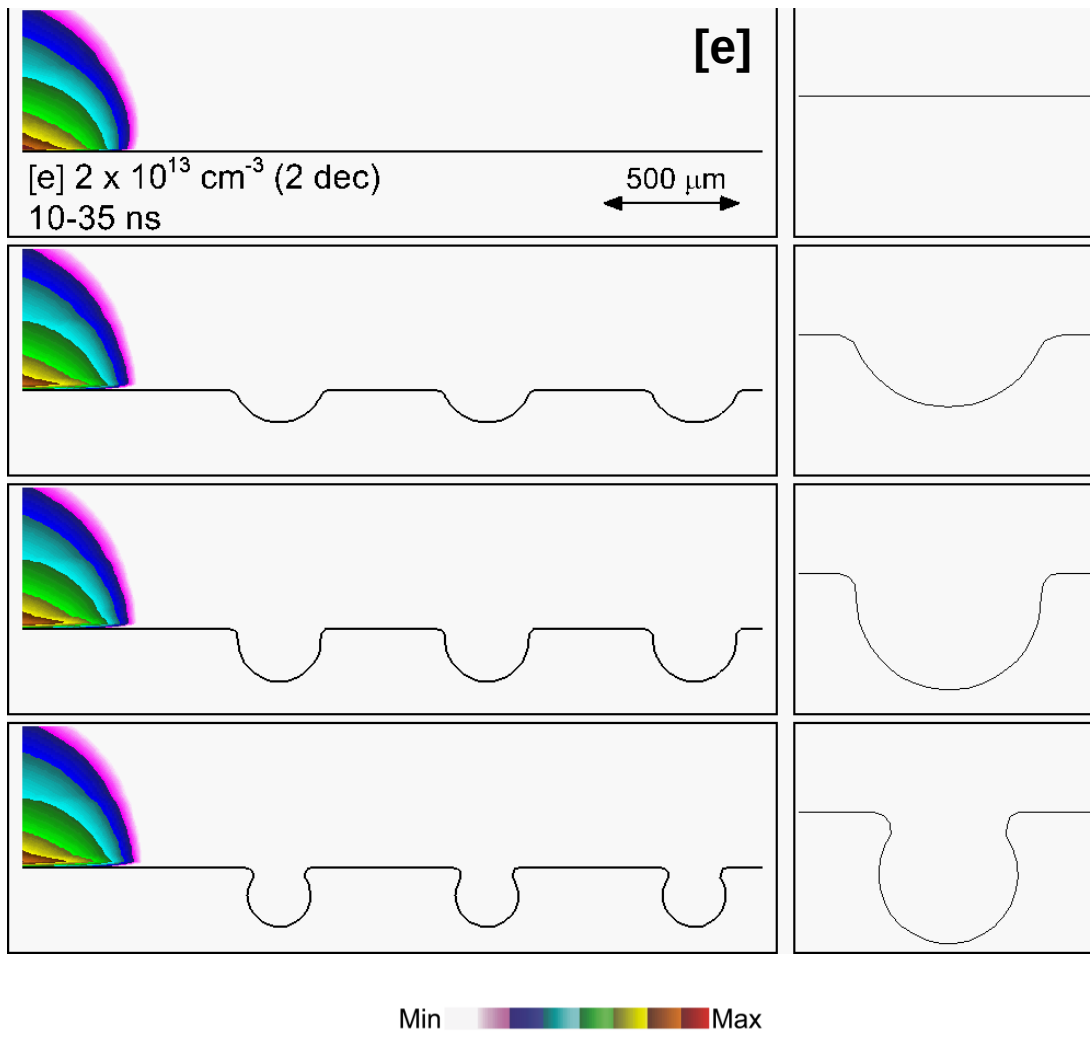
- Humid air, -24 kV (2.5 ns), $\epsilon_r=10$
- On time scale of ns pulse, water acts as a dielectric whose surface charges.
- Electrons immediately solvate but do not convect far from surface.
- Electric field enhancement at leading corner attracts SIW.
- Surface charging launches ionization wave over droplet.
- Photoionization at surface of droplet under SIW is 10^3 - 10^4 times more intense than in gas phase.
- Majority of immediate activation is VUV dominated.

APPs ACROSS POROUS MATERIALS



- Propagation of APPs across (or onto) porous materials is sensitive to the “cut” through the materials.
- Pores can be:
 - Subsurface
 - Barely opened
 - “Sliced open”.
- Pores are typically interconnected, enabling pore-to-pore propagation.
- Do “inverted” waves behave the same?
- How do plasmas get into pores?

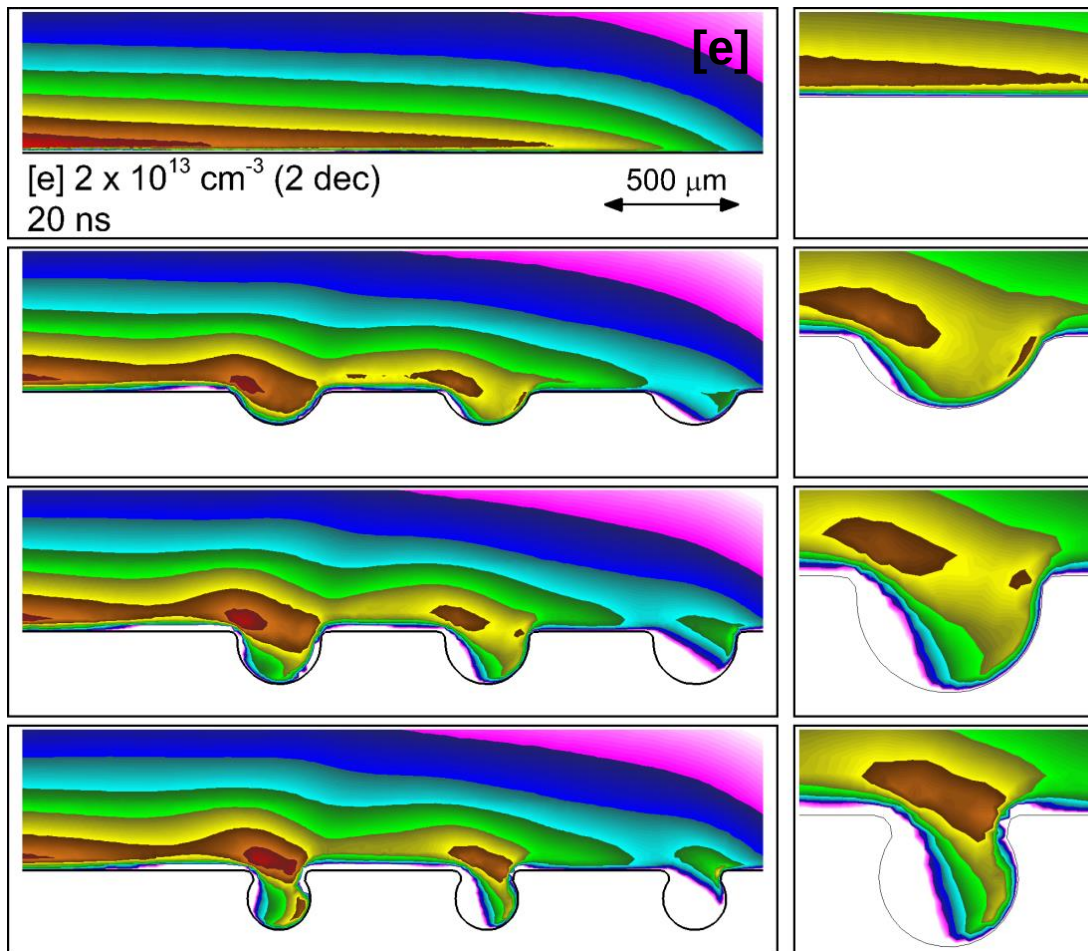
“CUT” THROUGH POROUS MATERIAL: [e]



- Humid air, -25 kV (step), $\epsilon_r=4$
- Electron density
- APPs react to cut pores similarly to wavy surfaces.
- Interaction is here more dominated by electric field enhancement at edges of cut pores.
- Overshoot of pore due to E-field enhancement on opposite edge leads to unexposed surfaces.
- Charging dominates in production of restrike SIWs.

