



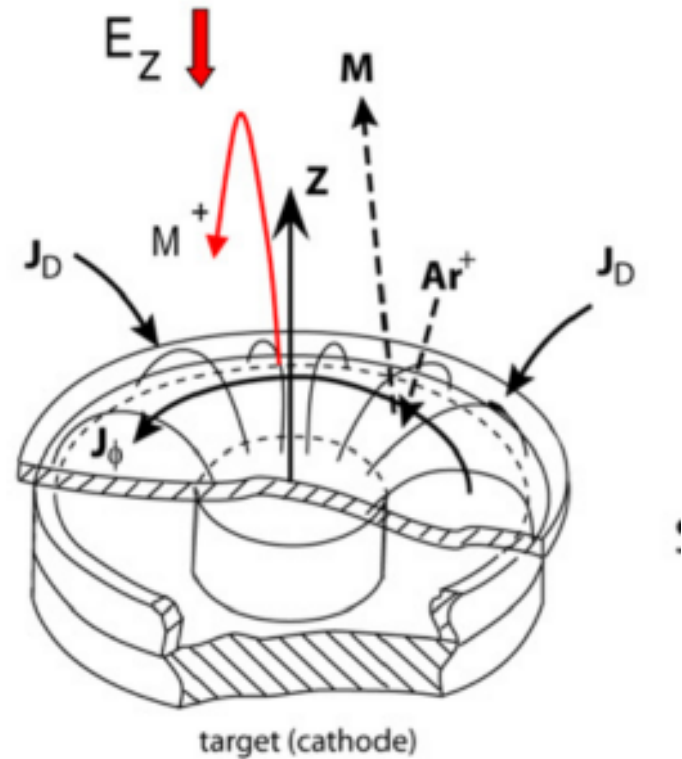
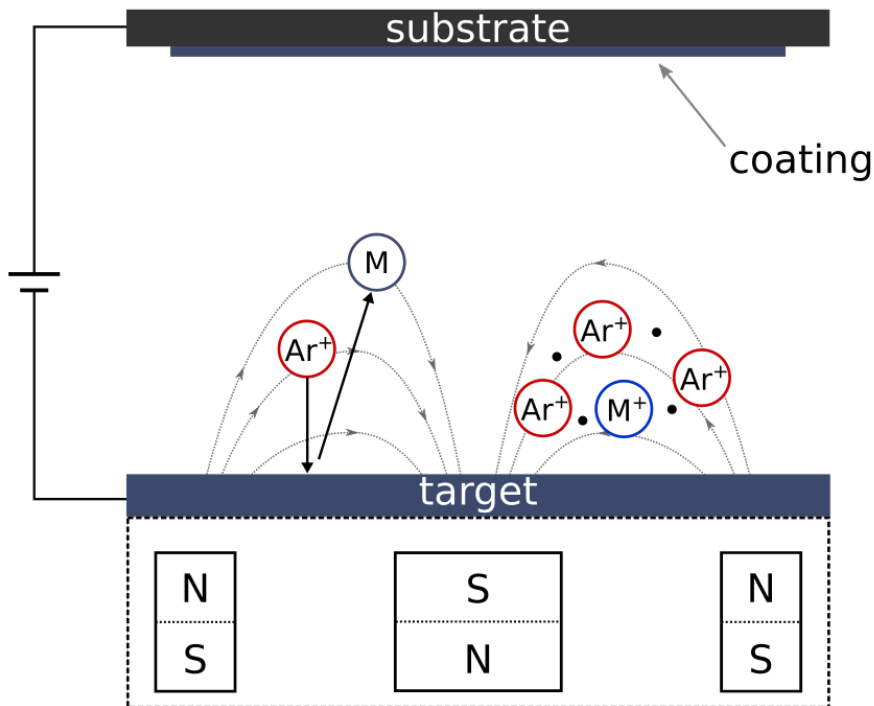
RUHR-UNIVERSITÄT BOCHUM

Azimuthal Ion Movement in HiPIMS: Hall Current or Spokes?

S Thiemann-Monjé, J Held, S Schüttler and A von Keudell

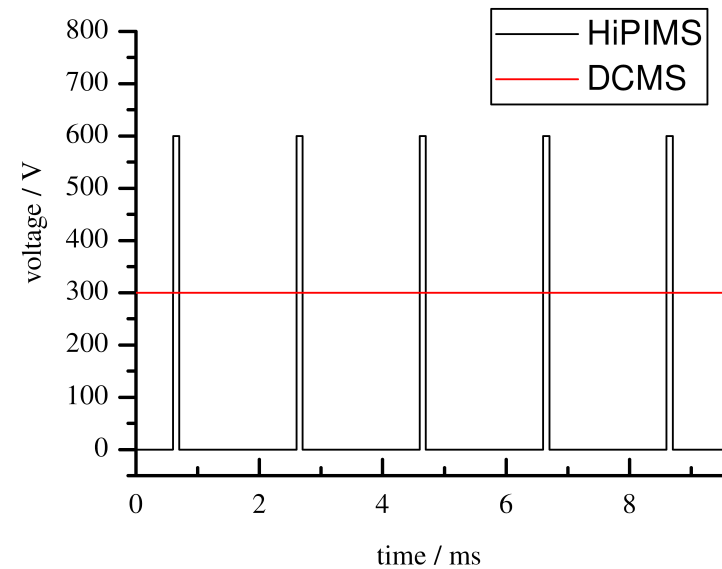
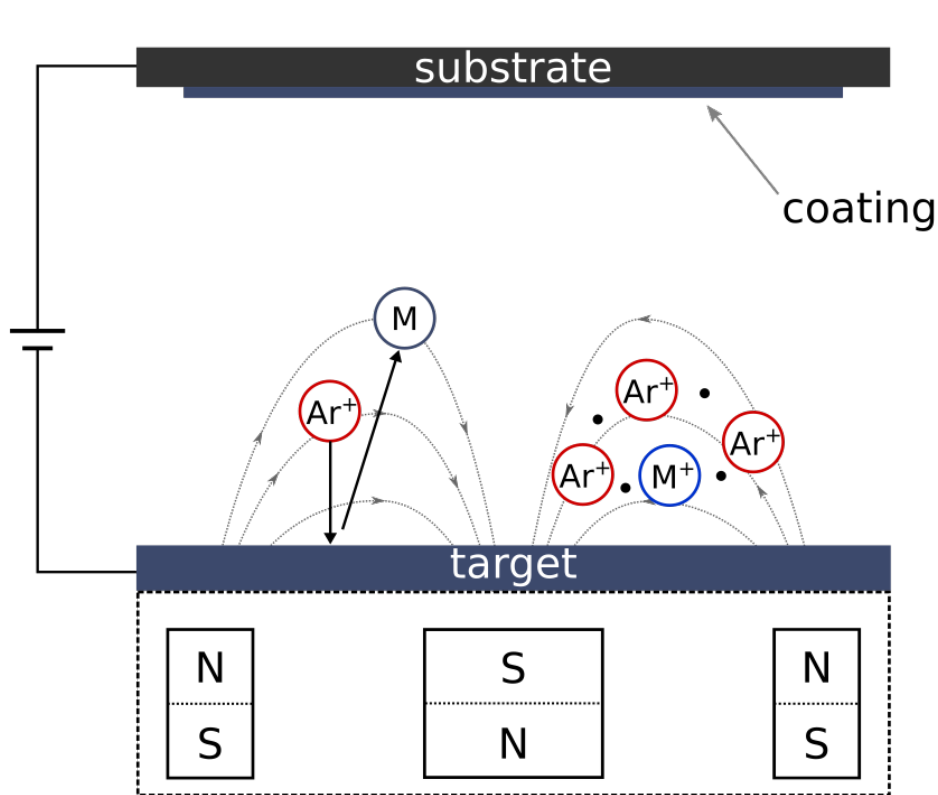


Introduction | High Power Impulse Magnetron Sputtering



N. Brenning et. al., PSST 21 025005 (2012)

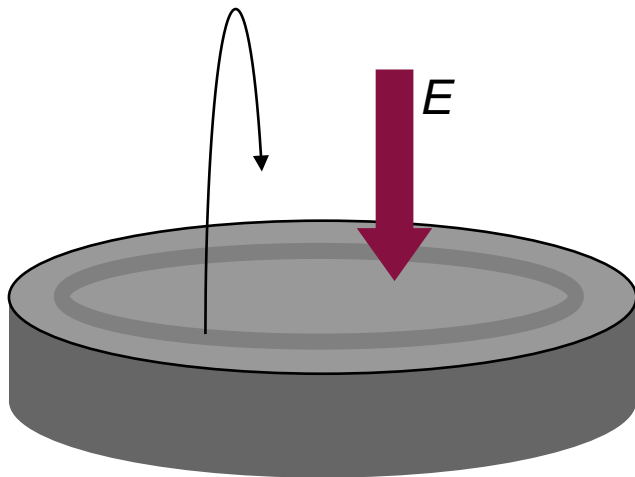
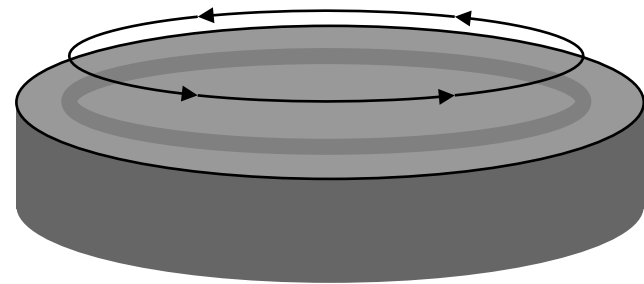
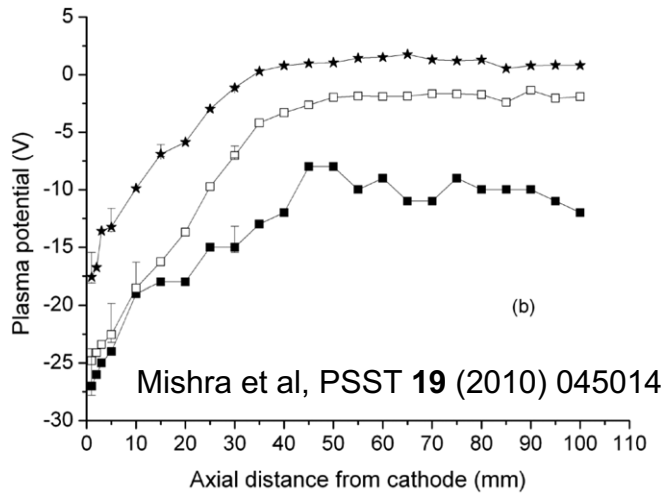
Introduction | High Power Impulse Magnetron Sputtering



HiPIMS:
Higher ionization degree → better coatings.

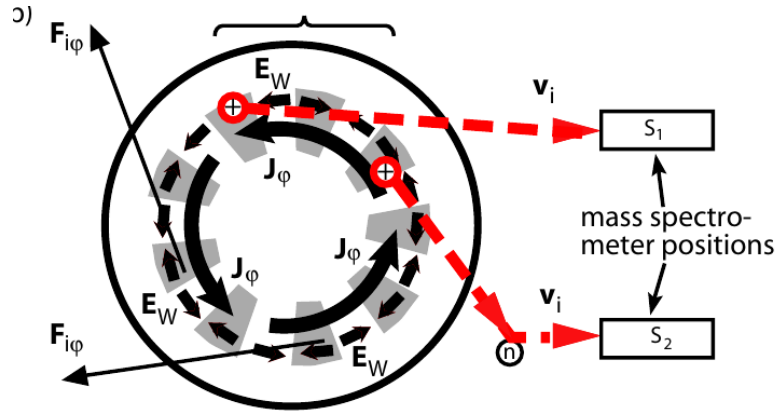
Ion movement is much more complex than neutrals.