Animation Slide

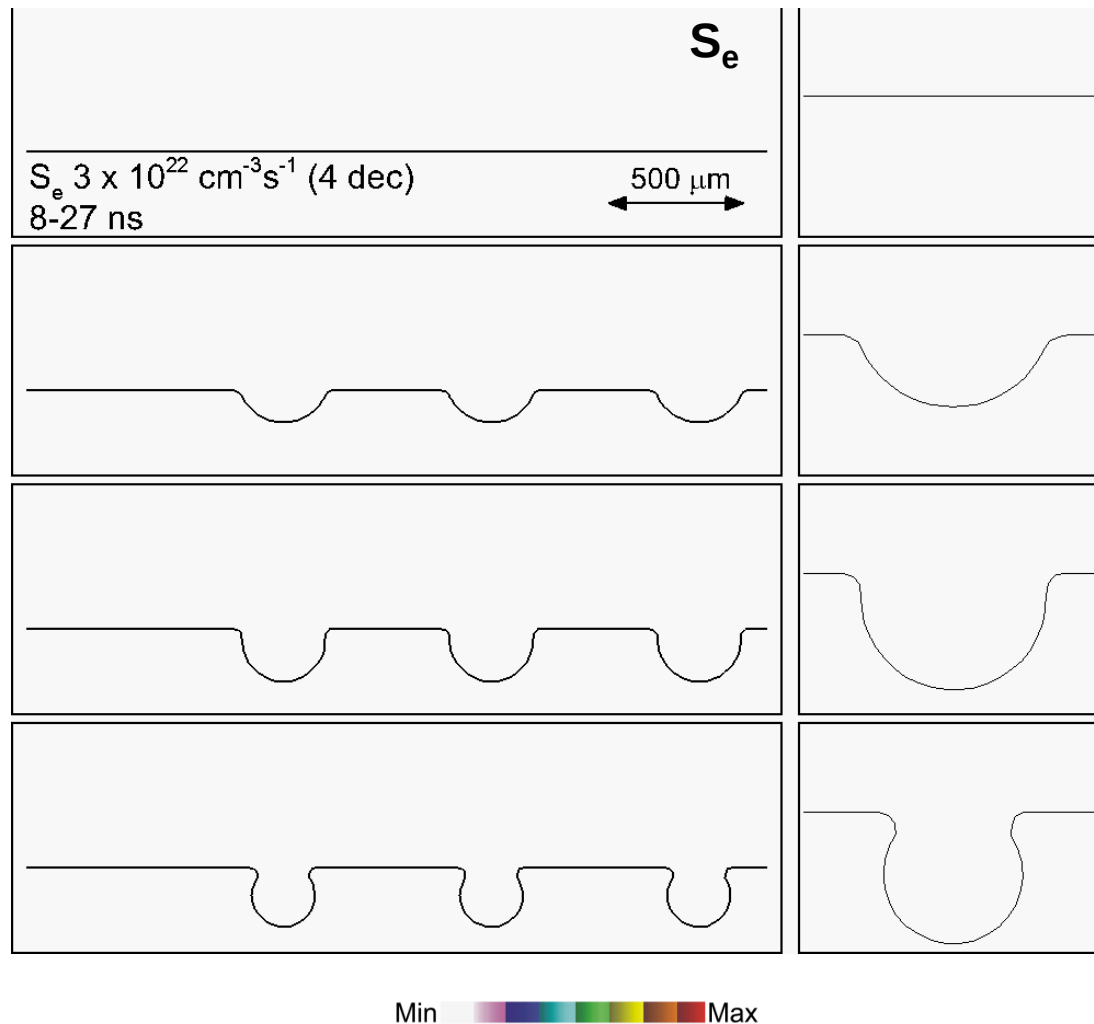
“CUT” THROUGH POROUS MATERIAL: [e]



Min  Max

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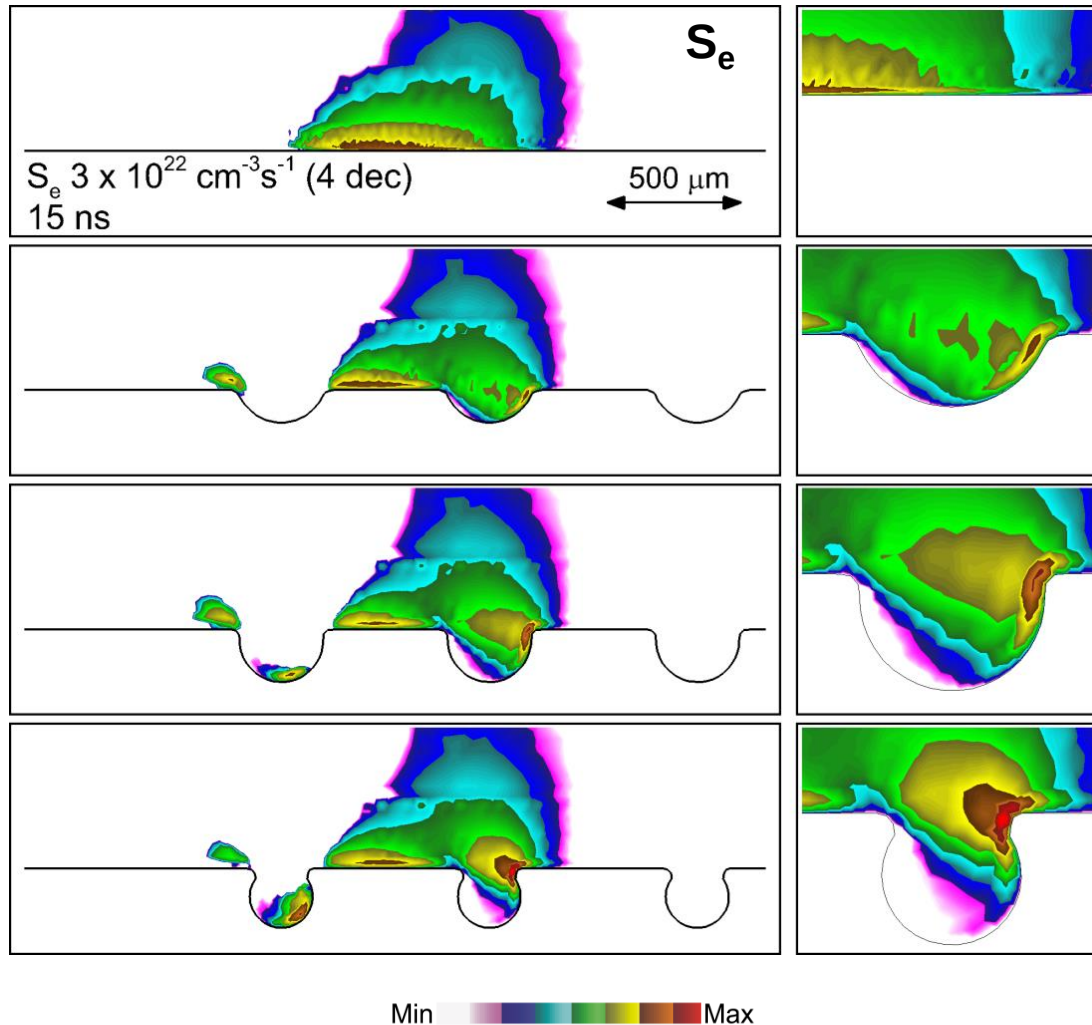
“CUT” THROUGH POROUS MATERIAL: S_e



- Humid air, -25 kV (step), $\epsilon_r=4$
- Ionization source.
- Sharp edges produced electric field enhancement and have larger capacitance.
- Intense ionization at “more than half pore” edges
- Charged induced restrike SIW provides full coverage of shallow pores.
- Restrike SIW in greater than half-pores are not able to fully cover the inner surface.

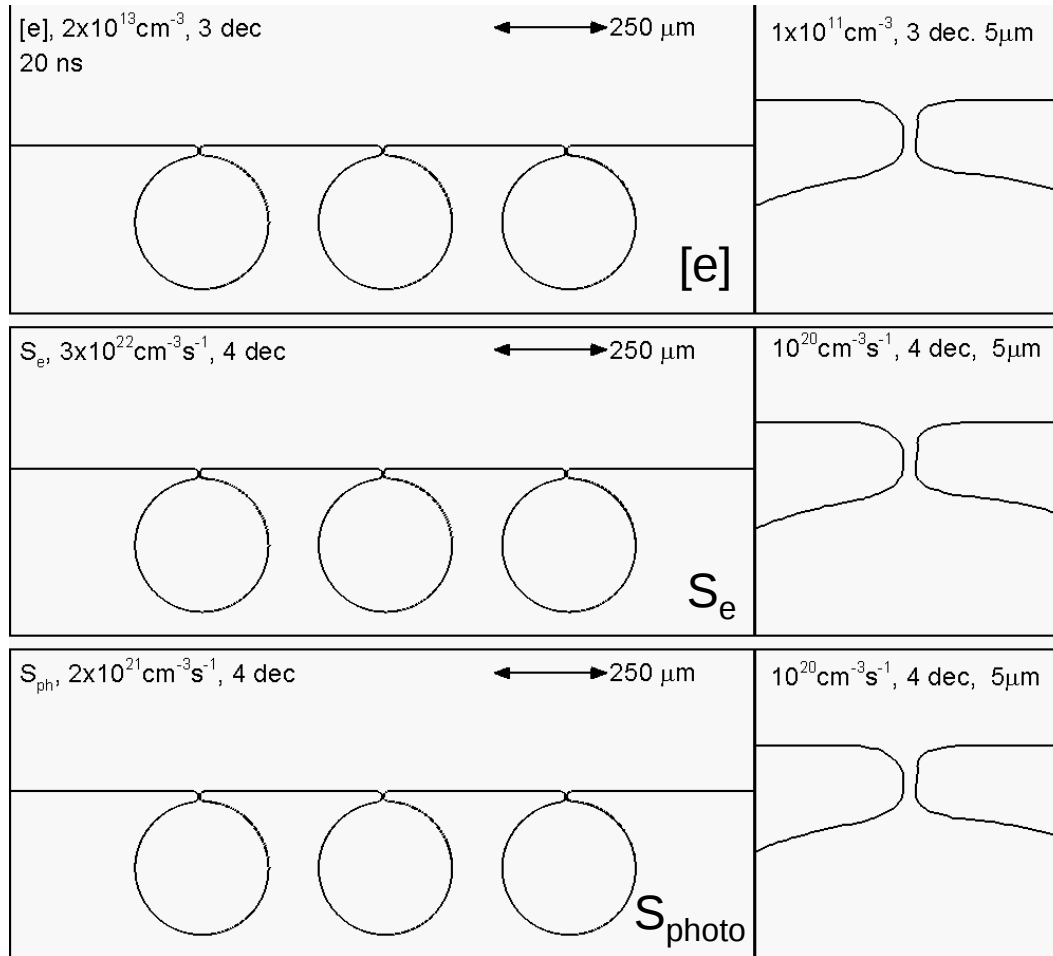
Animation Slide

“CUT” THROUGH POROUS MATERIAL: S_e



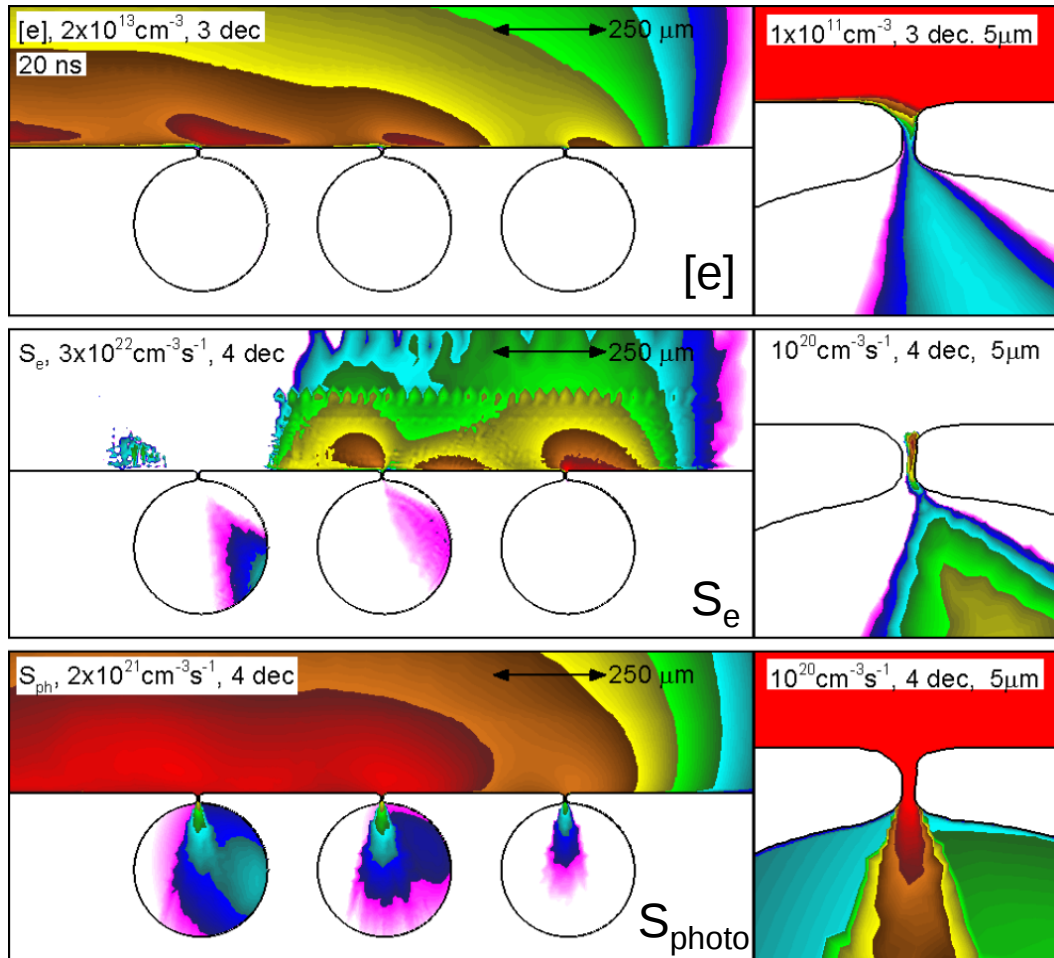
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SIW OVER SUBSURFACE (OPEN) PORES



- Humid air, -25 kV (step), $\epsilon_r=4$
- Flat surface with subsurface pores (300 μm)
- 5 μm opening to plasma to small to enable plasma to naturally flow ($\lambda_D = 3\text{-}4 \text{ μm}$). Leakage is non-ambipolar.
- VUV illumination through pore opening seeds plasma.
- Plasma “sees” subsurface pore through lower capacitance of underlying dielectric.
- Electric field enhancement at opening.

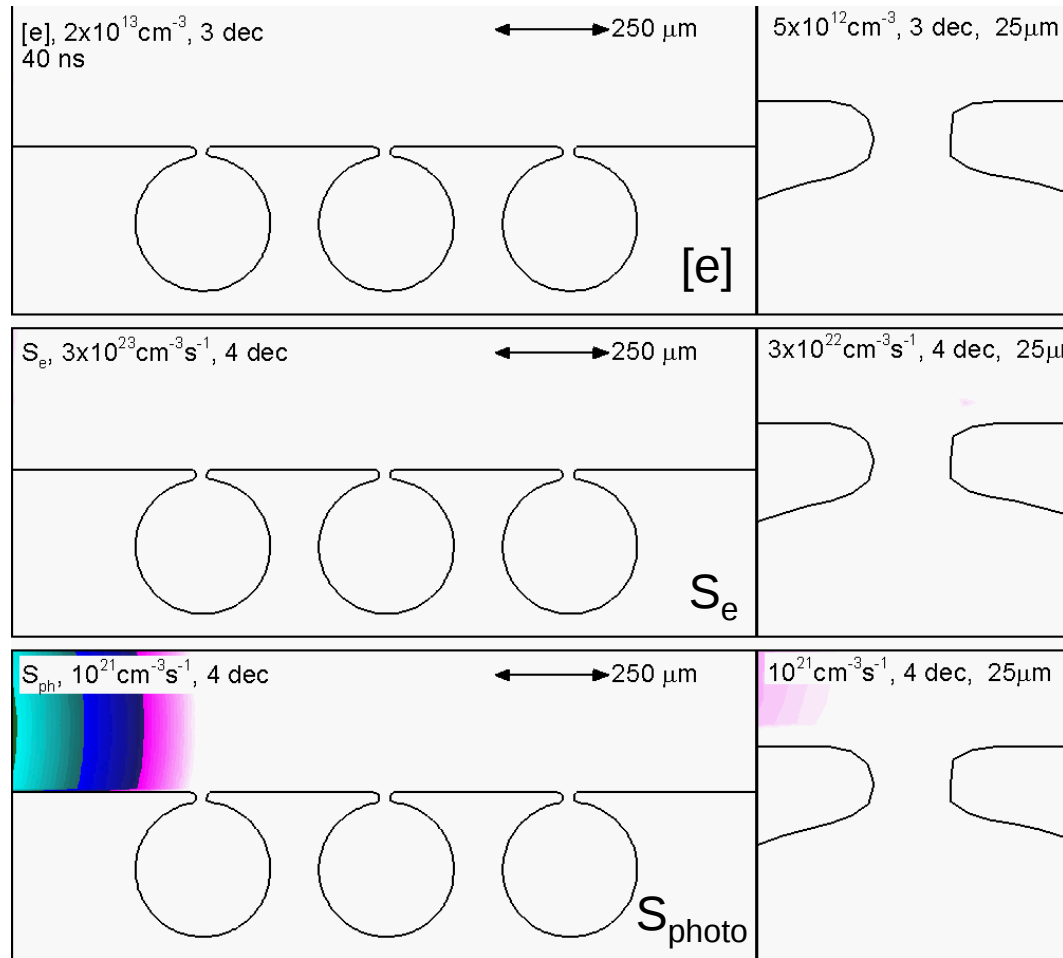
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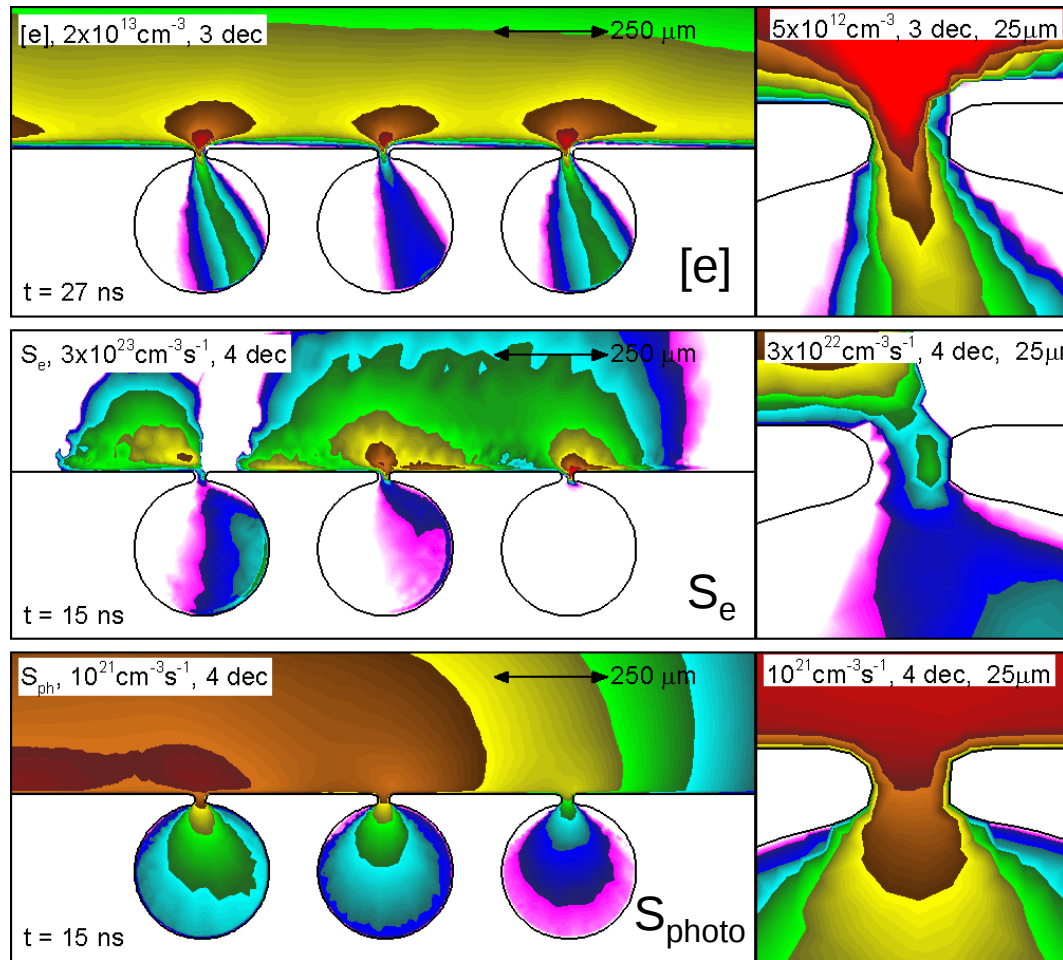
Min  Max