Introduction | Ion transport in HiPIMS

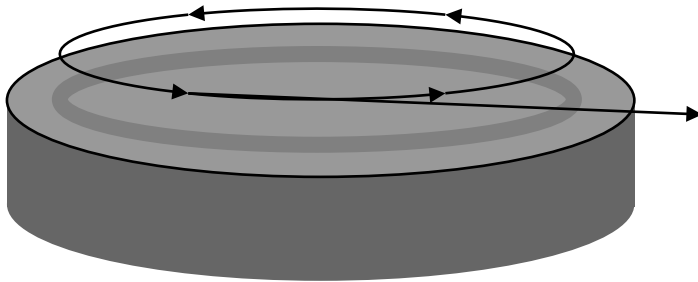
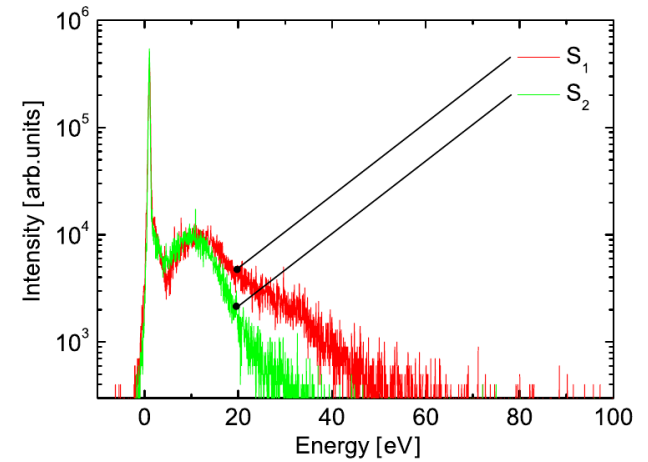


Ions also move in azimuthal direction.

Introduction | Azimuthal Ion Movement

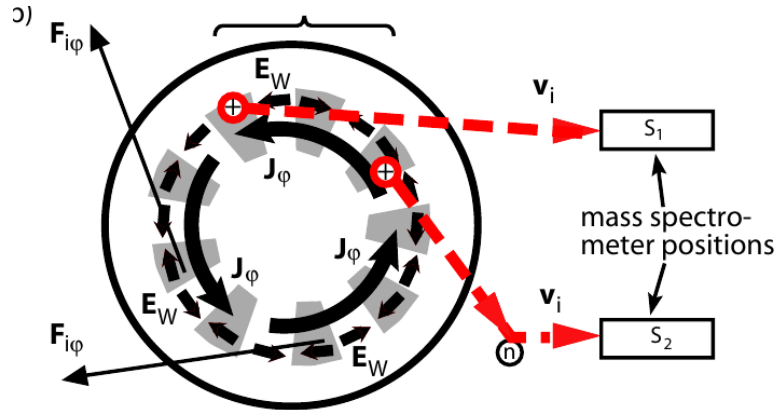


Lundin *et al*, Plasma Sources
Sci. Technol. **17** (2008) 035021



Ions are seemingly *flung out*
of the racetrack.

Introduction | Azimuthal Ion Movement: Why?



Lundin *et al*, Plasma Sources
Sci. Technol. **17** (2008) 035021

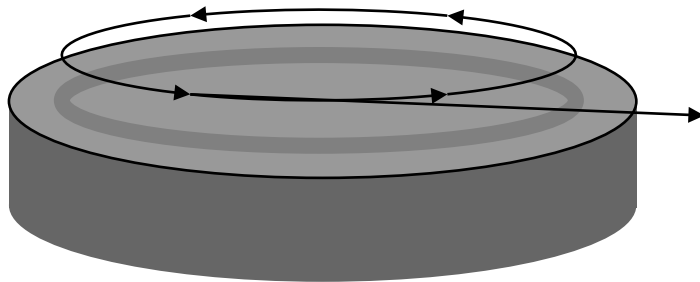
Hall current caused by electron drifts:

$$\mathbf{v}_D \equiv -\frac{\nabla p \times \mathbf{B}}{qnB^2} \quad \mathbf{v}_E \equiv \frac{\mathbf{E} \times \mathbf{B}}{B^2}$$

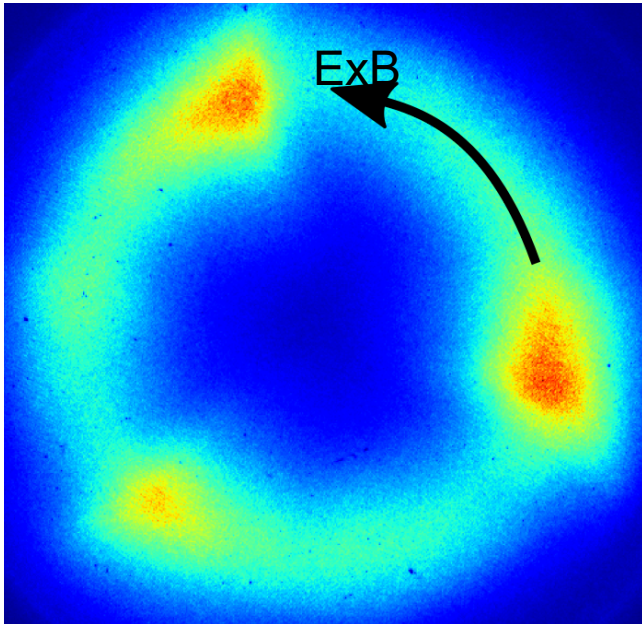
$$F_{i,\phi} = -\eta e^2 n (v_D + v_{E \times B})$$

In HiPIMS, the Hall current is usually
two or three times the discharge
current.

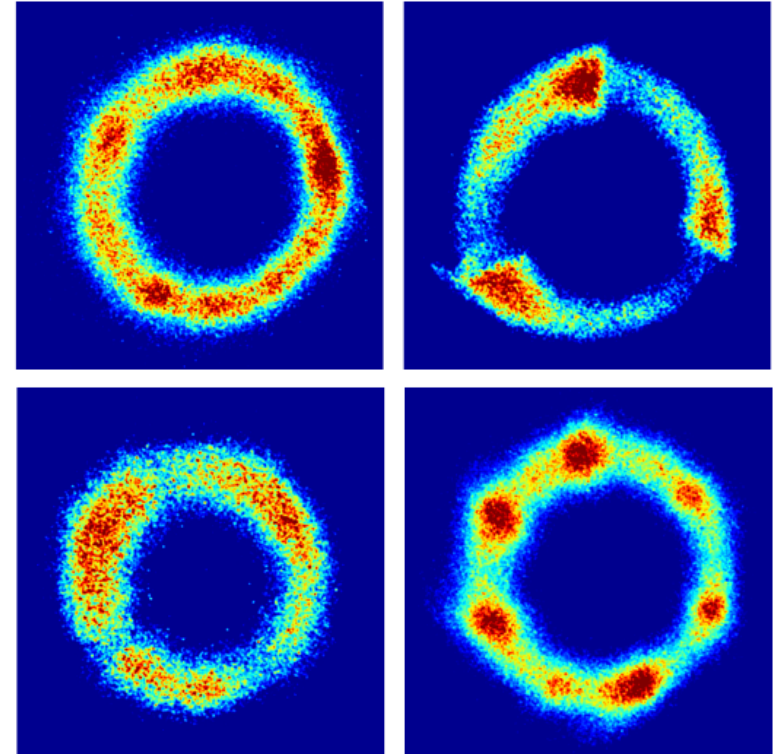
$$I_D = \sim 50 \text{ A} \rightarrow I_{\text{Hall}} = \sim 100 \text{ A}$$



Introduction | Spokes



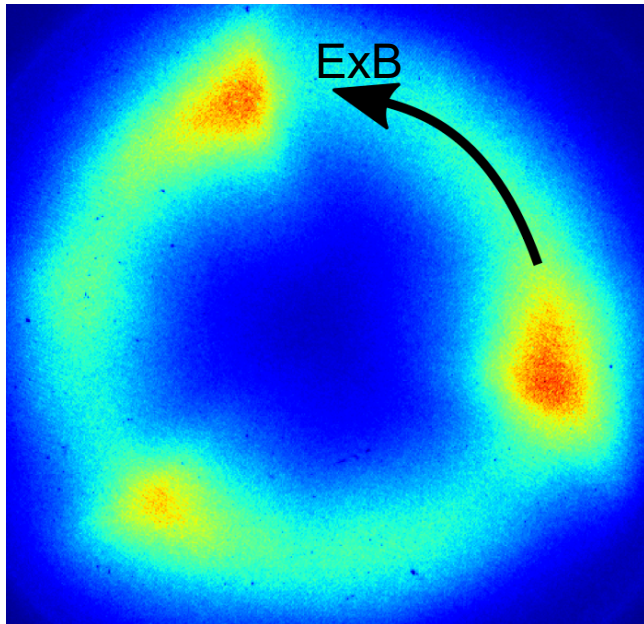
Rotating with 5 – 10 km/s in ExB direction



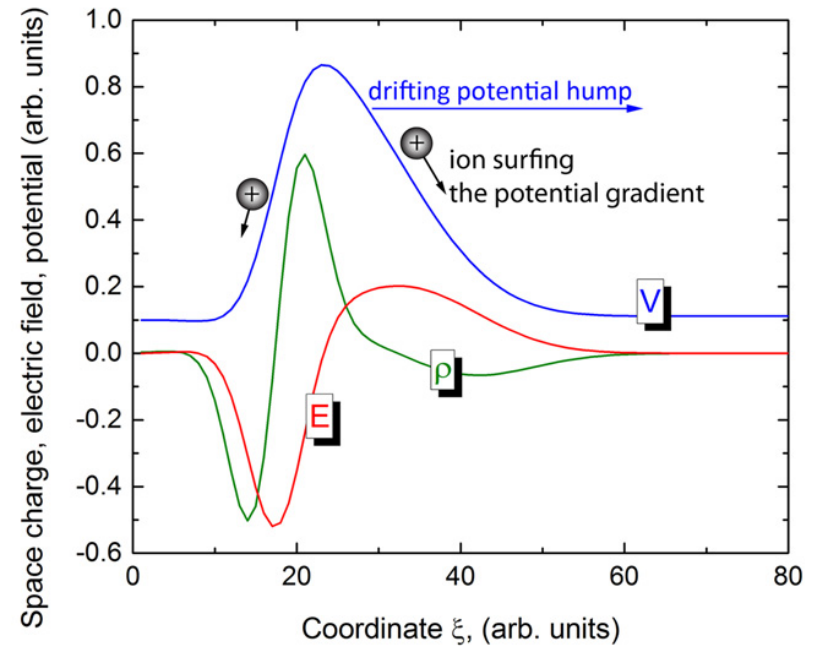
Hnilica, Klein, Šlapanská, Fekete, Vašina
J. Phys. D: Appl. Phys. **51** (2018) 095204

Spokes: wave phenomenon,
moving in ExB direction.