SIW OVER SUBSURFACE PORES: LARGE OPENING



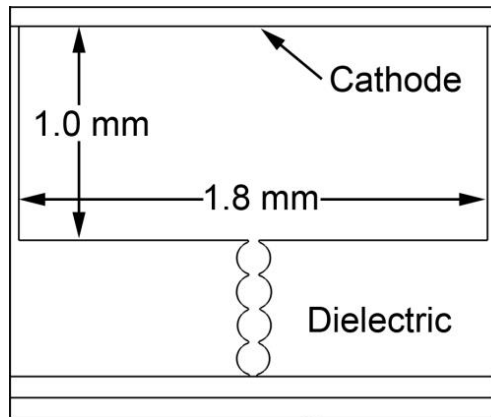
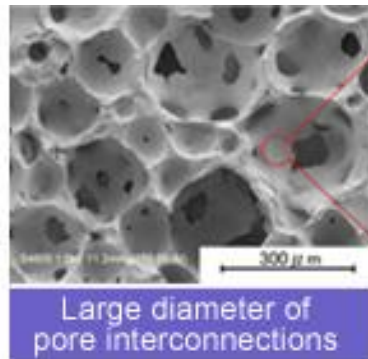
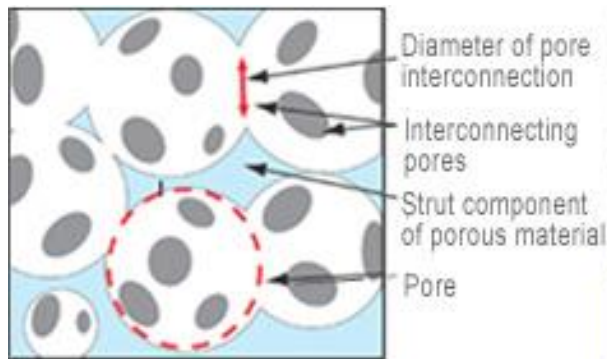
- Humid air, -25 kV (step), $\epsilon_r=4$
- Flat surface with subsurface pores (300 μm)
- 25 μm opening to plasma to exceeds Debye length and allows flow of plasma into pore.
- VUV illumination seeds plasma, though electron impact dominates ionization.
- Plasma “sees” subsurface pore through lower capacitance of underlying dielectric.
- Pores see each other through charging of surfaces.

SIW OVER SUBSURFACE PORES: LARGE OPENING

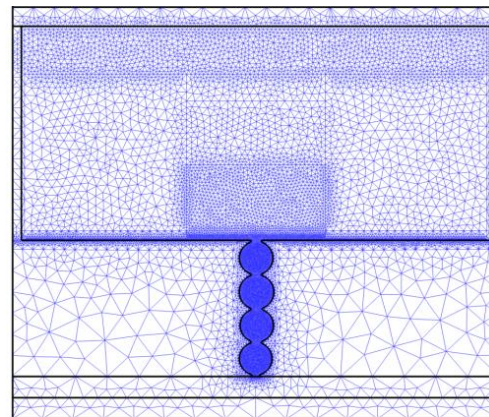


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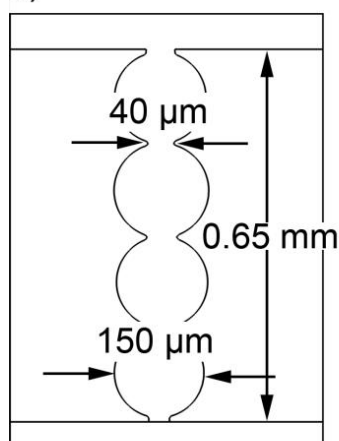
CONNECTED CHAINS OF PORES



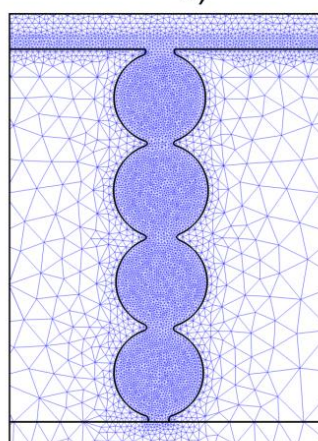
a) Ground



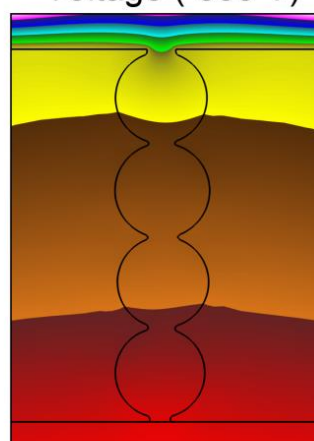
b) Voltage (-800 V)



c)



d)



e)

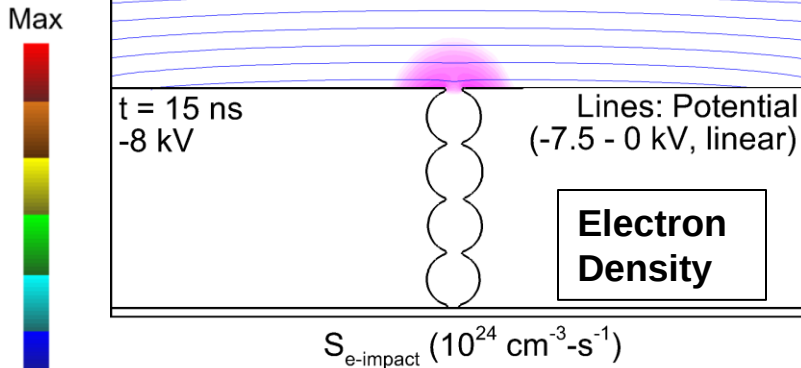
Min  Max

INTERCONNECTED PORES

- Treatment of internal surface of interconnected pores of interest to tissue scaffolding.
- Examine single chain of pores.
- Dielectric barrier discharge (DBD) in humid air ($N_2/O_2/H_2O = 78/21/1$).
- 1.0 mm gas gap, 0.65 mm dielectric layer ($\epsilon_r = 61$, wet bone).
- -8 kV, 15 ns, top electrode.

MICRODISCHARGE FORMATION

n_e (10^{15} cm^{-3} , 4-dec)



- Electron density, equipotential lines, and electron impact ionization source.
- Townsend avalanche begins near cathode, electrons impact and charge scaffold.
- Restrike propagates towards cathode.
- Surface ionization waves develop on scaffold.
- Conductive plasma shorts potential in pores, producing high electric field.
- Discharges take place inside pores.

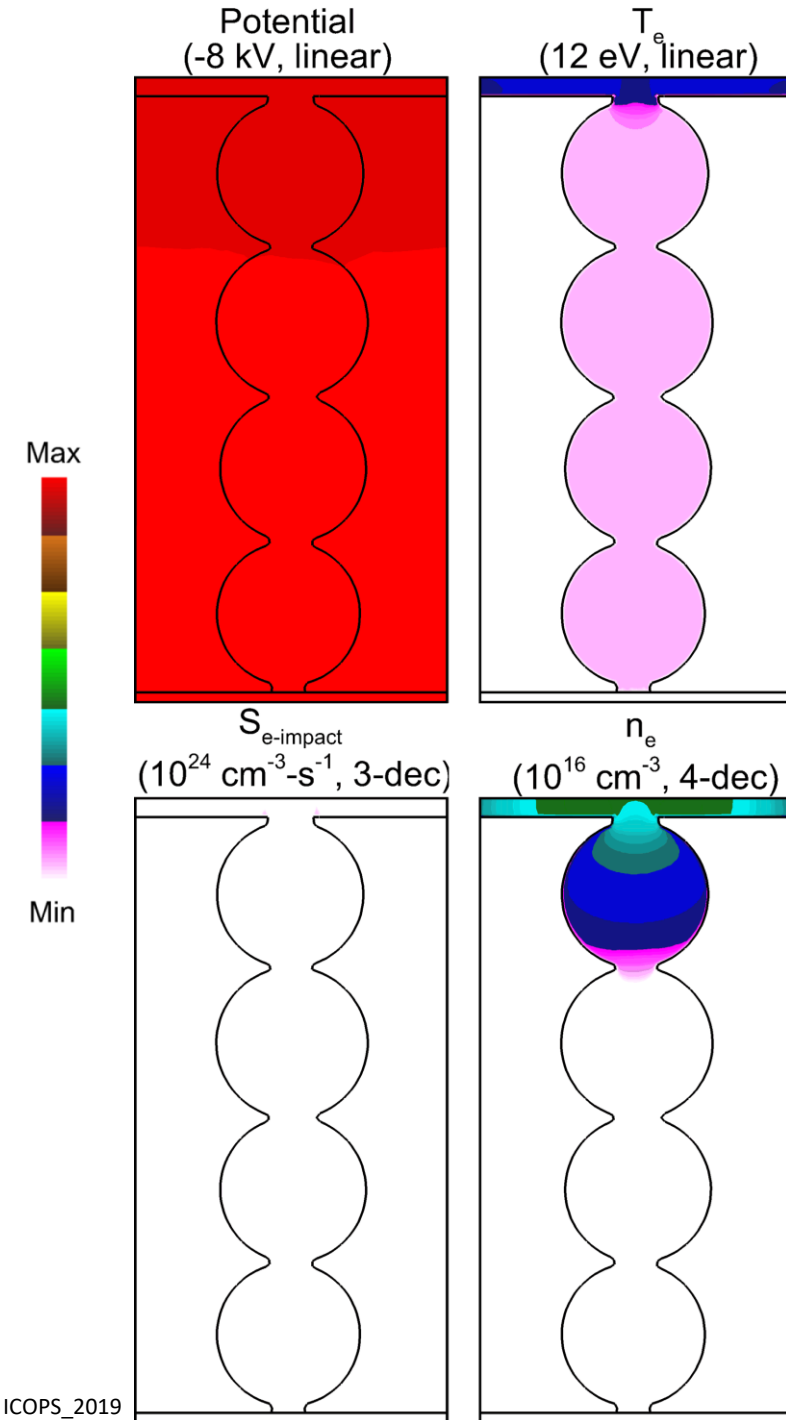
Animation Slide

PROPAGATION INTO PORES

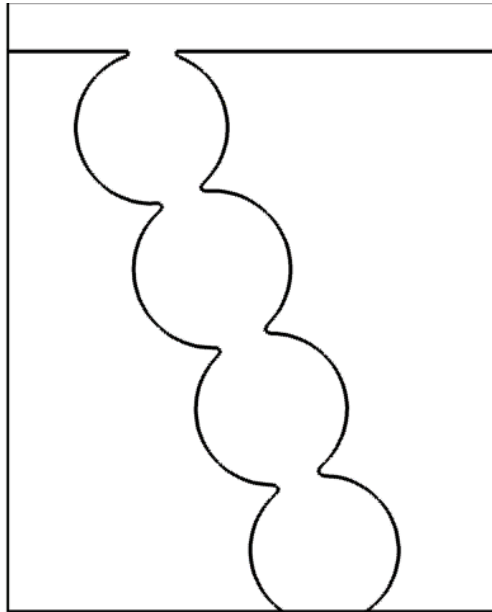
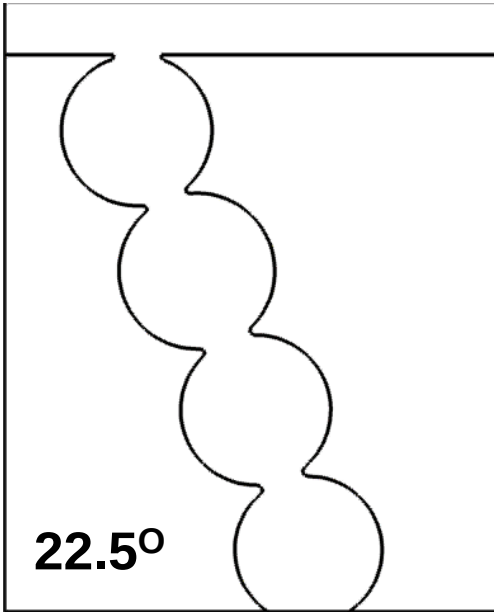
- Initially, electrons enter pores following electric field.
- Potential drop over small distance leads to high electron temperature (peak at $T_e \approx 11.7$ eV).
- Microdischarges form between pore connectors.
- Charging of connectors leads to formation of surface ionization waves (SIWs).

Animation Slide

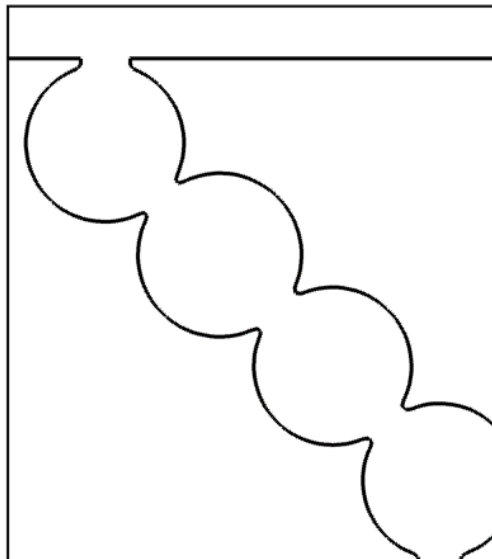
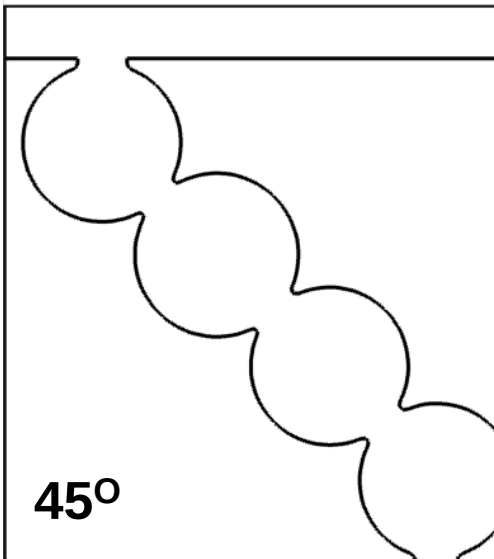
University of Michigan
Institute for Plasma Science & Engr.



n_e (10^{17} cm^{-3} 4 dec) S_e ($10^{27} \text{ cm}^{-3}\text{s}^{-14}$ dec)



n_e (10^{17} cm^{-3} 4 dec) S_e ($10^{27} \text{ cm}^{-3}\text{s}^{-14}$ dec)

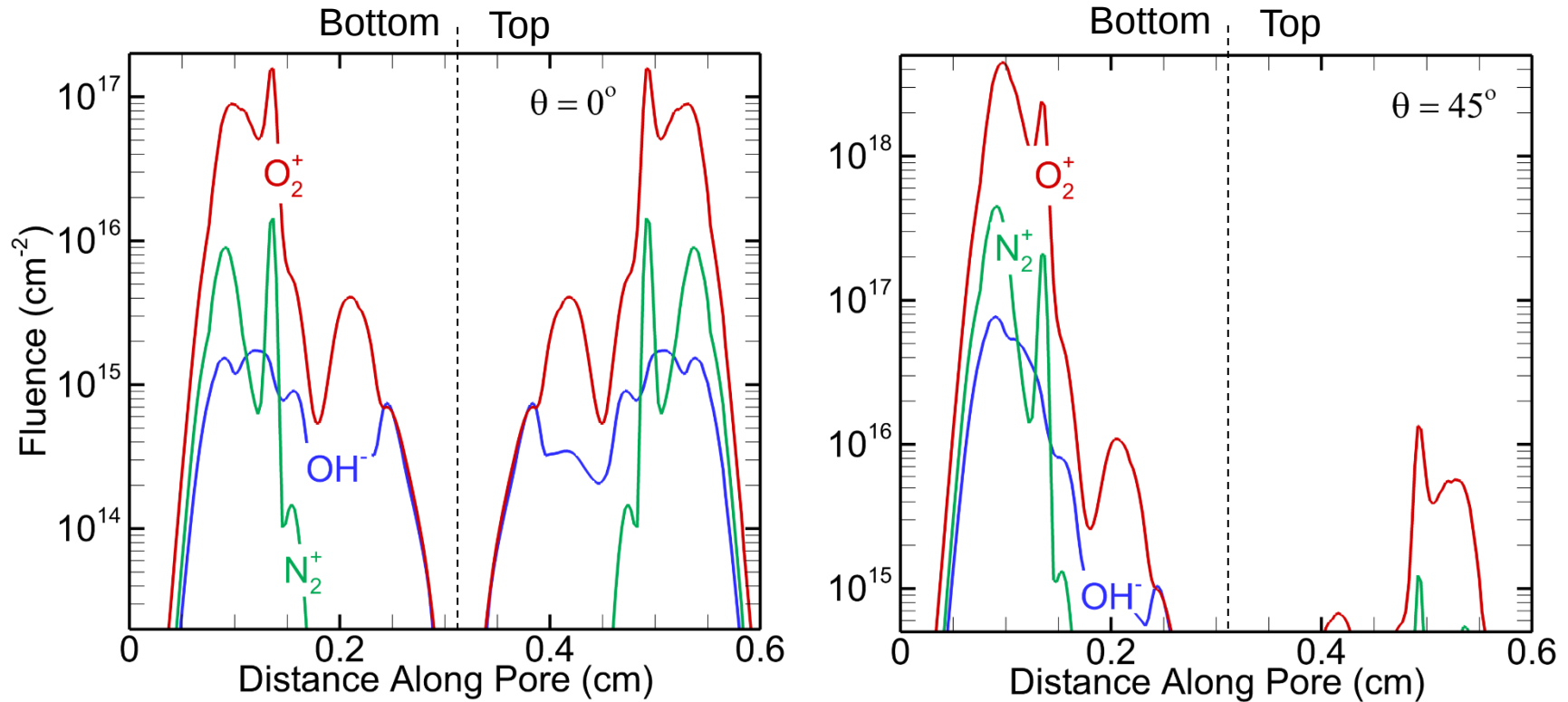


PORE ORIENTATION

- Original orientation of electric field is imprinted on IW propagation.
- Oblique pores are increasingly shadowed have IW on bottom surface.
- Restrike SIW fails to cover all surfaces.
- 1 atm, humid air, -8 kV , 5 ns