Introduction | Spokes



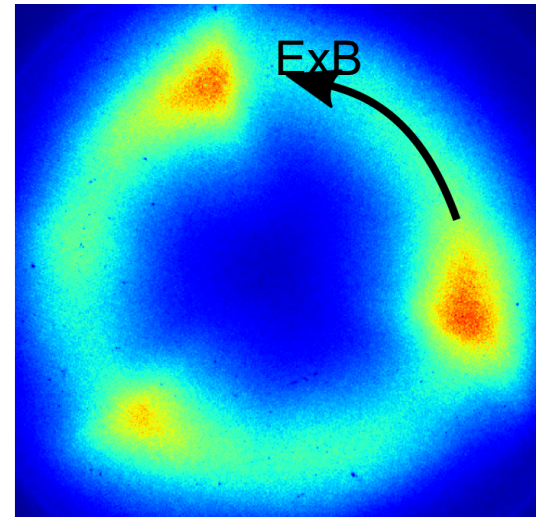
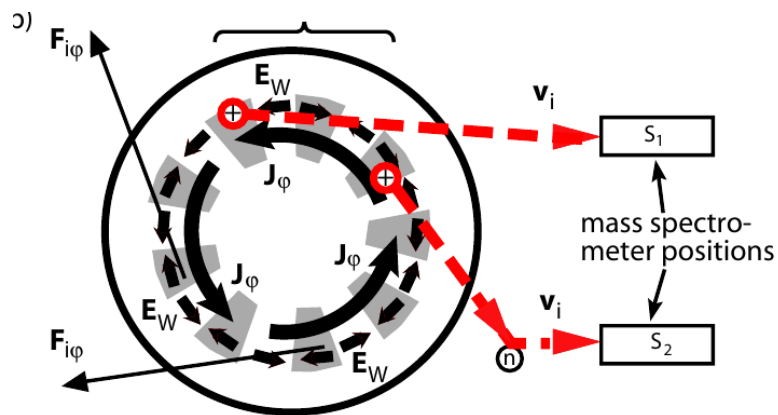
Rotating with 5 – 10 km/s in $E \times B$ direction



Panjan, Franz, Anders (2014) PSST **23** 025007

Spokes can also cause azimuthal ion movement.

Introduction | Azimuthal Ion Movement in HiPIMS

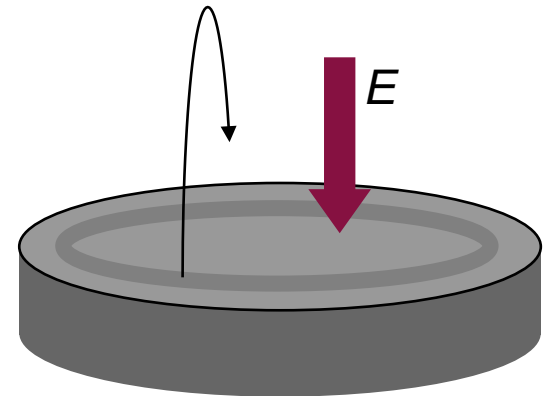
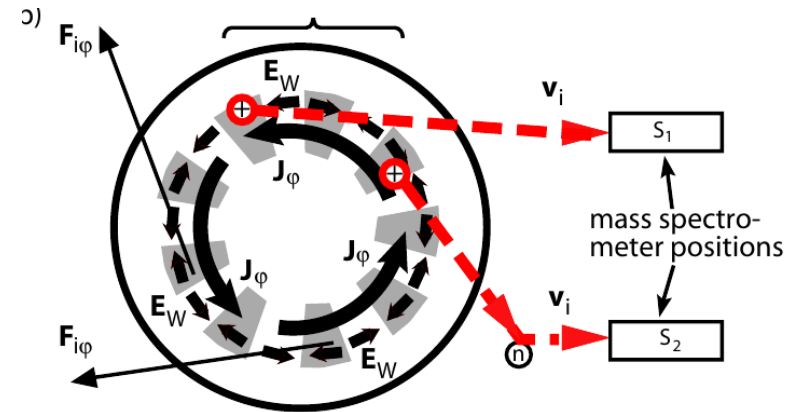


Hall current or Spokes?

Introduction | Azimuthal Ion Movement in HiPIMS

Hall current or spokes: What do we need?

1. Better measurements of ion movement.
2. An estimation of the Hall-current drag force
3. Quantitative measurements of Spoke electric field and density.



PART ONE

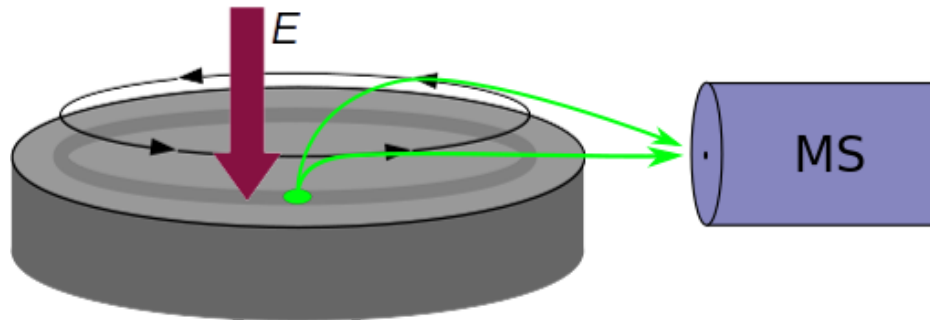
MEASUREMENT OF AZIMUTHAL ION VELOCITY

Measurement of Azimuthal Ion Velocity | MS measurements

Problems with mass spectrometry :

Measurement outside of plasma

Pre selection of particles due to time varying E-field



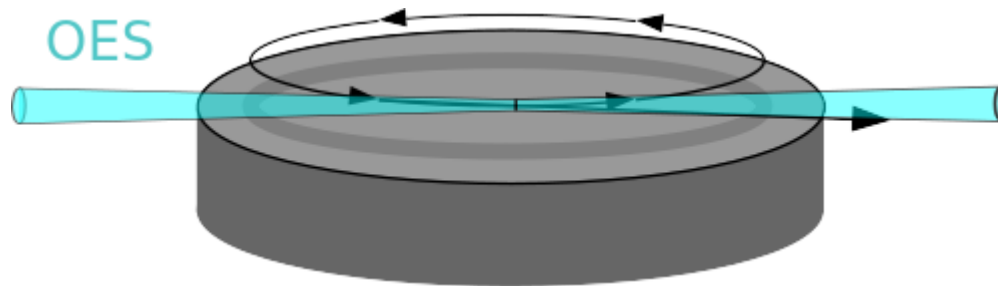
Measurement of Azimuthal Ion Velocity | OES

Using Optical Emission Spectroscopy:

non-intrusive

Spatially resolved

Time resolved



Measurement of Azimuthal Ion Velocity | OES

Optical Emission line profile:

Natural: ~ 0.1 pm

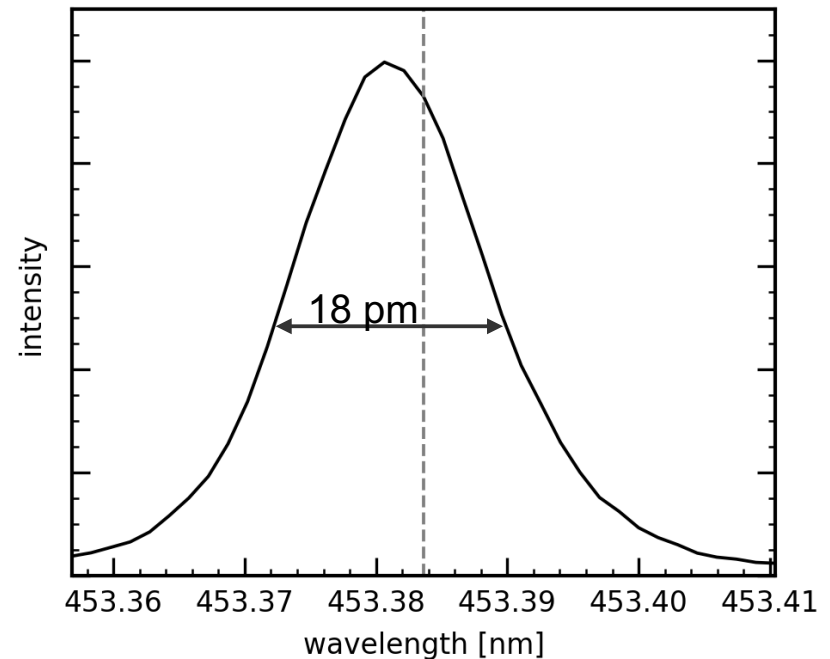
Stark: < 1 pm

Zeeman: ~ 1 pm

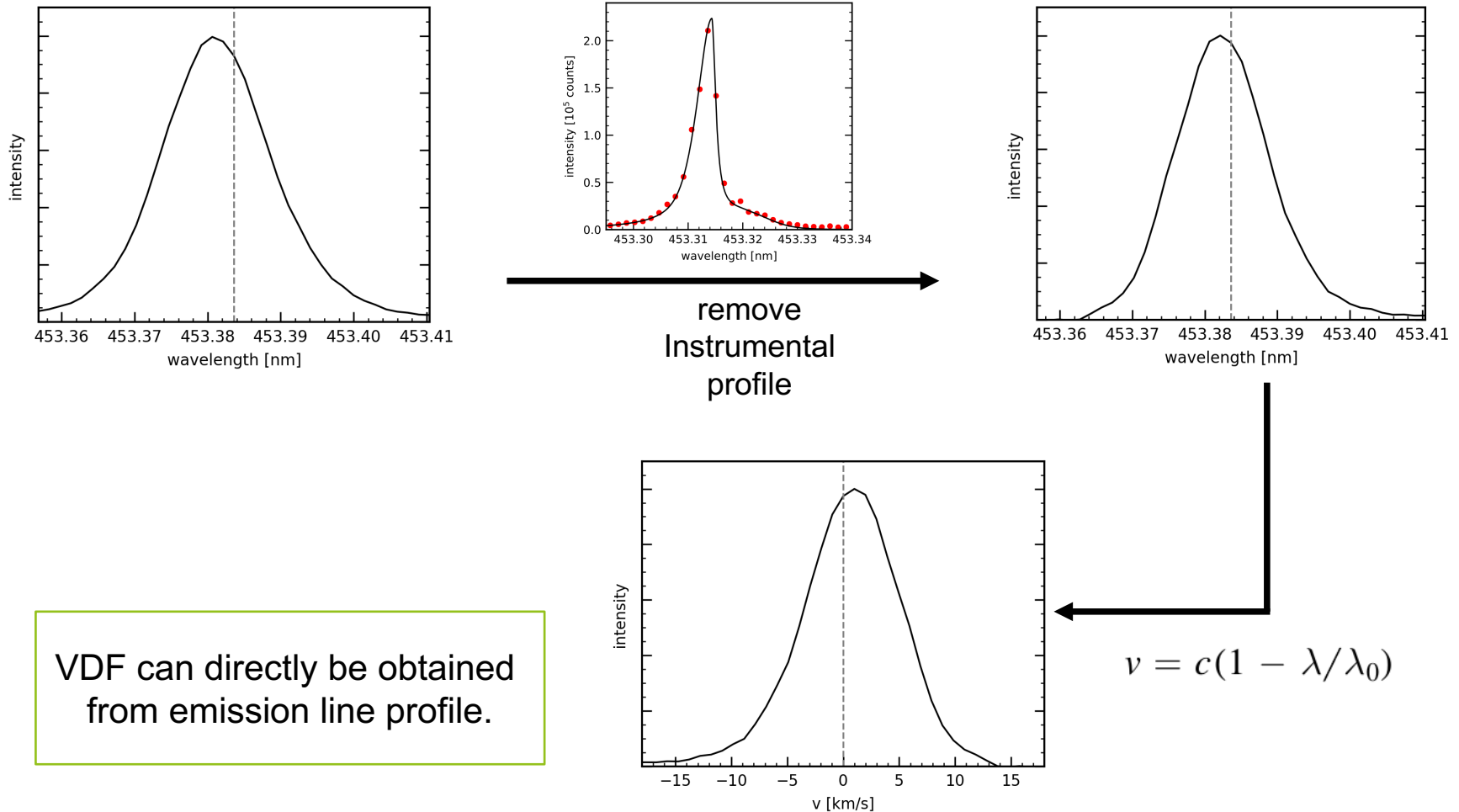
Instrument: ~ 6 pm

Doppler: ~ 15 pm

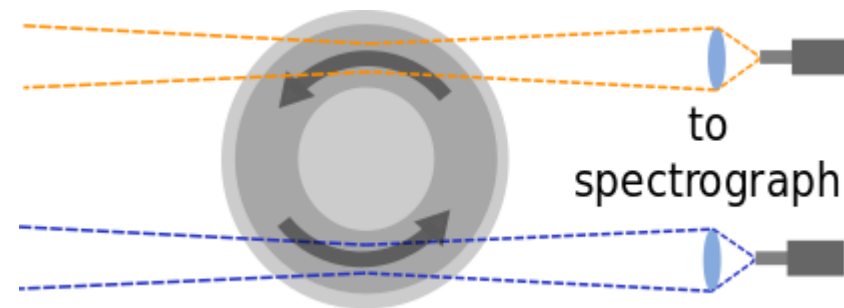
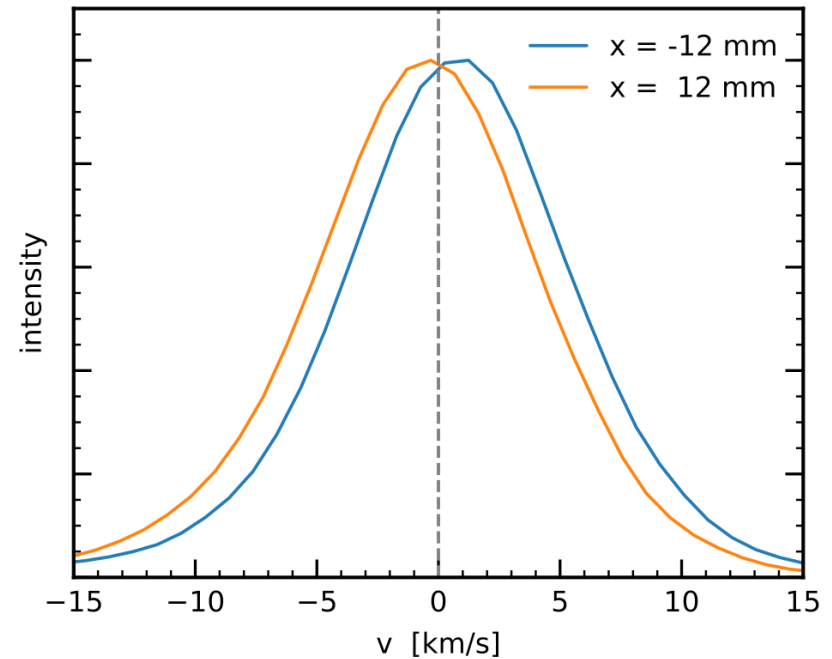
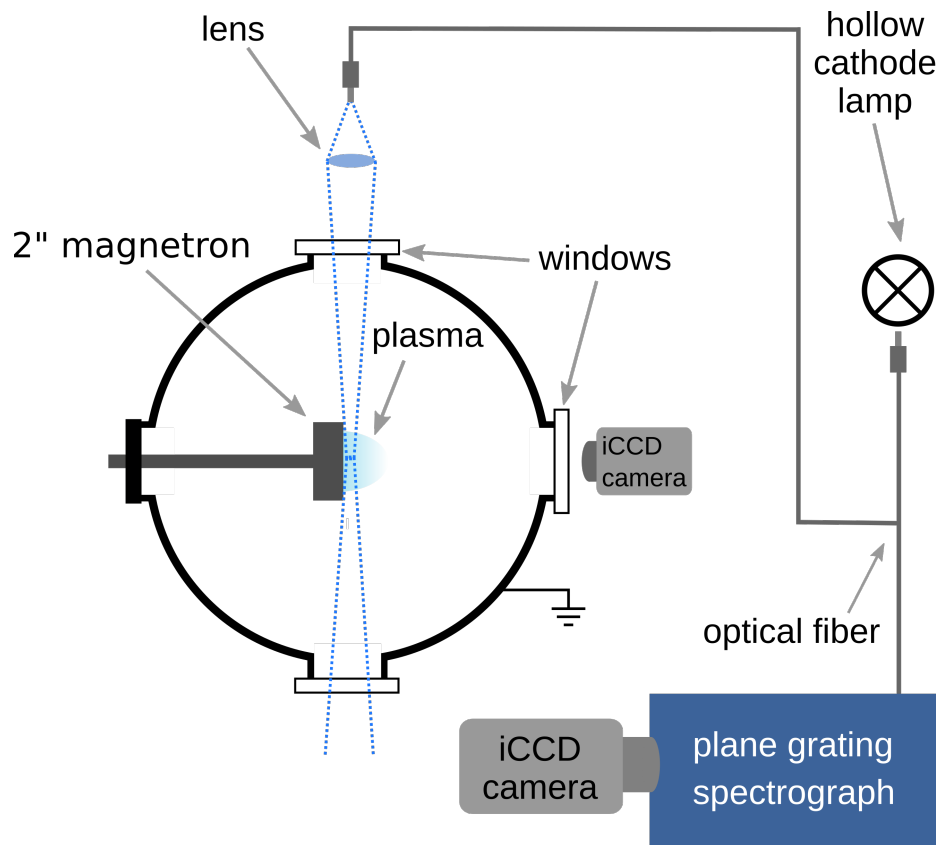
Line profile is dominated
by Doppler broadening.



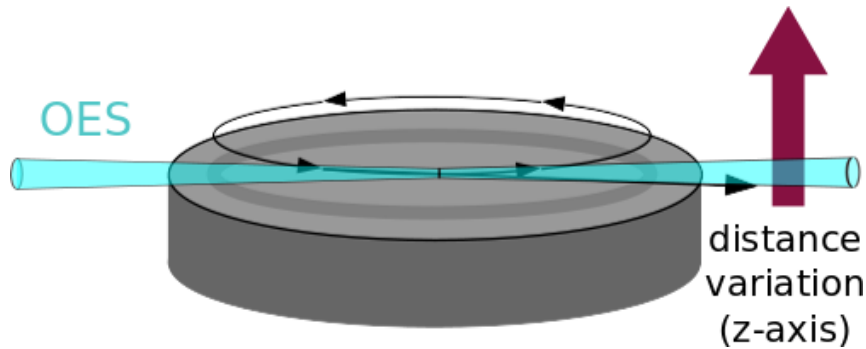
Measurement of Azimuthal Ion Velocity | OES



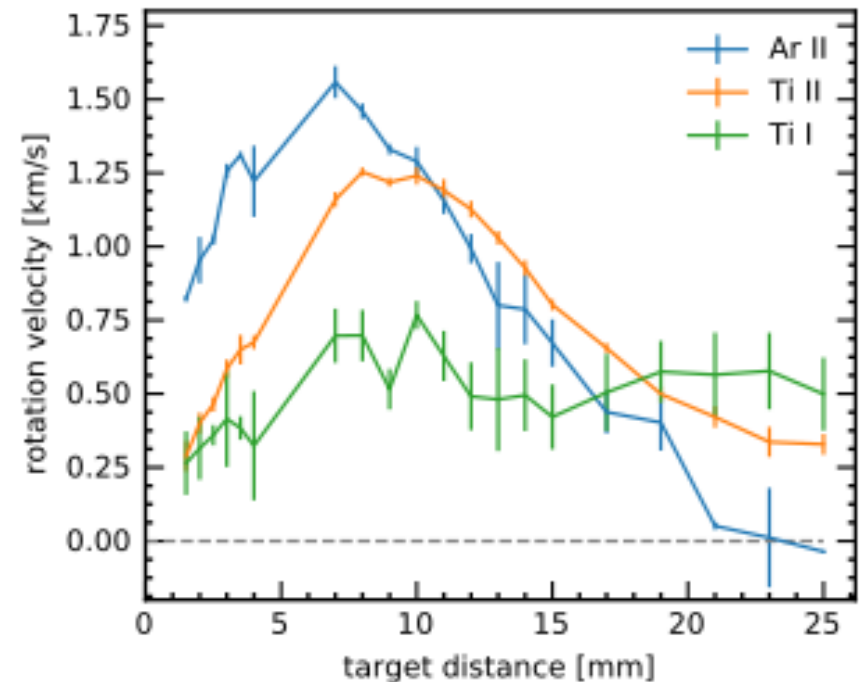
Measurement of Azimuthal Ion Velocity | OES



Measurement of Azimuthal Ion Velocity | OES



Maximum Ion rotation velocity is about 1.2–1.5 km/s at 8–10 mm distance from target



PART TWO

HALL CURRENT

Part Three | Azimuthal force

Parameters of Hall current drag force:

B – measured (plasma off)

T_e – measured 4.5 eV (z=8mm)

E – estimation needed

n_e – estimation needed

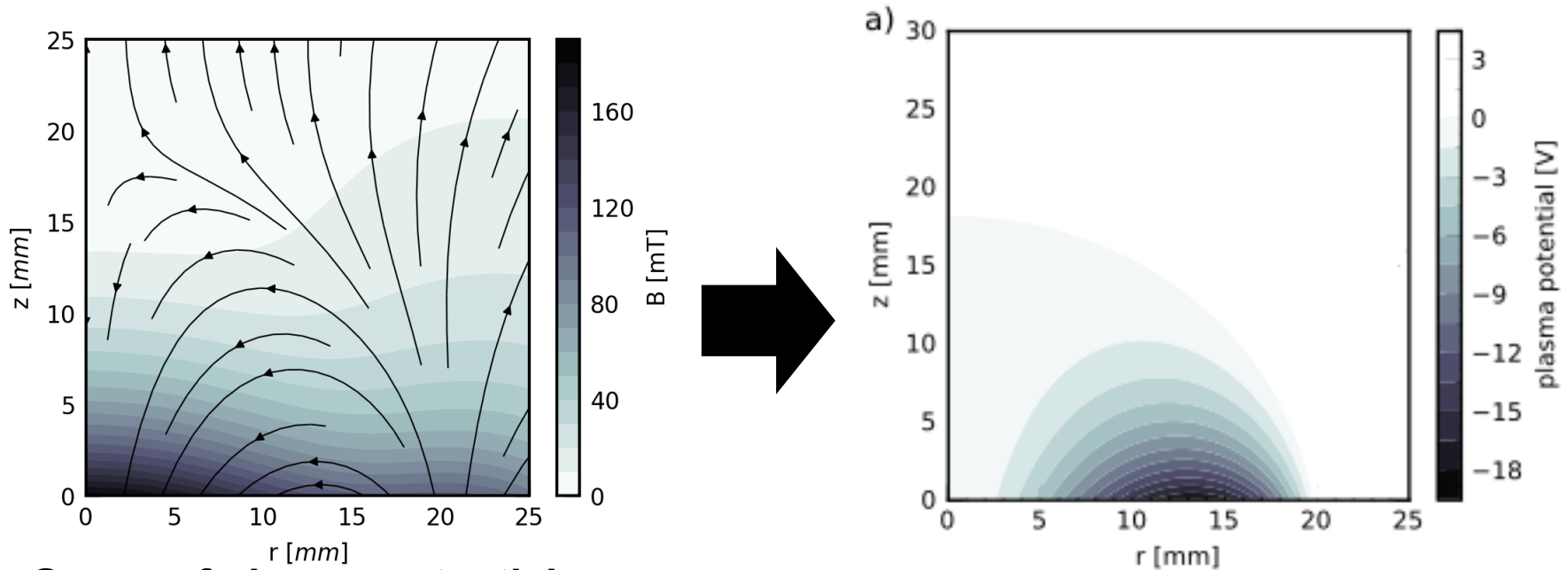
$$F_{i,\phi} = -\eta e^2 n (v_D + v_{E \times B})$$