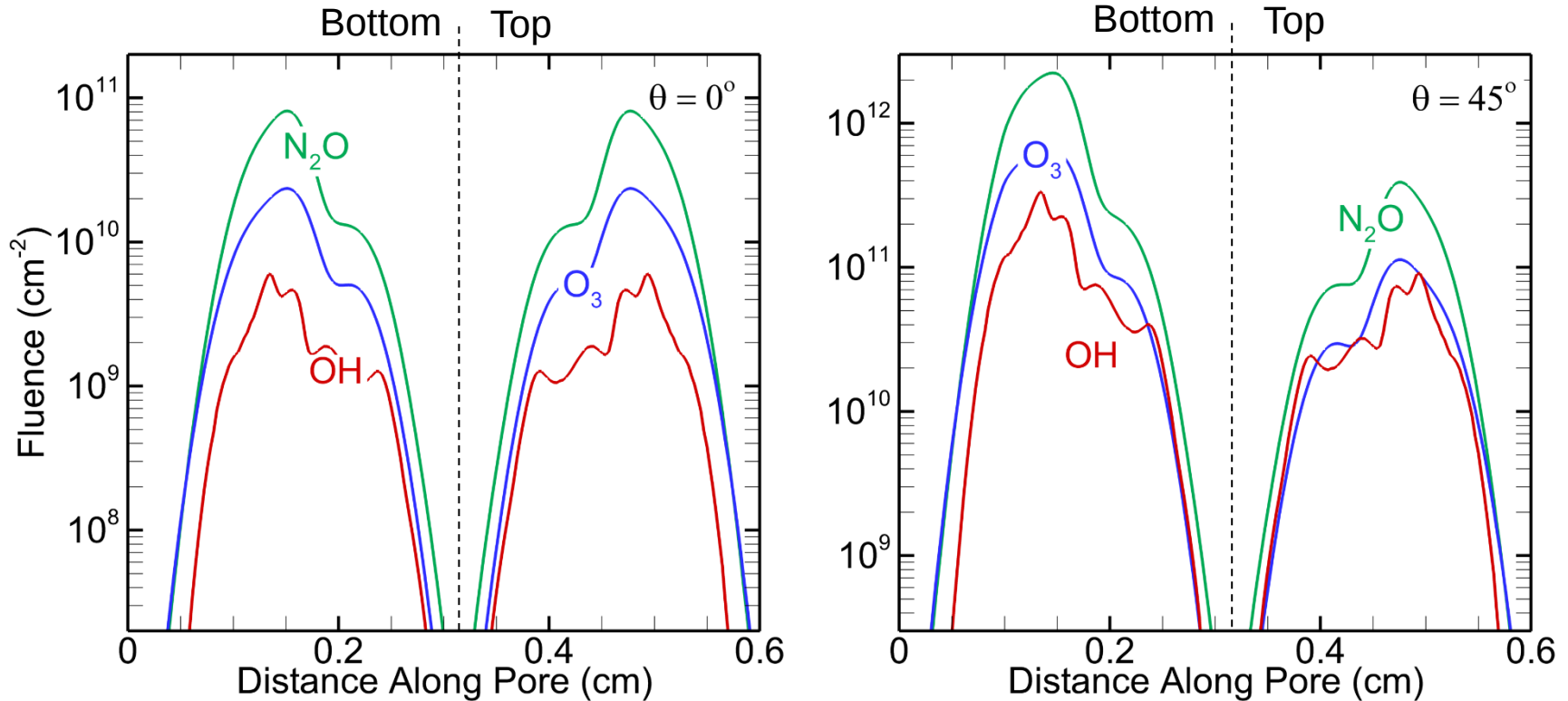
Animation Slide

FLUENCES: IONS



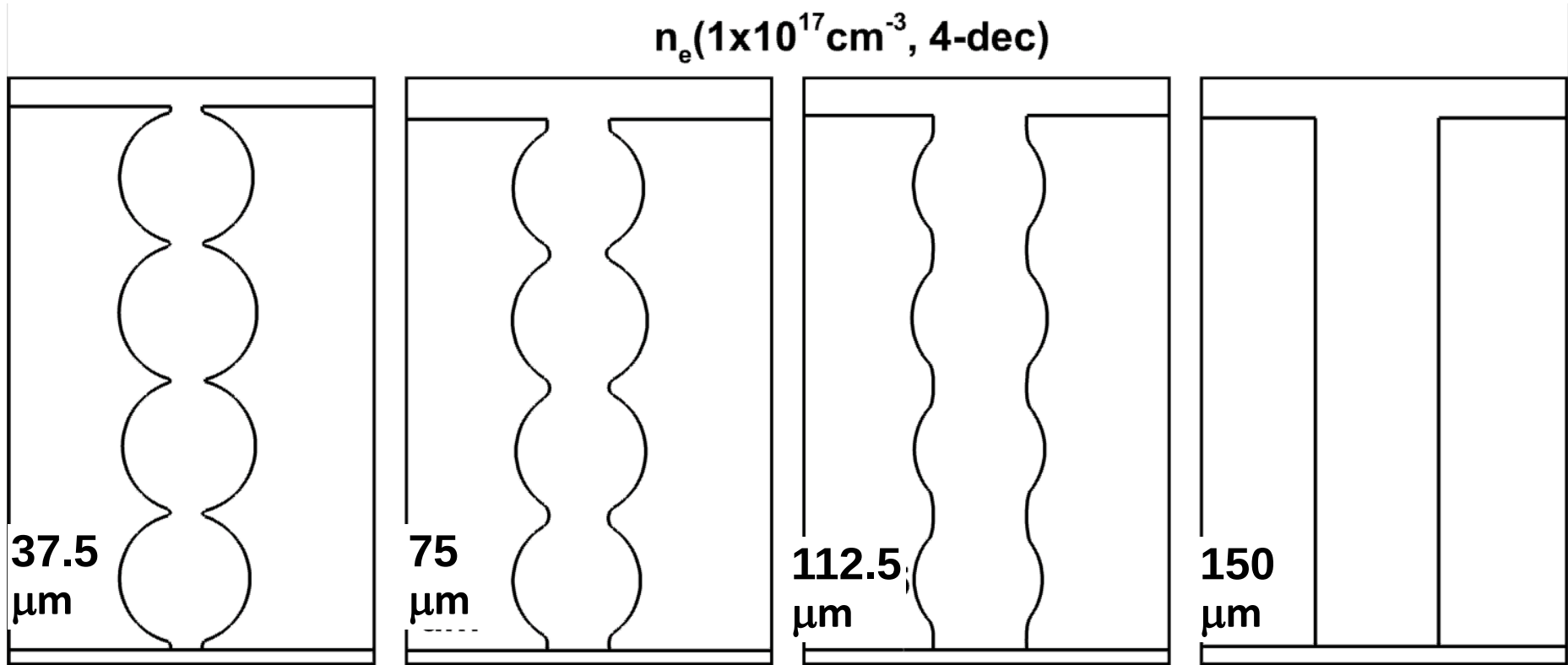
- **Distribution of fluences of selected ions along pore walls for incline angles of 0 and 45 degrees at $t = 10$ ms.**
- **Symmetric fluences to both walls at $\theta = 0$ deg.**
- **Ions recombine within microseconds, fluences drop.**
- **Pulse-on E-fields dictate final distributions of fluences.**

FLUENCES: RONS



- **Distribution of fluences of long-term species along pore walls for incline angles of 0 and 45 degrees at $t = 10$ ms.**
- **With time, asymmetry relaxes – species react in the gas-phase and fluxes of neutral RONS are driven by diffusion.**

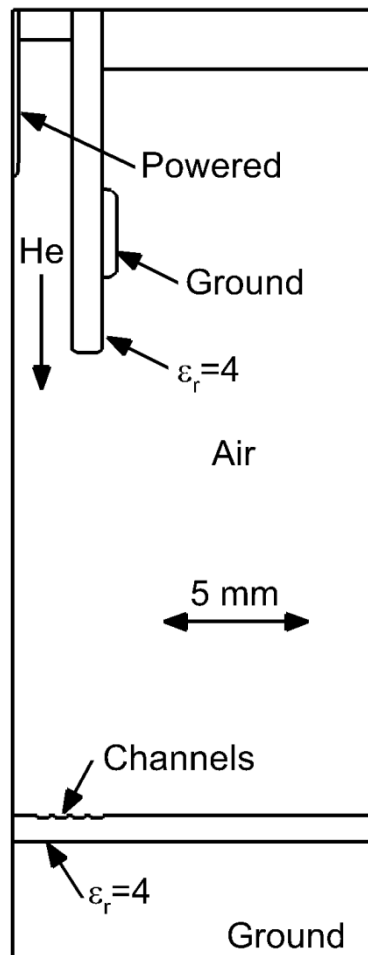
PORE OPENING – WAVY IN THE VERTICAL



- Larger pore opening results in less shadowing and less prominent restrike SIW.
- With opening larger than SIW thickness, SIWs become independent. Center of “pore” is not filled with plasma.
- 1 atm, humid air, -8 kV , 5 ns

JETS ON STRIPES

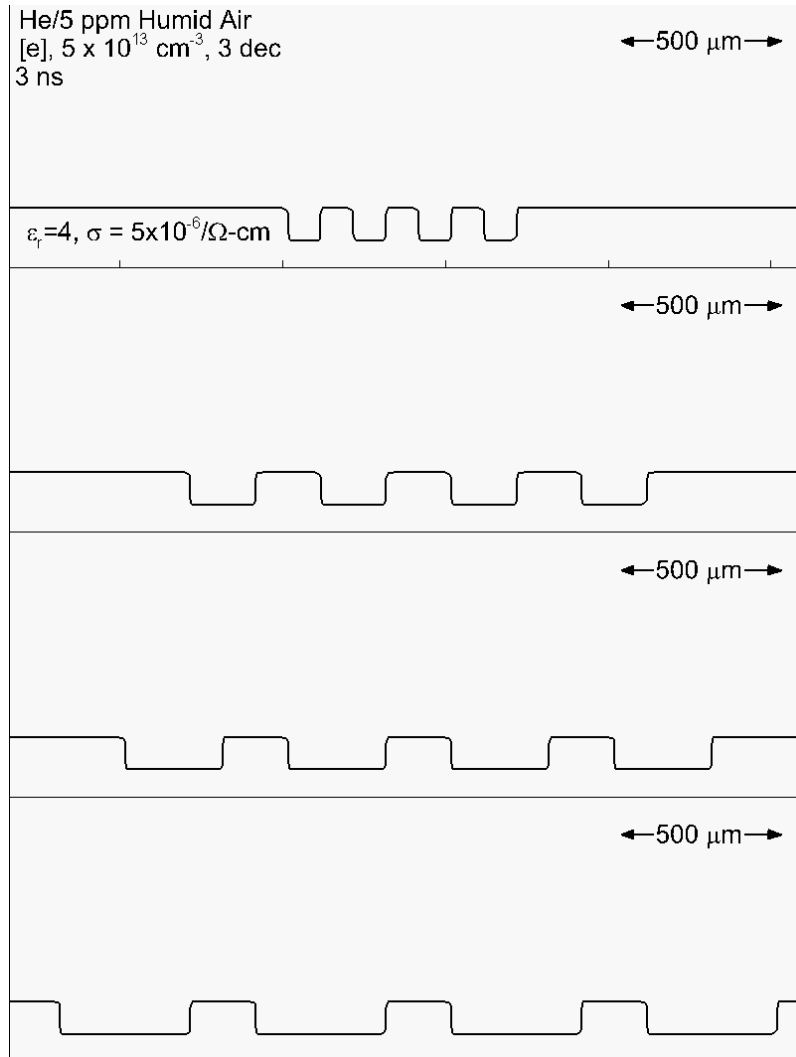
JET INTERACTIONS WITH CHANNELS



- UMD fabricated oxide-covered channels in Si

- Prelude to experiments performed with standard surfaces
- He/5 ppm ns pulsed plasma jet into humid air incident onto channels (-25 kV, 25 ns)
- Channels have varying width, spacing depth.

CHANNEL SPACING: PLASMA FILLING [e]

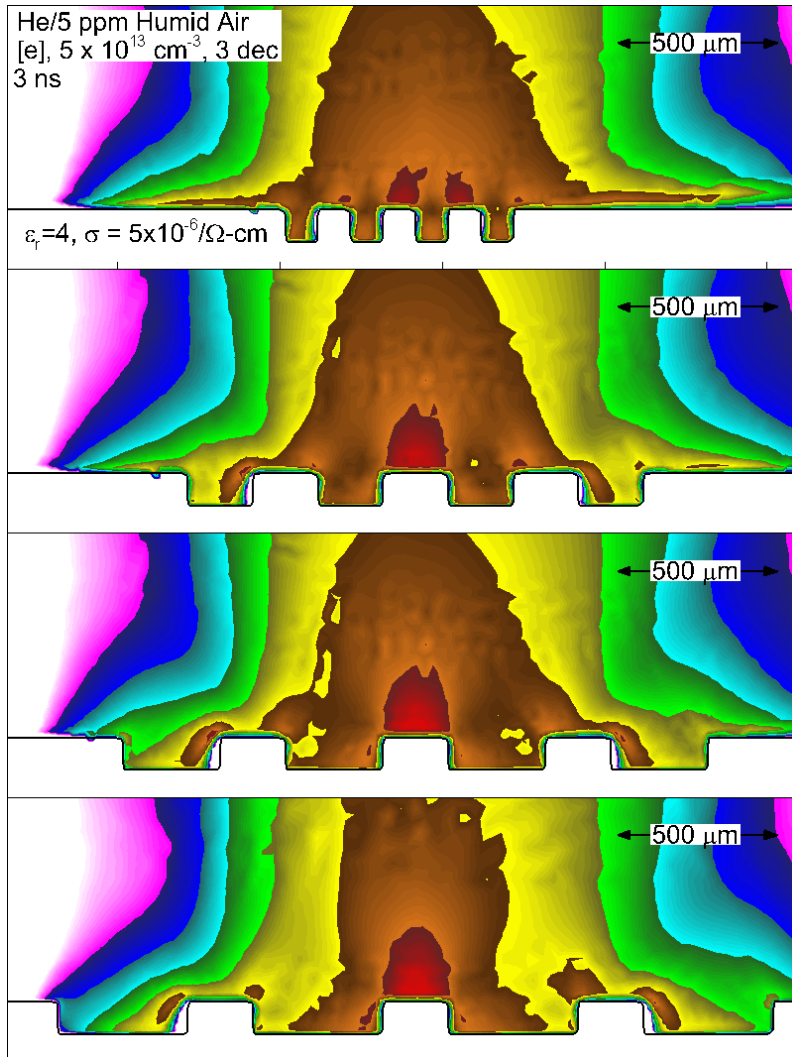


- He/5 ppm air imp. ns pulsed plasma jet into humid air incident onto channels (-25 kV, 25 ns)
- [e], $5 \times 10^{13} \text{ cm}^{-3}$ (3 dec)
- Channels depth constant (100 μm) with varying width.
- Ionization wave (IW) is attracted to apex of channels due to larger electric field above high permittivity material (smaller capacitance)
- Charging of apex diverts plasma into channels.
- Channels exceeding 250-300 μm do not uniformly fill.

Animation Slide

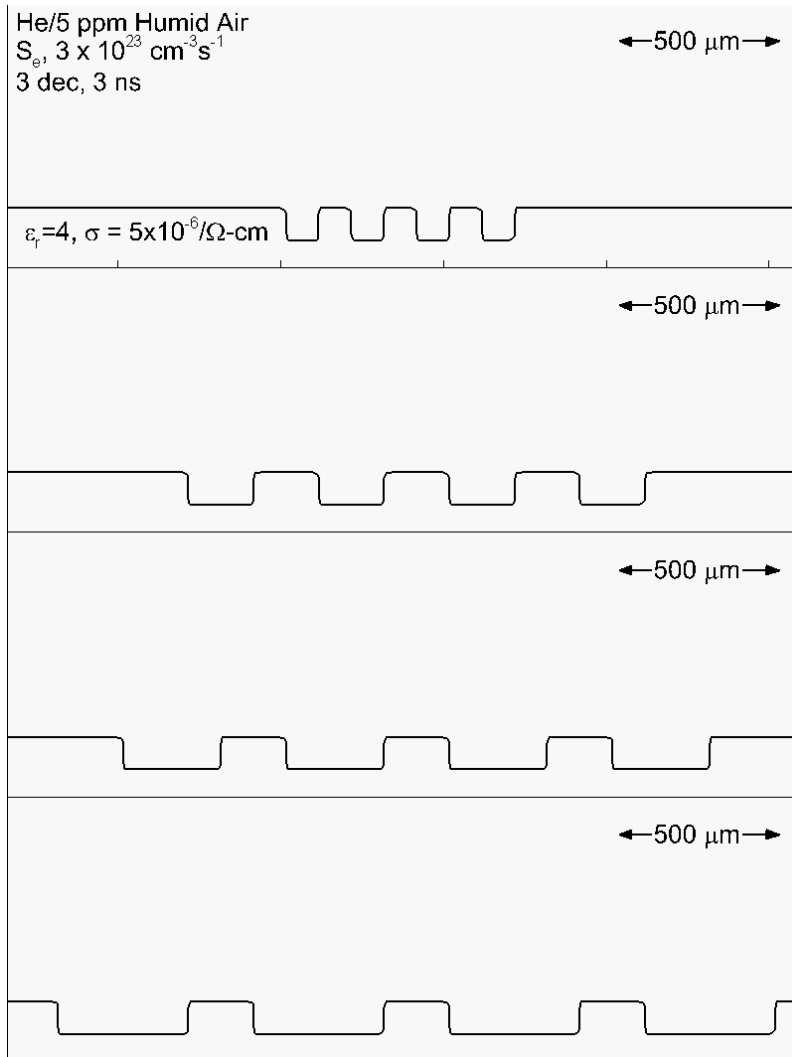
Min  Max

CHANNEL SPACING: PLASMA FILLING [e]



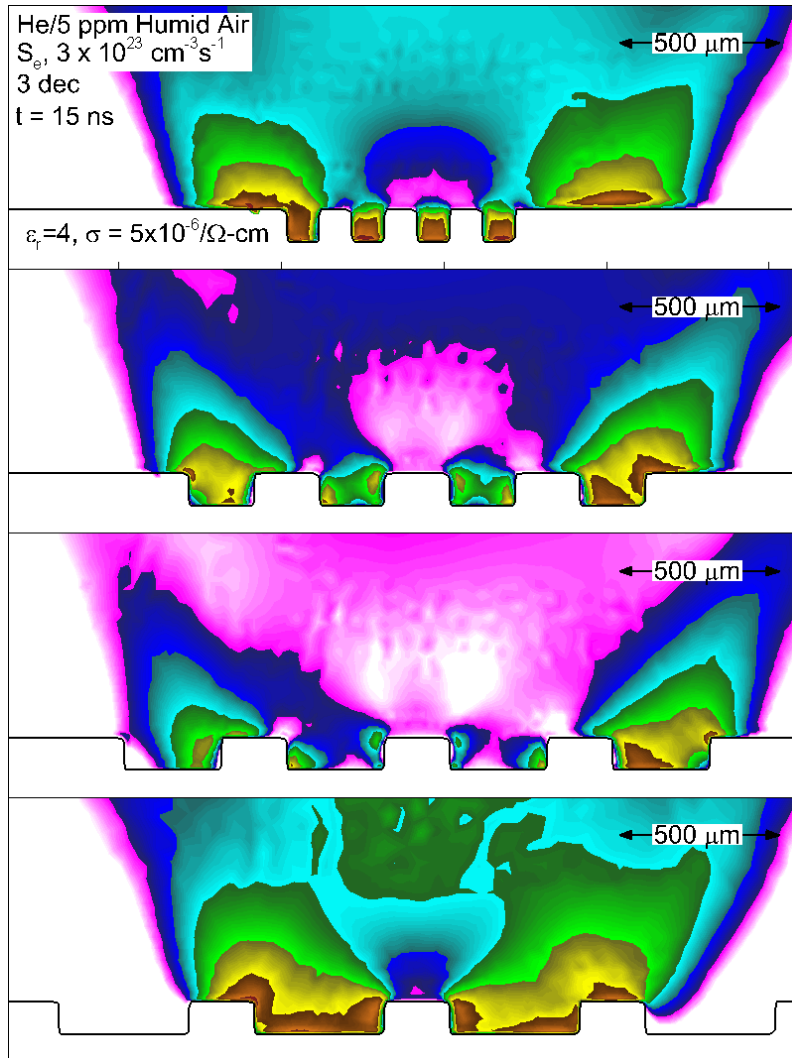
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CHANNEL SPACING: TRANSITION TO SIW - S_e