$$\mathbf{v}_D \equiv -k_B T \frac{\nabla n \times \mathbf{B}}{qnB^2}$$

$$\mathbf{v}_E \equiv \frac{\mathbf{E} \times \mathbf{B}}{B^2}$$

Estimations of E-field and density needed

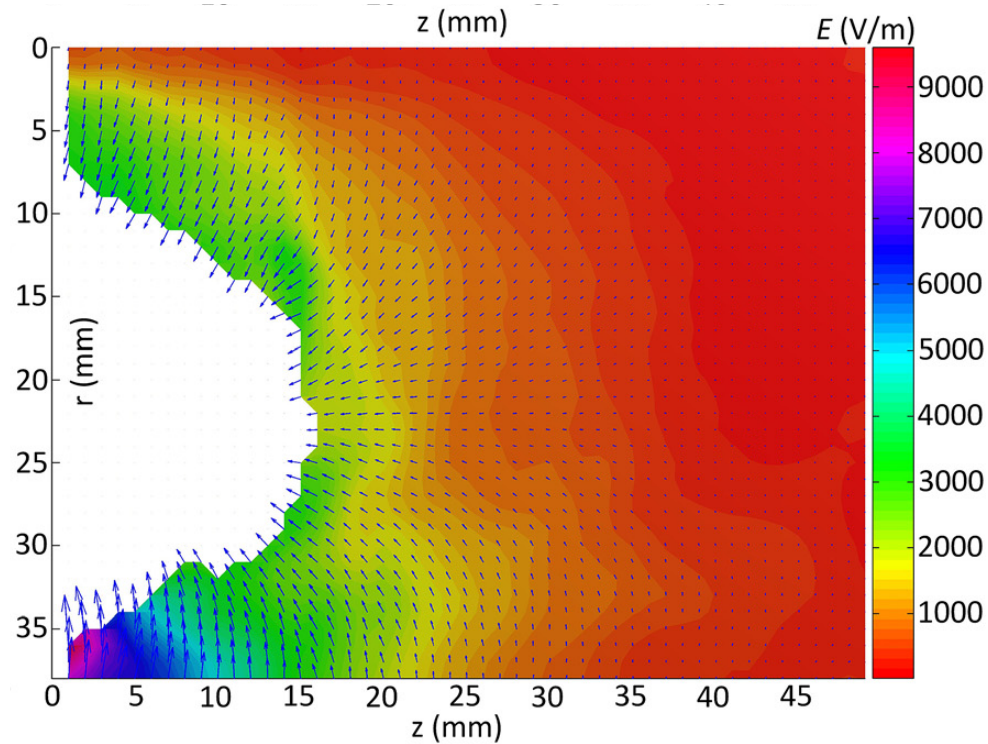
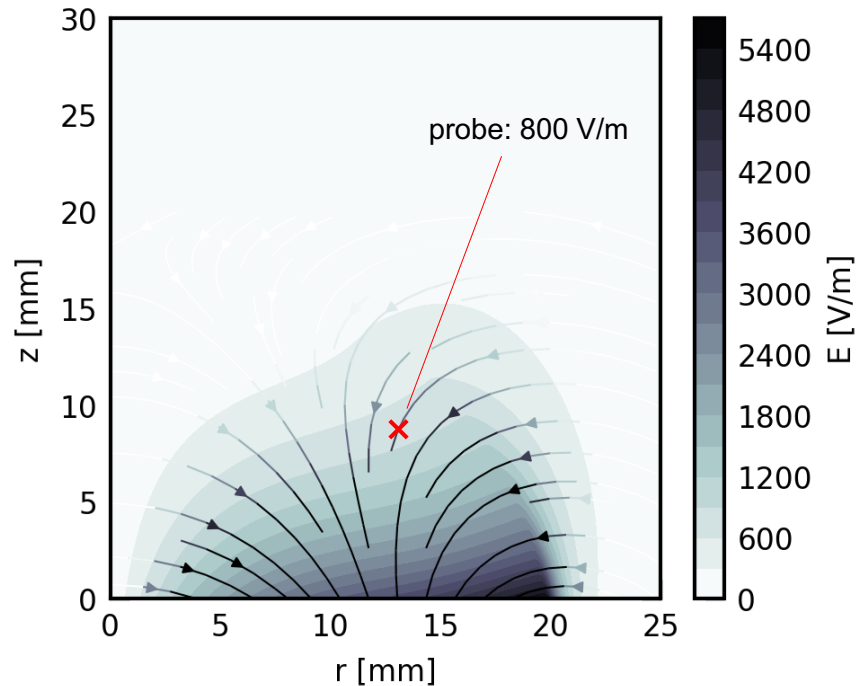
Part Three | Electric Field



Guess of plasma potential:

- No potential differences possible along B-Field lines
- Usage of magnetic flux coordinates Ψ introduced by Brinkmann and Krüger
[R. P. Brinkmann and D. Krüger, "Axisymmetric magnetically enhanced discharges described in terms of flux coordinates", Physics of Plasmas 27, 053504 (2020)]
- $\Psi = \text{const.}$ represents magnetic field lines
- Guess: plasma potential proportional to magnetic flux coordinates

Part Three | Electric Field



Rauch *et al* J Appl Phys (2012) 111 083302

E-field fits to measurements
but remains a weak estimation

Part Three | Ion Density

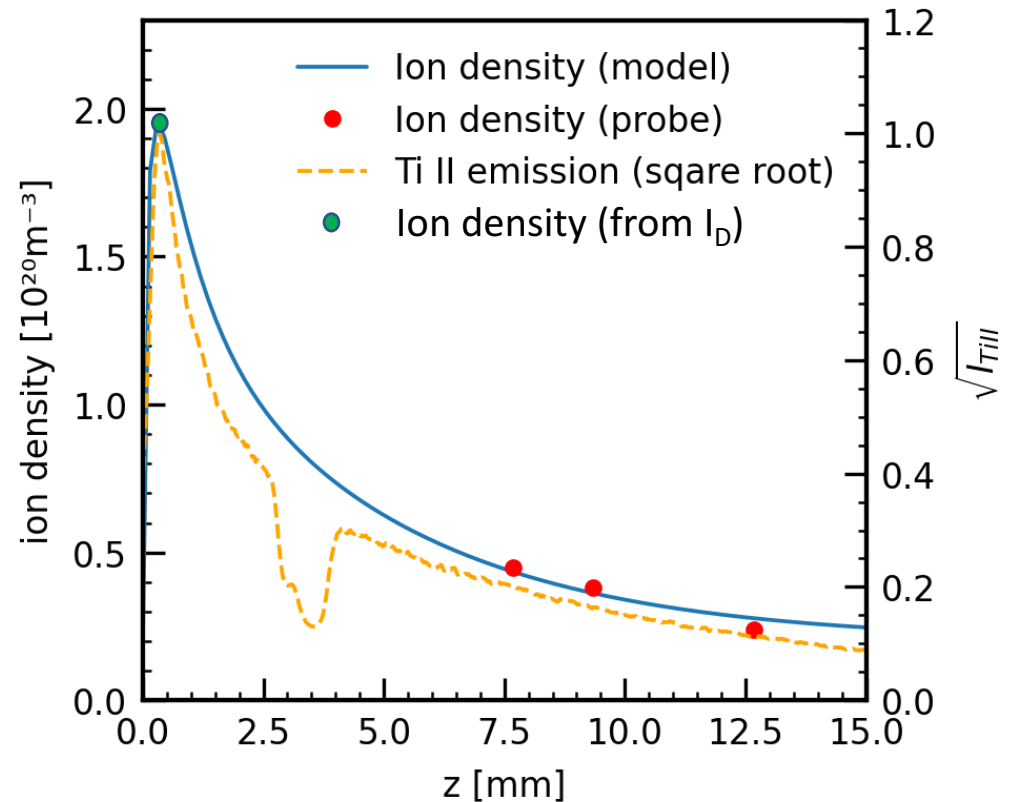
Model of Ion density:

Peak density from discharge current

Exponential decay to fit probe data

Shape similar as square root of Abel inverted emission

Suitable estimation reached

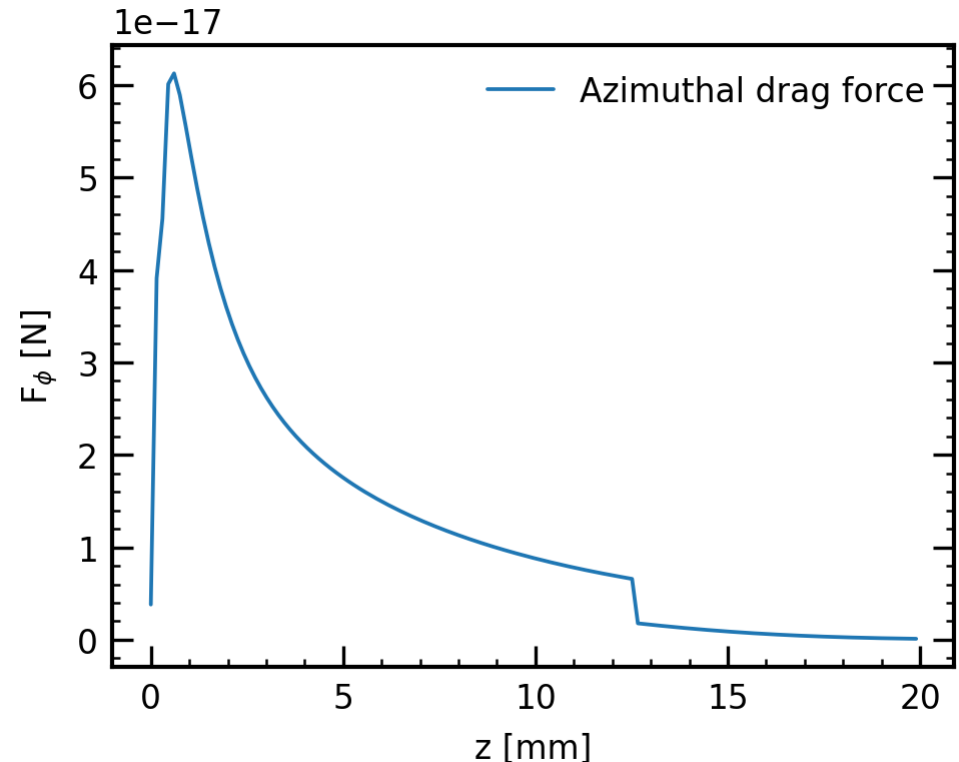


Part Three | Ion Density

$$F_{i,\phi} = -\eta e^2 n (v_D + v_{E \times B})$$

$$\mathbf{v}_D \equiv -k_B T \frac{\nabla n \times \mathbf{B}}{qnB^2}$$

$$\mathbf{v}_E \equiv \frac{\mathbf{E} \times \mathbf{B}}{B^2}$$



Shape of drag force mostly influenced by the density

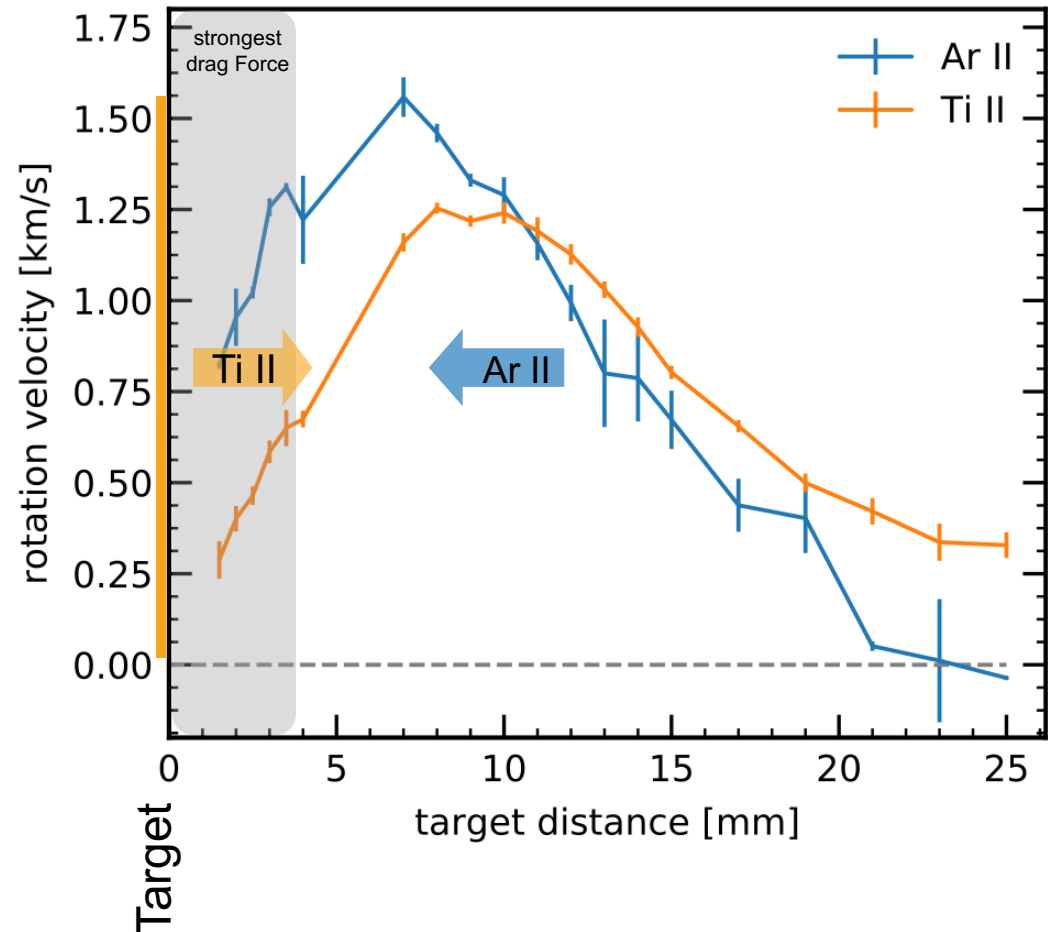
Measurement of Azimuthal Ion Velocity | OES

$z > 10$ mm:

- coupling of Ti II and Ar II due to collisions
- Slowing by Ion-Neutral collisions

$z < 10$ mm:

- Unknown residence time
- Ti II:
 - short residence time near target
 - accumulation of rotation velocity
- Ar II:
 - strong coupling with Ti II via collisions



Measurement of Azimuthal Ion Velocity | OES

Estimation of resulting velocity:

z = 8 mm:

- From [1] e⁻ confinement time:

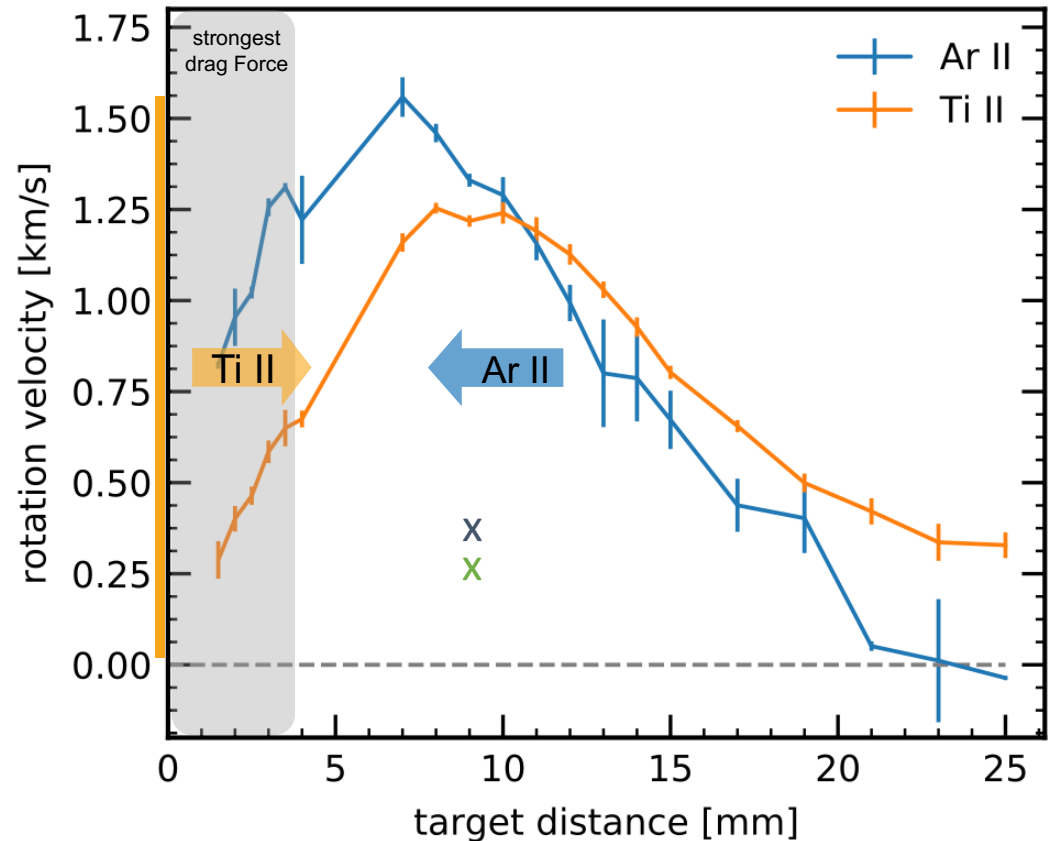
$$\tau_e(E) = 1.76 \times 10^4 \text{ eV} \frac{1}{E v_m(E)}$$

$$\tau_e(E) \approx 2 \mu\text{s}$$

- Same confinement time for electrons and ions

$$v_{\text{rot}} = \frac{F_{\text{drag}} \tau_i}{m_{\text{ion}}}$$

- Ti II: $v_{\text{rot}} \approx 250 \text{ m/s}$
- Ar II: $v_{\text{rot}} \approx 300 \text{ m/s}$

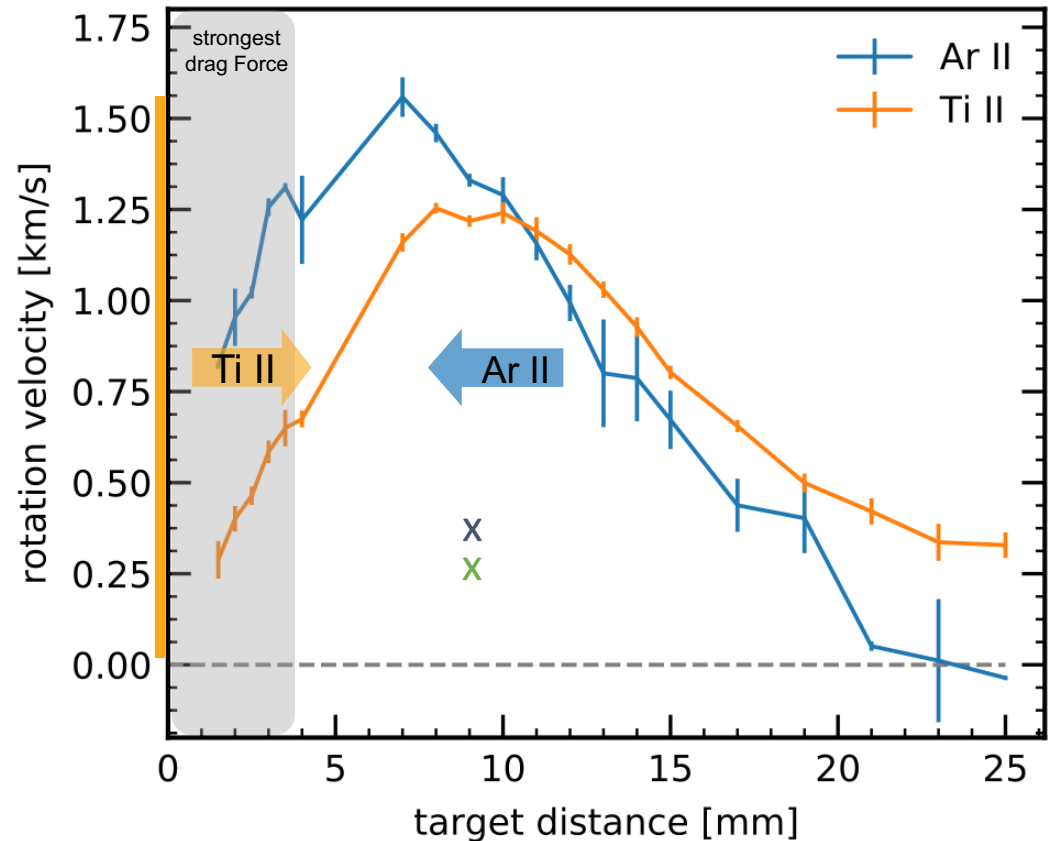


[1]: R. P. Brinkmann and D. Krüger, "Axisymmetric magnetically enhanced discharges described in terms of flux coordinates", *Plasmas* 27, 053504 (2020)

Measurement of Azimuthal Ion Velocity | OES

Problems:

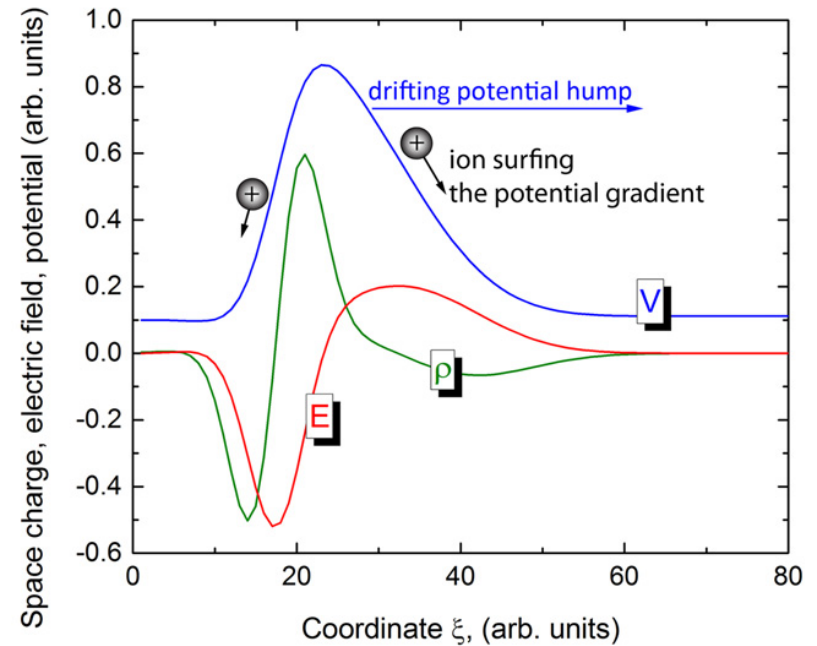
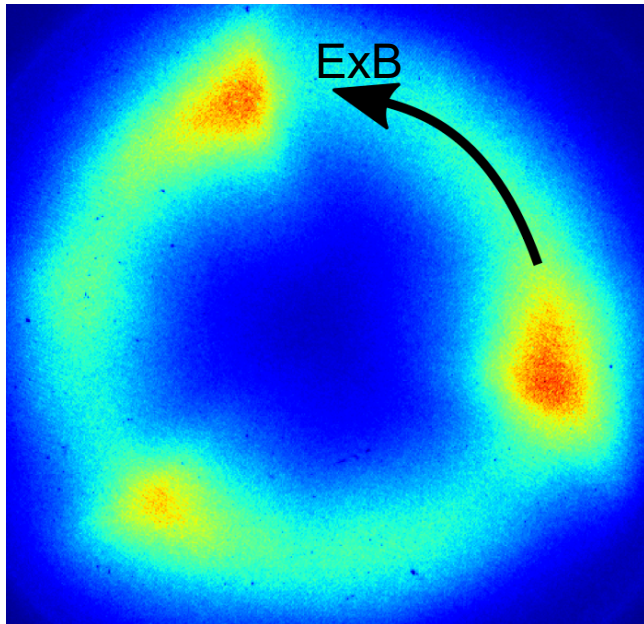
- accumulation of rotational velocity during movement in z direction
- Coupling of Ar II and Ti II due to collisions
- Spokes?



PART THREE

SPOKES

Introduction | Spokes



Panjan, Franz, Anders (2014) PSST **23** 025007

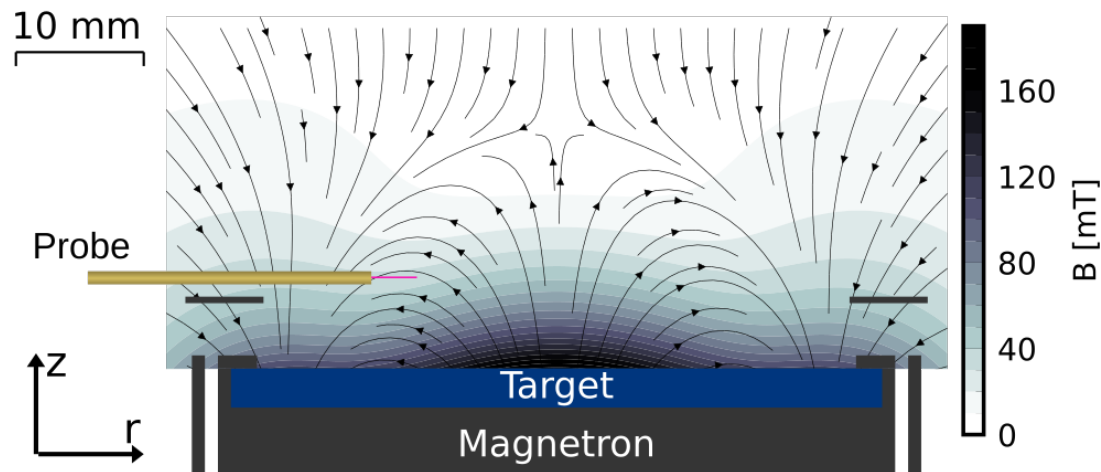
Rotating with 5 – 10 km/s in $E \times B$ direction

Spoke resolved probe
measurements needed

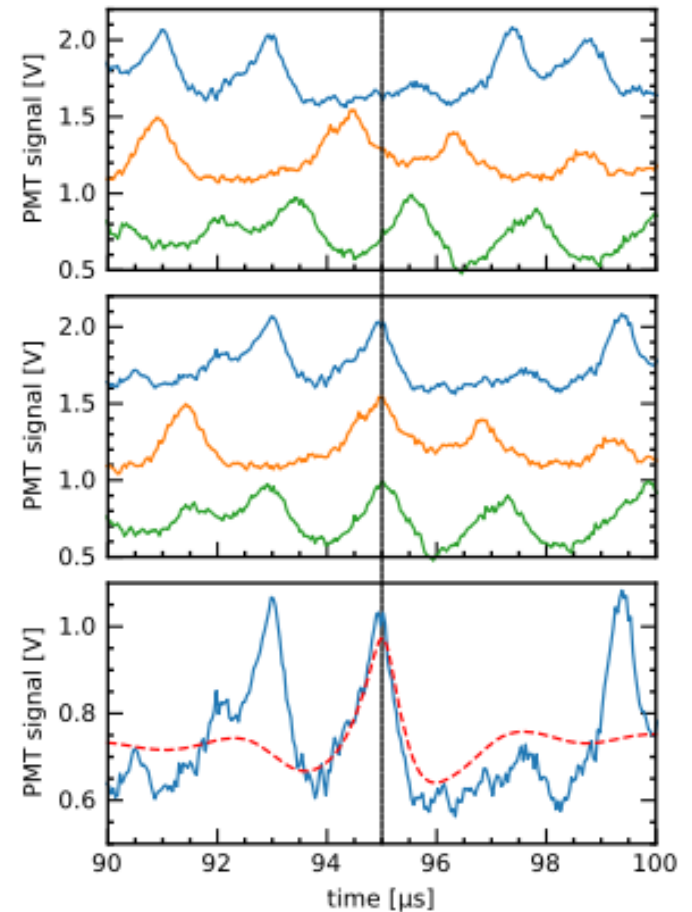
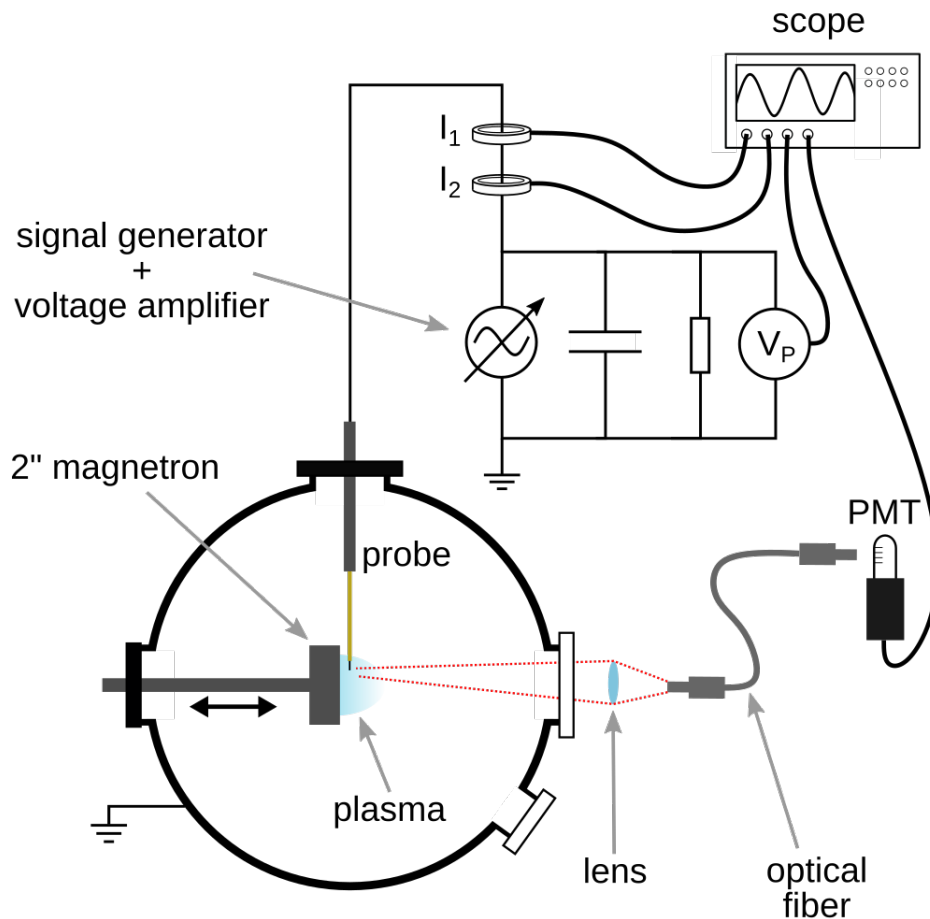
Part Two | Plasma Parameters? Langmuir Probe!

Two problems:

1. Magnetic fields
→ Magnetized Probe Theory
2. Spoke synchronization

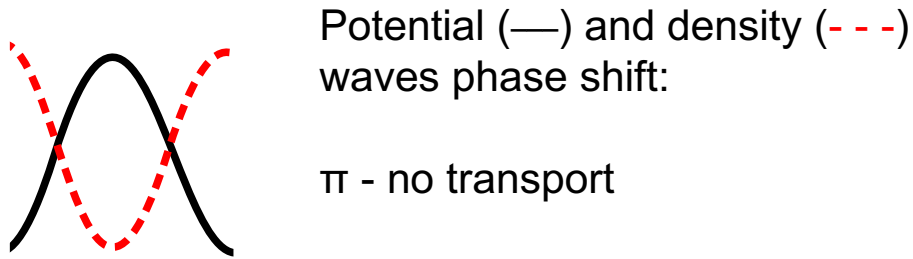


Part Two | Spoke Synchronization

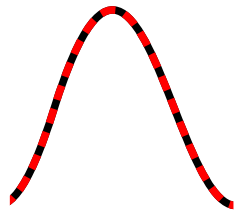


Time shift averaging to preserve spoke information under averaging

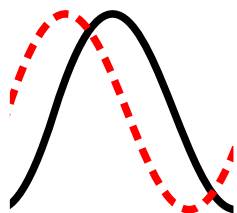
Part Two | Plasma Parameters



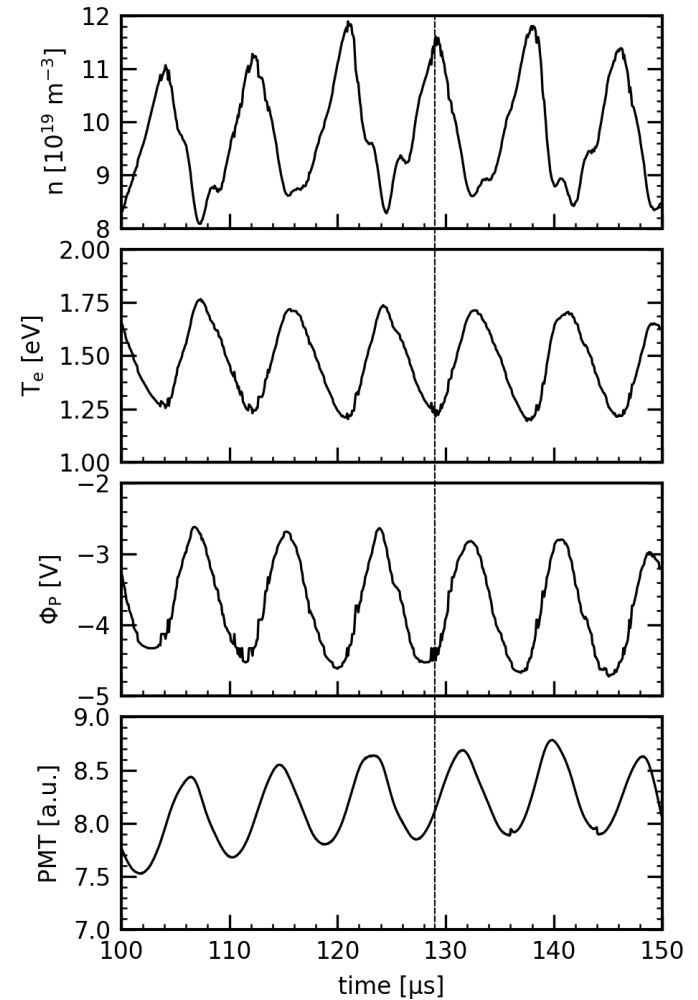
π - no transport



0 - no transport

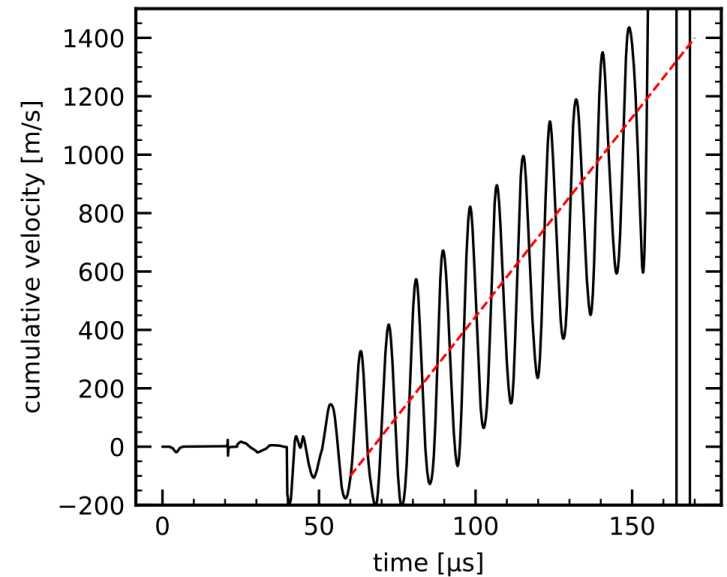
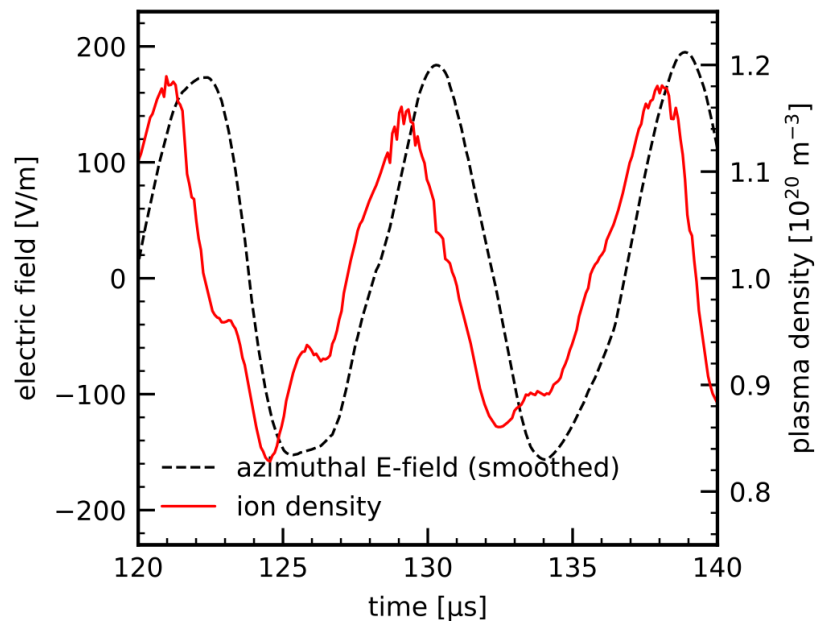


$\pi/2$ - transport!



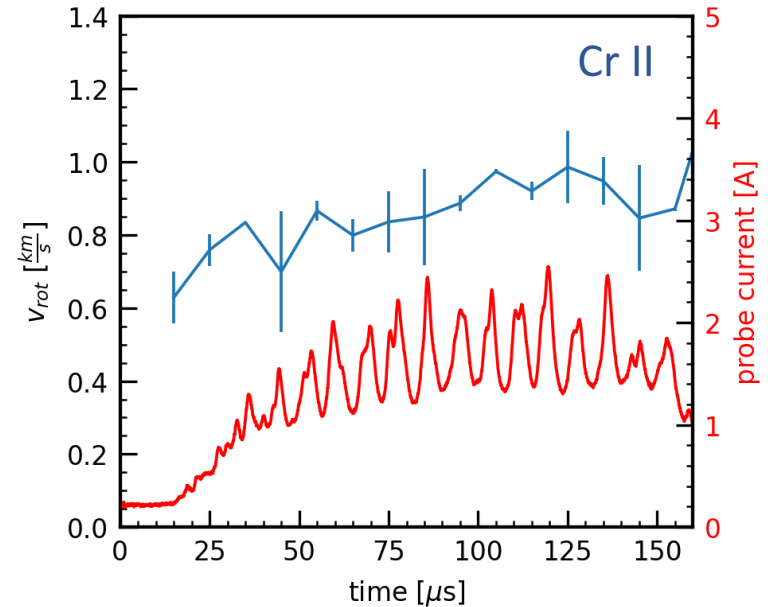
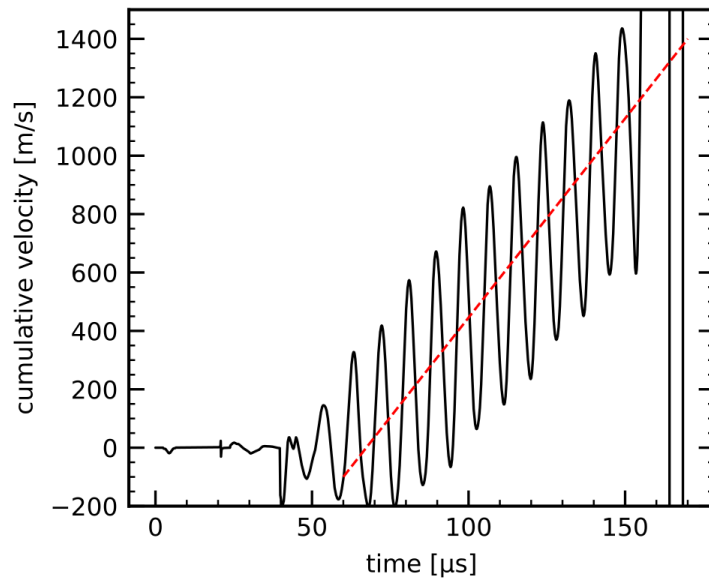
Part Two | Plasma Parameters

$$E_{\phi} = \frac{1}{v_s} \frac{dV_P}{dt} \quad \Gamma_{\phi} = n_i \frac{eE_{\phi}}{m} t$$



Spokes cause azimuthal ion rotation
in the order of 1 km/s, *eventually*.

Part Two | Plasma Parameters



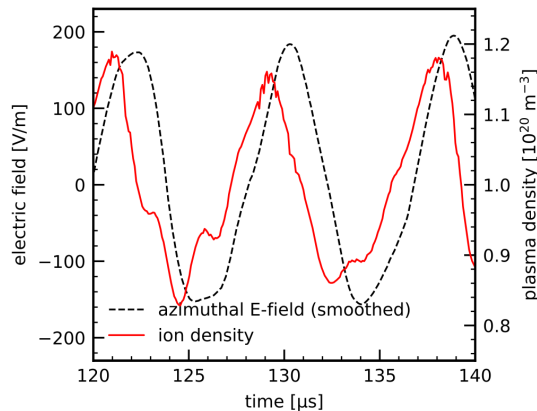
No cumulative effect over
whole discharge visible

CONCLUSION

HALL CURRENT OR SPOKES?

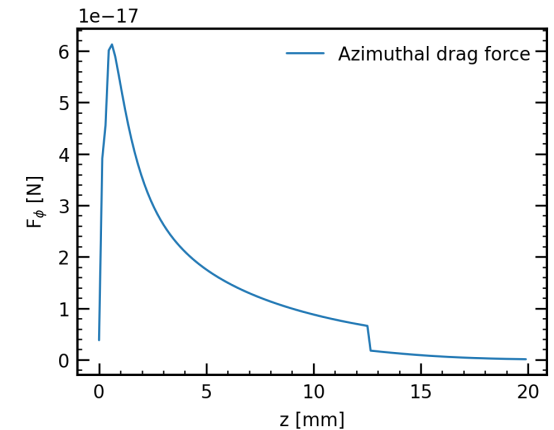
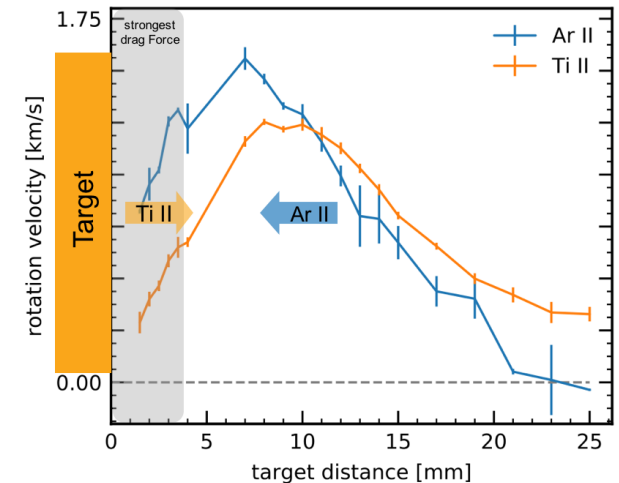
Summary

A better measurements of ion movement, based on high-resolution OES, is possible.
We find azimuthal velocities in the order of 1 km/s for ions.



Spoke-synchronized Langmuir probe measurements show strong oscillations in ion density, electron temperature and plasma potential. The shift between density and potential fluctuations strongly implies a net transport in the direction of spoke movement.

Estimations of the electron hall-current show almost constant values within the magnetic trap region. The shape of the azimuthal drag force therefore mimics the shape of the density. Having a maximum near the surface and decaying to further distances



Conslusion | Azimuthal Ion Movement in HiPIMS

**Hall current or
spokes?**

Probably both!

V Layes *et al* Appl Phys Lett (2017) **110** 081603

J Held *et al* PSST 2018 **27** 105012

J Held and A von Keudell: Plasma Chem Plasma Process (2019) **38**

J Held *et al* PSST (2020) **29** 025006