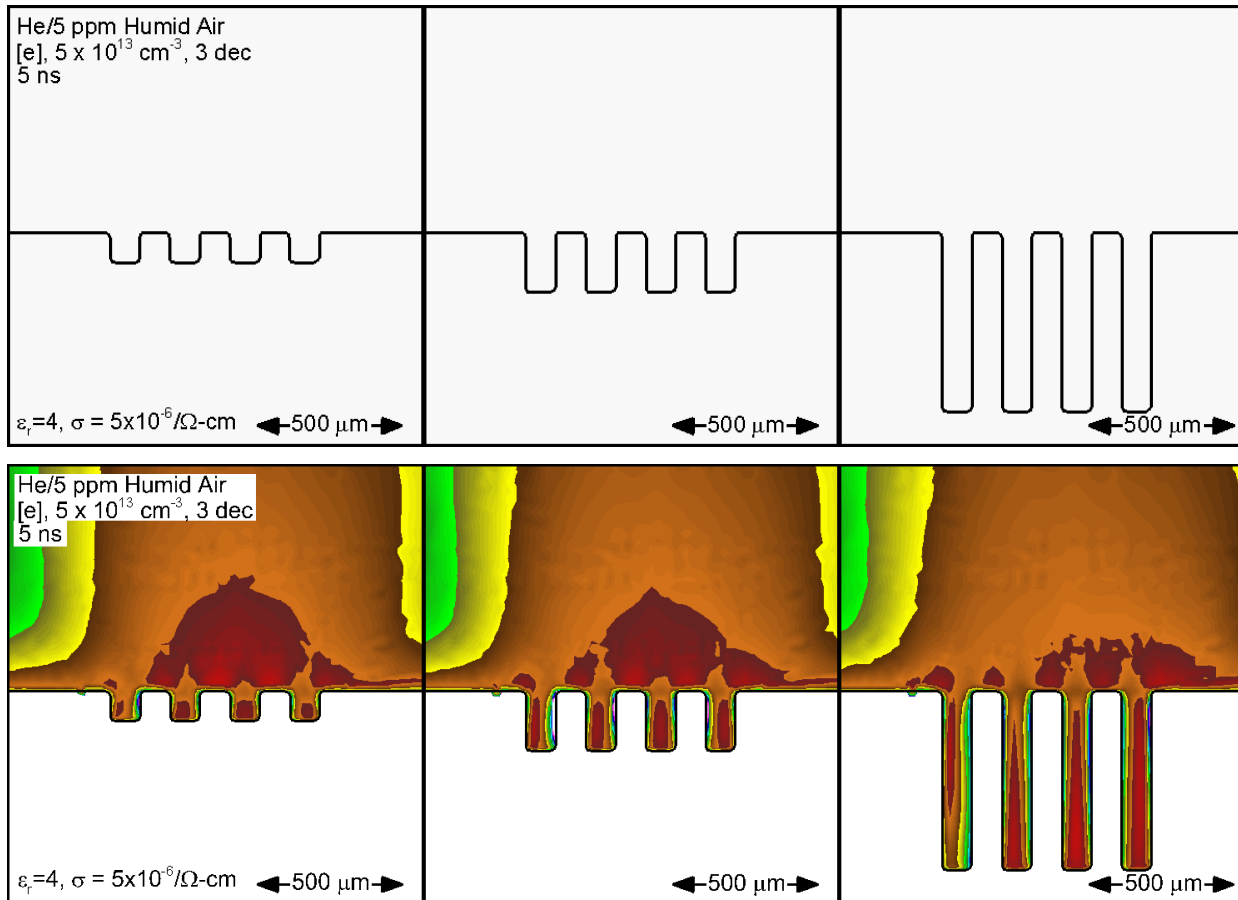
- He/5 ppm ns pulsed plasma jet into humid air incident onto channels (-25 kV, 25 ns)
- Ionization source $3 \times 10^{23} \text{ cm}^{-3} \text{ s}^{-1}$ (3 dec)
- IW penetrates the depth and fills narrower channels.
- Wider channels support small surface ionization waves.
- Electric field enhancement at corners maintains small sources after SIW passes – discharging of capacitance.

CHANNEL SPACING: TRANSITION TO SIW - S_e



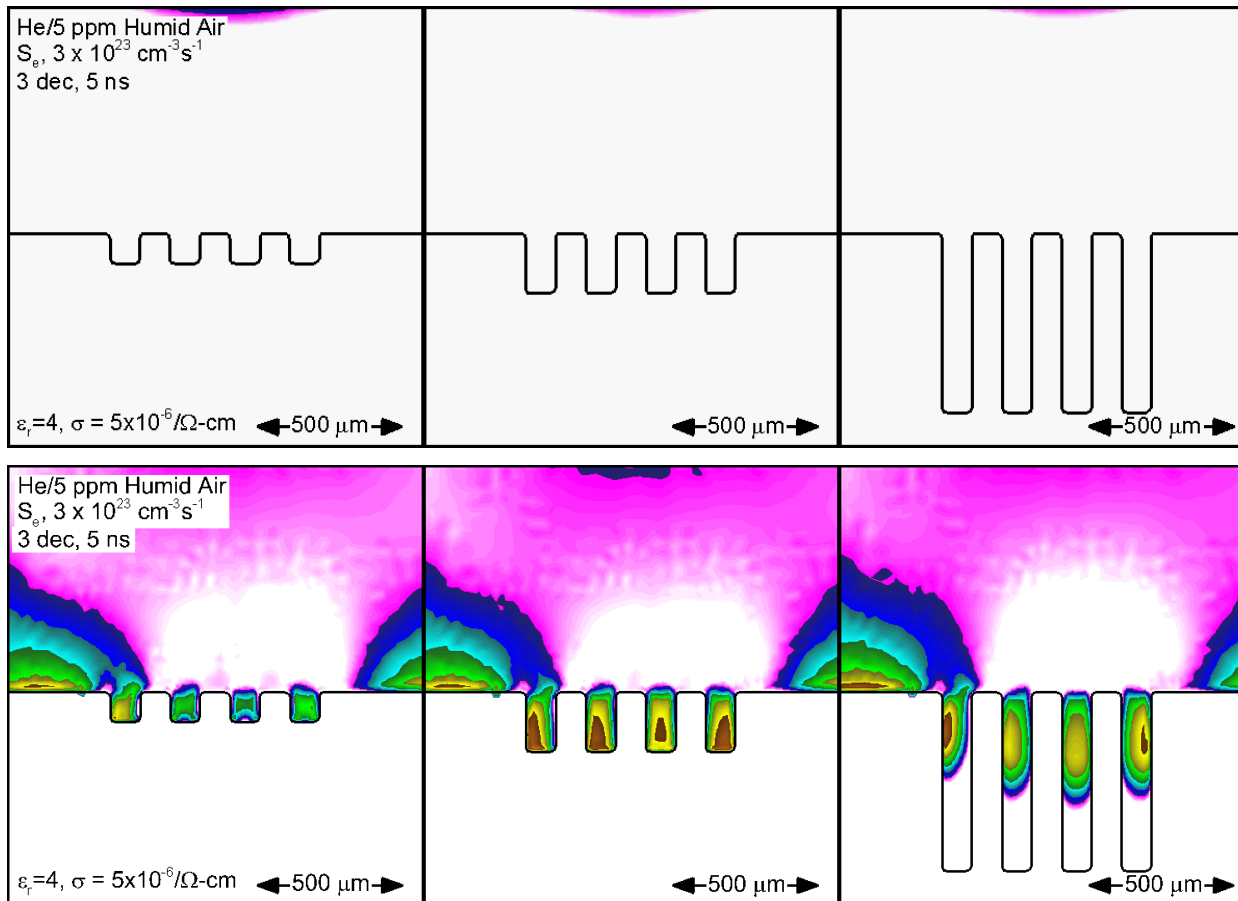
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CHANNEL DEPTH-ELECTRON DENSITY



- He/5 ppm plasma jet into incident onto channels (-25 kV, 25 ns)
- Electron density $5 \times 10^{13} \text{ cm}^{-3}$ (3 dec)
- IW penetrates channels – as SIW for deeper channels.
- Charging of channels provides energy source to sustain ionization.
- Channels see each other, pushing plasma to opposite walls.

CHANNEL DEPTH-IONIZATION SOURCE



- He/5 ppm plasma jet into incident onto channels (-25 kV, 25 ns)
- Ionization source $3 \times 10^{23} \text{ cm}^{-3} \text{ s}^{-1}$ (3 dec)
- IW penetrates channels.
- Aspect ratio > 1 , transition to SIW.
- Charging of channels provides energy source to sustain ionization.
- Channels see each other, pushing SIW to opposite walls.

CONCLUDING REMARKS

- Geometrically and materially complex surfaces have synergistic interactions with atmospheric pressure plasmas (APPs)
- The short mean free paths of electrons and ions localize surface complexities, that in turn intensify feedback.
- The surface-hugging nature of SIWs can be compromised by surface roughness that (combined with high ϵ_r) produces critically large local field enhancement that initiates hopping.
- Surface charging is not independent of topology and is responsible for spatially dependent restrike SIWs...
- ...And the relationship between APPs and complex dielectric surfaces is all about charging.
- Reactant fluxes of neutral reactive species are far more uniform than ions.
- However, the topology dependent differential charging of surfaces likely plays a role in catalytic reactivity of those neutrals.